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1 Formation of cratonic lithosphere during the initiation of

2 plate tectonics

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9 ABSTRACT

10 The Earth's oldest near-surface material, the cratonic crust, is typically underlain 11 by thick lithosphere (>200 km) of Archean age. This cratonic lithosphere likely thickened 12 in a high compressional stress environment, potentially linked to the onset of crustal 13 shortening in the Neoarchean. Mantle convection in the hotter Archean Earth would have 14 imparted relatively low stresses on the lithosphere, whether or not plate tectonics was 15 operating, so a high stress signal from the early Earth is paradoxical. We propose that a 16 rapid transition, from heat-pipe mode convection to the onset of plate tectonics, generated 17 the high stresses required to thicken the cratonic lithosphere. Numerical calculations are 18 used to demonstrate that an existing buoyant and strong layer, representing depleted 19 continental lithosphere, can thicken and stabilize during a lid-breaking event. The peak 20 compressional stress experienced by the lithosphere is 3–4x higher than for the stagnant 21 lid or mobile lid regimes immediately before and after. It is plausible that the cratonic 22 lithosphere has not been subjected to this high stress-state since, explaining its long-term

stability. The lid-breaking thickening event reproduces features observed in typical
Neoarchean cratons, such as lithospheric seismological reflectors and the formation of
thrust faults. Paleoarchean 'pre-tectonic' structures can also survive the lid-breaking
event, acting as strong rafts, that are assembled during the compressive event. Together,
the results indicate that the signature of a catastrophic switch, from a stagnant lid Earth to
the initiation of plate tectonics, has been captured and preserved in the characteristics of
cratonic crust and lithosphere.

30 INTRODUCTION

31 The surviving remnants of the Archean crust appear to have been formed under 32 conditions of a low geothermal gradient (Burke and Kidd, 1978) best explained by a pre-33 plate tectonic, stagnant, heat-pipe mode of mantle convection characterized by vertical 34 tectonics and low convective stresses (Moore and Webb, 2013; Rozel et al., 2017). This 35 crust is underlain by the thickened remnants of lithosphere of similar age (Pearson et al., 36 1995; O'Reilly et al., 2001) that form the stable, buoyant continental tectosphere (Jordan, 37 1975). Long term stability of the cratonic lithosphere requires high strength in addition to 38 buoyancy to survive erosion by mantle convection (Lenardic and Moresi, 1999). We 39 propose that cratons formed above regions of lithospheric foundering, predicted to have 40 occurred when the Earth switched from the heat-pipe mode to a mobile-lid, plate tectonic 41 regime (Moresi and Solomatov, 1998). We argue that stresses of this magnitude have not 42 been reproduced since the Archean, explaining the long-term mechanical stability of the 43 cratons and lack of modern examples.

44 The initial stability of the cratonic lithosphere requires high compressive stresses
45 (Cooper et al., 2006) to overcome its strength and buoyancy to thicken it to the now

46	observed 200-300 km (Pasyanos et al., 2014), which cannot be generated by any steady-
47	state Archean mantle regime. Mobile lid convection involves heat-loss though strong
48	oceanic geothermal gradients, as a result of surface migration and recycling,
49	characteristic of plate tectonics. In its absence, convection operates in the stagnant lid
50	mode (Moresi and Solomatov, 1998), which is dominated by conduction through a
51	globally uniform boundary layer (lid). The heat-pipe regime is a variant of this stagnant
52	lid mode, in which upward melt transport into and through the lid is the dominant heat-
53	loss mechanism and can generate continental crust and highly depleted lithosphere
54	(Moore and Webb, 2013; Rozel et al., 2017), consistent with the observed cratonic
55	composition (O'Reilly et al., 2001). However, this lithosphere would initially have a
56	modest thickness, consistent with estimated melting depths of <100 km (Lee, 2014). An
57	additional thickening mechanism is therefore required.
58	Cratonic mantle lithosphere thickening appears to have occurred through thrusting
59	analogous to plate tectonic collisional zones (Bostock, 1998), also observed in
60	Neoarchean crustal deformation (Percival et al. 2006). Once plate tectonics begins, the
61	mobile lid regime is able to generate higher stresses than the heat-pipe regime, but
62	subduction zone and convective stresses would still be significantly lower in the warmer
63	Archean mantle compared to today (van Hunen and van den Berg, 2008; Sandu et al.,
64	2011). Cratonic lithosphere formed as a result of thickening by steady-state mobile lid
65	convection is potentially unstable to rising mantle stresses over time in a cooling Earth
66	(Cooper et al. 2007). A transient, high-stress event occurring during the early mobile lid
67	regime is therefore required to thicken and stabilize strong lithosphere that can sustain
68	stability for billions of years.

69	The last significant recorded cratonic deformation (Percival et al., 2006; Van
70	Kranendonk et al., 2007) typically occurred at a similar time (within 300 Ma) to the
71	proposed initiation of plate tectonics at \sim 3 Ga (based on data summarized by
72	Hawkesworth et al., 2017). This switch from the heat-pipe regime to mobile lid (plate
73	tectonics) would momentarily generate extremely high stresses, as the heat-pipe thermal
74	boundary layer was catastrophically recycled into the mantle (Moore and Webb, 2013).
75	This paper examines the manner in which proto-crust and depleted lithosphere thicken
76	and stabilize during lid collapse and the transition to mobile-lid convection, forming
77	cratonic nuclei of considerable strength.
78	DYNAMIC MODELING
79	We solve Stokes flow and temperature advection-diffusion using Underworld in a
80	2D, Cartesian domain 8700×2900 km in size (Moresi et al., 2007, see the GSA Data
81	Repository for further details of the numerical modeling methodology). We modeled
82	internally heated whole mantle convection (Fig. 2), with a Rayleigh number (Ra) of
83	5×10^8 , a viscosity range of five orders of magnitude and a resolution of 12 km in the
84	horizontal direction. In the vertical direction, the upper mantle grid spacing is refined to
85	6km.
86	Our models begin with a 72.5 km harzburgite layer with a density of 3269
87	kg/m^{-3} (relative to an assumed mantle density of 3310 kg/m^3), consistent with
88	xenoliths (O'Reilly et al., 2001). Our focus is the dependence of craton stability on the
89	deep lithospheric buoyancy and strength. The influence of the continental crust, which
90	could have formed previously through melting and recycling through the heat-pipe lid, is
91	generally ignored. The lithosphere has a finite plastic strength modeled by a depth-

92	dependent Von Mises criterion, representing the increased strength of olivine at high
93	pressure (Karato, 2010) and is assumed to be dry and melt-depleted compared to the
94	mantle which is 8x weaker. The stagnant lid can only be recycled if its strength is limited
95	by a low yield stress of 10-50 MPa (as in Moresi and Solomatov, 1998), which can be
96	generated through dynamic grain-size reduction and a subsequent switch to diffusion
97	creep (Rozel, 2012). A yield strength of 0.6 MPa/km is set, assuming that this weakening
98	has occurred. A short damage-dependent model demonstrates that the lid-breaking event
99	and craton stabilization occur in a similar way without this simplification (data
100	repository).
101	Melting buffers the mantle temperature, which is approximated by capping the
102	geotherm at the depth-dependent solidus. This produces an intermediate geotherm which
103	lies between the cool, extrusive heat-pipe end member (Moore and Webb, 2013) and the
104	warmer, intrusive end-member (Rozel et al., 2017). When the mantle cools down
105	sufficiently for melting to switch off (following Moore and Webb, 2013), the lithospheric
106	thickness is no longer controlled by the solidus and a large wavelength lithosphere-
107	asthenosphere boundary (LAB) slope can form, typical of the stagnant lid regime (Moresi
108	and Solomatov, 1998). This LAB slope is associated with high stresses which trigger an
109	overturn event, subjecting the harzburgitic layer to an episode of high compressive stress
110	(Figs. 1 and 2). The continental layer is shortened by <80% to form <300 km thick
111	'cratonic nuclei' during the foundering event.
110	CTDECC HICTORY AND CTADILITY OF THE ODATONIC LITHOCOUPDE

112 STRESS HISTORY AND STABILITY OF THE CRATONIC LITHOSPHERE

The lid-breaking event is associated with high horizontal compressional stresses
of ~150MPa in the buoyant lithosphere above the zones of mantle downwelling (Fig. 1).

These maximum stresses are 3–3.5x higher than the heat-pipe and stagnant lid regimes
(~30 MPa) prevailing immediately before the initiation of lid-breaking and the mobile lid
(plate tectonics) regime that follows.

118 The lid is recycled (except the continental layer) in fragments initially 1500km -119 2000 km wide. Each produces a local stress pulse which lasts 10-50Ma and contributes to 120 the formation of localized shear zones and thrust stacking within the cratonic lithosphere. 121 Thrusting occurs locally until the cratonic lithosphere stabilizes, incrementally 122 assembling a complete craton with up to three pulses over 100 - 350 Ma (central craton in 123 Figure 2B; one pulse has generated a single nucleus by 193 Ma and two more by 609 Ma), predicting in a wide range of craton formation ages globally. Due to the lithospheric 124 125 strength at high pressure and its cold pre-tectonic geotherm, stabilization occurs as soon 126 as the lithosphere is thickened to a critical depth (as in Cooper et al. 2006). No further 127 cooling or annealing mechanisms are required for stability, which would require longer 128 lived compressional regimes than typical for the hotter mantle. 129 Once all of the cold, dense lower lid has been recycled, no subsequent high stress 130 pulses occur, as a thick boundary layer can no longer form in the mobile lid (plate 131 tectonics) regime. Models are run for \sim 500 Ma after the lid-breaking event, in which the 132 modeled craton stress remains below $\sim 30\%$ of the thickening stress. The cratons remain 133 stable within the convecting mobile lid regime (e.g., craton stability after 609 Ma, Fig.

134 2B), despite an initial period of craton warming and mantle cooling.

The stress state evolution from the modeled Archean mobile lid convection to the modern Earth is estimated by scaling convective stress as a function of mantle temperature, which is estimated to have cooled by ~200 °C (Herzberg et al., 2010).

138 Mobile lid convective stress varies in proportion to the mantle viscosity (the Frank-139 Kamenetskii approximation of temperature-dependence is assumed, with E = 12) and the convective boundary layer thickness which scales with $Ra^{-\frac{1}{3}}$ (Turcotte and Oxburgh, 140 141 1972; Moresi and Solomatov, 1998). This scaling predicts modeled stresses which are 142 2.5x higher (~75 MPa) in the modern Earth than in the Archean (Fig. 1). The modern 143 cratonic lithospheric stress-state would therefore have only reached $\sim 60\% - 80\%$ of the 144 stress experienced during its Archean thickening event. 145 The maximum lithospheric stress during Earth's evolution over the last 3 Ga 146 should also be reflected in the evolution of maximum orogenic crustal thickness. The 147 maximum thickness of orogenic crust is limited to isostatic equilibrium with its local 148 compressional stress state. The largest crustal thickness is proportional to the highest 149 tectonic stress and provides an upper bound for the cratonic lithosphere stress state. 150 Dhuime et al. (2015) calculated the billion year evolution of juvenile upper-plate 151 continental crust thickness at subduction zones, based on geochemical isotopic 152 compilations, which we use as a proxy for tectonic stress. After ~3Ga, stresses increased 153 by 2.2x by ~ 1.5 Ga, which is initially more rapid than the scaling estimate (Fig. 1). 154 However, the subduction zone stress peaked at this time, such that it is still only \sim 70% of 155 the stress experienced during craton formation.

The simplified stress evolution is also consistent with the mantle traction stresses and lithospheric stress due to topography (<50 MPa) and net lithospheric stresses near subduction zones (~100 Mpa, Lithgow-Bertolloni and Guynn, 2004). This estimate depends on the lithosphere-mantle stress balance, without which subduction traction could theoretically generate <200 Mpa.

161 These stress estimates support our hypothesis that the high lithospheric stresses 162 which built the cratons have not been reproduced since their formation. The thermal 163 boundary layer may have thickened, but the cratonic nuclei are predicted to be 164 undeformed unless weakened. The importance of weakening cratons through slab-derived 165 fluid influx is debatable (Lee et al. 2011). The calculated stress evolution indicates that 166 weakening was critical in the past, though subduction zone traction stresses may have 167 risen to magnitudes for which this may be becoming less necessary in anomalously high 168 stress regions.

PRESERVATION OF PRE-PLATE TECTONICS CRUST 169

170 The simplified numerical model is an end-member in which the buoyancy and 171 strength of continental crust is negligible. A strong compressional stress-state consistent 172 with both shallow and deep thrusting is generated (Fig. 2 and Cooper et al., 2006). 173 Thrusting is recorded in Neoarchean crustal structures (e.g., Yilgarn and Superior 174 cratons; van der Velden et al., 2006; Percival et al., 2006), but its absence is a key 175 characteristic of Paleoarchean crust. Cratonic nuclei which formed >3 Ga (Pilbara and 176 Kaapvaal cratons) are instead dominated by buoyancy-driven domes (Van Kranendonk et 177 al., 2007), with no evidence of craton thickening. This can be explained by assuming that 178 lid breaking events generated the earliest thrusting and only occurred post-3Ga, which 179 agrees with the predicted onset of tectonics (Hawkesworth et al., 2017) and allows 180 Paleoarchean crust sufficient time to have 'stabilized'. Additional models (Fig. 2d) test 181 this by including an embedded 580 x 36.25 km crustal fragment of Paleoarchean crust, 182 stronger and more buoyant than the continental lithosphere with a plastic strength of 1.5 183 MPa/km and a density of 2877 kg/m^{-3} .

184 The buoyancy of the crustal block, applicable to felsic crust of any age, prevents 185 its deep burial. It is transported onto thick cratonic lithosphere (Fig. 2d) with no internal 186 plastic deformation and relatively modest viscous thickening (which would be further 187 reduced if strain-localization were modeled). Neoarchean crust was forming and 188 therefore still weak enough for the lid-breaking event to generate thrusts (preserved as a 189 synmagmatic combination of horizontal and vertical tectonics in the Yilgarn, Zibra et al., 190 2014). The structural and geochemical contrasts, for example between Yilgarn and 191 Pilbara crust, may then be the consequence of the large-scale convective regime they 192 formed in, whereas their lithospheric roots formed through similar thickening processes 193 during the initiation of tectonics.

EVIDENCE IN THE GEOLOGICAL RECORD

195 While our cratonization models are primarily designed to reconcile mantle regime 196 dynamics with the mechanical properties of the cratons, they also reproduce observed 197 features of cratons. Thrusting of the lithosphere and displacement of internal layering on 198 a large scale occurs in the models and can reproduce observed dipping lithospheric 199 reflectors (Bostock, 1998), as well as mid-lithospheric discontinuities (Calò et al., 2016). 200 Rapid Neoarchean events involving the amalgamation of multiple cratonic nuclei 201 (Percival et al., 2006) can also be generated by the lid breaking episode, which provides a 202 mechanism for suturing terranes in a relatively low stress tectonic environment. In the 203 models, individual mantle downwellings generate cratonic nuclei and suture them to form 204 larger assembled cratons. This has the combined effect of transporting crust laterally by 205 <4000 km, suturing terranes which would have contrasting geological evolution due to 206 their initial isolation, as observed in the Superior Province (Percival et al., 2006). Finally,

Neoarchean crust preserves horizontal tectonic features which can be explained by the
lid-breaking mechanism. The thrust sheets would reflect the compressional stress-state
(described earlier). Late-stage felsic magmatism (Percival et al., 2006) could be generated
as hydrated mafic crust was rapidly buried to >100 km (Fig. 1, Bédard, 2006), the
displacement also assisting with its recycling through the lithosphere.

212 The strongest evidence for craton formation during a catastrophic mantle 213 transition is the anomalous thickness, strength and billion year stability of the preserved 214 cratonic lithosphere, which places constraints on the evolving mantle dynamics since the 215 Archean. These features are reconcilable with craton formation during the proposed lid-216 breaking event, which would have generated a pulse of high stress, anomalous compared 217 to lithospheric stresses generated both during the Archean and today. The lid-breaking 218 event involves strong coupling between the asthenosphere, mantle lithosphere and crust, 219 providing opportunities to support our hypothesis with Archean crustal observations. The 220 simplified models demonstrate that our craton formation model is plausible, while there 221 is now the opportunity to more thoroughly explore the role of dynamic rheology and 222 crustal deformation. The lid-breaking model provides a plausible framework for 223 interpreting unique Neoarchean lithospheric and crustal features, which are otherwise 224 difficult to explain purely with secular adjustment to modern-style subduction and plate 225 tectonic models.

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- 328

329 **FIGURES**



331

332	Figure 1. Earth's evolving lithospheric stress state. The maximum stress during the
333	modeled stagnant lid regime is relatively low and increases gradually due to declining
334	heat-production. When widespread melting switches off, the stagnant lid can develop a
335	boundary layer slope and stresses increase rapidly, reaching the stagnant lid yield stress
336	and triggering the lid-breaking event and the initiation of plate tectonics, at \sim 3Ga
337	(Hawkesworth et al., 2017). Archean tectonic stresses are likely to be $\sim 60\%$ smaller than
338	today, based on the variation of crustal thickness over time (Dhuime et al. 2015) and
339	mantle convection scaling with $200^{\circ}C$ of cooling. The modeled lid-breaking event
340	generates pulses of anomalously high stress (ranging <300 Ma after lid-breaking) which
341	form strong cratonic nuclei. These anomalous stresses have not been reproduced since.
342	The lid-breaking event explains cratonic lithosphere formation in an otherwise low-stress
343	environment and the following stability at the billion year time-scale.

344



346 Figure 2. Typical model evolution from the stagnant lid regime (A), to the initiation of



348	change). The maximum stress evolution of the left craton is shown in Figure 1 (central
349	panel) and is representative. The data repository contains a movie of this model and the
350	stress evolution for the other two cratons. Individual thickening events take <30 Ma,
351	while multiple cratonic nuclei are amalgamated into the large central craton within 100
352	Ma. An enlarged frame of the thickened lithosphere (C) demonstrates the thrust structures
353	preserved after thickening. (D) An embedded Palaeoarchean crustal fragment, assumed to
354	have stabilized before the lid-breaking event, records relatively small shortening. This
355	explains the preservation of vertical tectonic structures, whereas significant thrusting
356	occurs in the weaker Neoarchean crust (shown schematically).
357	
358	1GSA Data Repository item 2018xxx, further details of numerical modeling
359	methodology and movie of typical model evolution, is available online at
360	http://www.geosociety.org/datarepository/2018/ or on request from
361	editing@geosociety.org.