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Modeling of High Frequency Current Transformer Based Partial Discharge Detection in High Voltage Cables

Xiao Hu, W. H. Siew, *Senior Member, IEEE*, Martin D. Judd, *Senior Member, IEEE*, Alistair J. Reid, *Senior Member, IEEE* and Bojie Sheng

Abstract--Partial discharge (PD) testing of high voltage (HV) cables and cable accessories has been implemented predominantly using high frequency current transformers (HFCTs) as PD sensors. PD currents initiating at PD sources are coupled onto cable conductors and travel away from the PD sources and will be detected by HFCTs installed at cable terminations. In this paper, based on combining finite-difference time-domain (FDTD) modeling and transfer function theory, a hybrid modeling approach is proposed to investigate the processes of PD coupling and detection involved in HFCT-based PD testing of HV cables. This approach allows exciting a PD event anywhere in FDTD models of the cables and predicting output from HFCTs some distance away. Implementation of the method is illustrated using an 11 kV XLPE cable. Moreover, a “direct measurement” method to obtain original PD pulses as the excitation source waveform is presented. The modeling approach introduced here will facilitate studies on the relationship between measured PD signals and those excited at PD sources, which can potentially give useful insight into the basic mechanisms behind PD detection in cables.

Index Terms--Finite-difference time-domain (FDTD), transfer functions, high frequency current transformer (HFCT), partial discharge (PD), power cables.

I. INTRODUCTION

PARTIAL discharge (PD) is indicative of local defects and degradation in high voltage insulation and is a popular insulation diagnostic tool. Measurement of PD has been widely carried out for high voltage (HV) cables to allow both detection and location of insulation defects within the cables. The PD sensor used in the measurement is frequently high frequency current transformers (HFCTs), in particular in on-line measurement [1]. The HFCTs are usually clamped at

cable terminations where the cable inner and outer conductors are accessible. PD current pulses flowing through the cable conductors will be detected by the HFCTs and converted to corresponding voltage pulses at the output terminals of the HFCTs.

PD currents initiating at PD source sites will be subject to significant changes before detection. Attenuation of the PD current pulses has been found as a function of propagated distance in the cables [2]. Furthermore, effects resulting from coupling of the PD currents from the PD sources onto the cable conductors and detection by HFCTs can also make appreciable differences. These effects make it difficult to interpret PD measurement results from the field. To investigate the relationship between PD currents excited at PD sources and PD measurement results from HFCTs in laboratory is both difficult and limited since some defect types cannot easily be reproduced experimentally. Exploring this problem through modeling is more versatile and can potentially provide useful insight.

Finite-Difference Time-Domain (FDTD) methods are used in electrodynamics as a numerical modeling technique to solve electromagnetic (EM) wave problems. Since the first application to EM waves by Yee in 1966 [3], it has grown in popularity for solving time-domain EM problems. The FDTD technique was found to be a useful tool within the field of studies on HV insulation systems and PD monitoring techniques in 2000 [4]. Since then, FDTD modeling has been used to investigate PD propagation problems in GIS [5-7] and transformers [8-10]. FDTD modeling has also been found effective in assisting with optimizing design and installation of UHF PD sensors for GIS and transformers [11,12]. In terms of applications to PD measurement for HV cables, some work has been seen using FDTD modeling to investigate sensitivity of field PD sensors [13,14] and high frequency properties of HV cables [15,16]. However, not much is reported in recent years. Probably due to the fact that modeling of frequency-dependent media such as the semi-conducting (semicon) layers of the cables is not supported by the basic FDTD method. Moreover, the widely used PD sensor like HFCTs is a kind of inductive coupling sensor which is usually made by winding many turns of thin wires onto a toroidal coil. This means that there are going to be too much fine details if proceeding to create physical representations of HFCTs in FDTD models. In this paper, a hybrid modeling approach based on combining FDTD modeling and transfer function

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theory is proposed to investigate the processes involved in HFCT-based PD testing of HV cables. Implementation of the modeling approach is illustrated by taking an 11 kV XLPE cable as an example. Steps and characteristics of the modeling approach are summarized to facilitate further applications.

II. FDTD METHOD AND TRANSFER FUNCTION THEORY

FDTD is based on numerically approximating and solving Maxwell's equations. By gridding a finite region with a sufficient number of electric (\mathbf{E}) and magnetic (\mathbf{H}) field cells, Maxwell's equations are solved in a time and space interleaved manner so that time-domain electromagnetic fields throughout the simulated region will be obtained. Being a time-domain technique, FDTD's solutions of which transient responses can be converted to the frequency-domain are able to cover a wide frequency range with one single simulation run, which is ideal for simulating PD pulses, a kind of broadband pulse. For carrying out practical FDTD simulations, both open source and commercial software packages are available. Commercial software usually has a powerful graphical user interface and is both able to deal with complex problems and easy to use. In this work, the commercial FDTD modeling package of XFDTD (ver. 7.1) [17] was employed for all the FDTD simulations and a Dell workstation with Xeon CPU of 2.67 GHz and 12 GB RAM have been employed to carry out the FDTD modeling.

Transfer functions are used in signals and systems to describe completely the input-output characteristics of a physical system [18]. In the frequency-domain, a transfer function $H(\omega)$ can be expressed as

$$H(\omega) = \frac{Y(\omega)}{X(\omega)} \quad (1)$$

where $Y(\omega)$ and $X(\omega)$ are the Fourier transform of the input and output of the system, respectively. Given the transfer function $H(\omega)$ of a system the output Y of the system to any input X can be obtained through multiplying $H(\omega)$ by the input in the frequency-domain and applying the inverse Fourier transform to the product if the output in the time-domain is sought.

III. FDTD MODELING OF HV CABLES

A 20 m 11 kV XLPE cable manufactured in 2010 according to BS 7870-4.10 was under investigation. The structure and dimensions of the cable are shown in Fig. 1, where r_1 and r_5 are the inner radius of cable components 1 and 5, respectively. Three-dimensional (3D) models of the cable were created using the built-in 3D modeling tool of XFDTD. The outer conductor of the cable is made of helically applied circular wires. Including the helical wires into the FDTD model can prove to be modeling intensive as well as computing intensive. There has been research considering detailed geometries of cable outer conductors to investigate their effects on high frequency properties of the cables [16,19,20], but in many other examples, these details are not considered. In this study, both inner and outer conductors were created as solid regardless of the detailed geometries.

In defining the electric properties of the materials for the

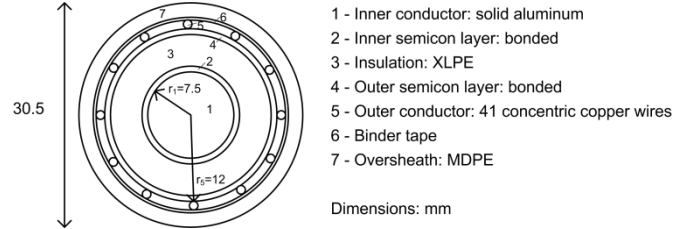


Fig. 1. Structure of the 11 kV XLPE cable sample. MDPE is acronym of medium-density polyethylene. The thickness of the inner semicon, insulation and outer semicon layers is 0.4 mm, 3.4 mm and 0.7 mm, respectively.

TABLE I
3-pole Debye function's parameters for material definition of the cable semicon layers within the FDTD model.

Parameter	Inner semicon	Outer semicon
σ_{dc} (S/m)	0.112	0.143
ϵ_{∞}	60.8	21.801
τ_1 (μ s)	0.0056	0.0098
ϵ_{s1}	435.16	74.412
τ_2 (μ s)	0.12	0.67
ϵ_{s2}	806.26	220.194
τ_3 (μ s)	1.23	114717
ϵ_{s3}	1200.84	220.181

cable structure within the FDTD model, properties of the inner and outer semicon layers were not measured because the semicon layers are bonded to the insulation and it is difficult to get proper samples for measurement. Instead, estimations of the semicon properties were made based on published semicon property data. The complex relative permittivity measured (at 25 °C) for the conductor screen and the insulation screen of a 12 kV XLPE cable in [21] was taken as the initial estimations. An optimization method [22] was used to further improve the accuracy of the estimations. TABLE I lists the final parameters used to define the semicon layers' properties in the FDTD model. The parameters are defined by a 3-pole Debye function [23] with which the complex relative permittivity of the semicon layers ϵ^* can be expressed as

$$\epsilon^*(\omega) = \epsilon_{\infty} + \sum_{i=1}^3 \frac{\epsilon_{si} - \epsilon_{\infty}}{1 + j\omega\tau_i} + \frac{\sigma_{dc}}{j\omega\epsilon_0} \quad (2)$$

where ϵ_{∞} is the permittivity at infinite frequencies and ϵ_{si} is the static permittivity from the i th Debye relaxation process and its relaxation time is τ_i , and σ_{dc} is the DC conductivity, and $\epsilon_0 = 8.854 \times 10^{-12}$ F/m is the permittivity of air. The cable conductors and XLPE insulation within the FDTD model were specified as perfect conductors and an ideal dielectric with $\epsilon_r = 2.3$ [24] and $\sigma = 0$ S/m, respectively. This means that the skin effect in the conductors and the dielectric loss were not considered in the FDTD model. It is usually not practical to consider the skin effect in conductors within FDTD models because the skin effect decreases the wavelength in conductors with frequency, and at very high frequencies the wavelength can become too small to satisfy a general rule of thumb of the FDTD method, i.e., 10 FDTD cells per wavelength. The dielectric loss in the cable insulation was not considered because the employed modeling package did not provide an appropriate material definition [17]. However, for XLPE-insulated cables, it is the semicon layers that are dominating the high frequency characteristics of the cables [25,26]. It will

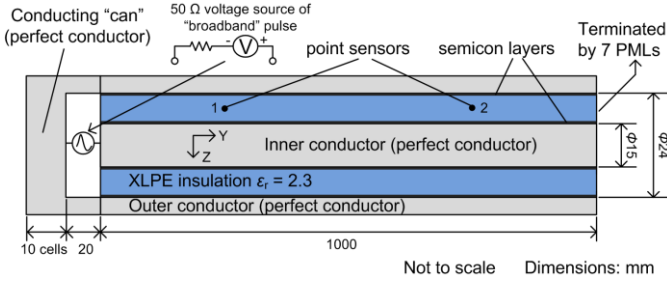


Fig. 2. Axial cross-section of the FDTD model of the 11 kV XLPE cable. Two point sensors 1 and 2 axially spaced for some distance were put within the insulation to record electromagnetic fields at corresponding positions.

be shown later that while only the losses in the semicon layers were represented, expected agreement between the prediction of the FDTD model and measured high frequency attenuation of the cable was achieved within a frequency range of 200 MHz.

IV. EVALUATING ACCURACY OF THE FDTD MODEL

A. Modeling of TDR Measurement

In order to evaluate the accuracy of the FDTD model, modeling of time-domain reflectometry (TDR) measurement on the cable was carried out. Fig. 2 shows an axial cross-section view of the FDTD model where a TDR input module was created at one end of the cable. The other end of the cable was terminated with 7 perfect matched layers (PMLs) which are one of the commonly used boundary conditions of the FDTD method and were used here to avoid reflection from the open end. The TDR source was simulated using a lumped 50 Ω voltage source with the waveform of a “broadband” pulse of which an example is shown in Fig. 3. The “broadband” pulse is a specific waveform definition in the modeling package that is a Gaussian pulse with a bandwidth able to cover the highest frequency content possible for the specific FDTD space [17]. The highest possible frequency is subject to the maximal cell dimension L_{\max} in meters. For this model, in Cartesian coordinates, an FDTD cell size of $(\Delta x, \Delta y, \Delta z) = (0.20, 20, 0.20)$ mm was applied making $L_{\max} = \Delta y = 20 \cdot 10^{-3}$ m. The highest possible frequency can be calculated as $f_{\max} = c / (10 \cdot L_{\max})$ where c is the speed of light in free space, i.e., $3 \cdot 10^8$ m/s, and thus $f_{\max} = 1.5$ GHz. The spectrum of the “broadband” pulse will have less than 10 dB attenuation at the highest frequency f_{\max} . With electromagnetic fields recorded by both point sensors (refer to Fig. 2) in FDTD simulations, the cable propagation constant γ_c can be worked out (e.g., up to f_{\max}). The electromagnetic fields at the two points, EM_1 and EM_2 , are related by

$$\text{FFT}[EM_2] = \text{FFT}[EM_1] \times e^{-\gamma_c(\omega)(L_2-L_1)} \quad (2)$$

where FFT is the fast Fourier transform, and L_1 and L_2 are axial distances from point sensors 1 and 2 to the TDR input, respectively. $L_1 = 0.1$ m and $L_2 = 0.9$ m were used in this model.

B. FDTD Simulation Results

The FDTD simulation was carried out requesting results of a sufficient length of time so that both the point sensors can capture the complete TDR input pulse. Fig. 4 shows the radial

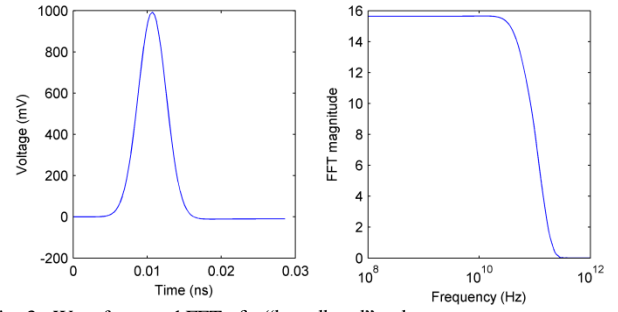


Fig. 3. Waveform and FFT of a “broadband” pulse.

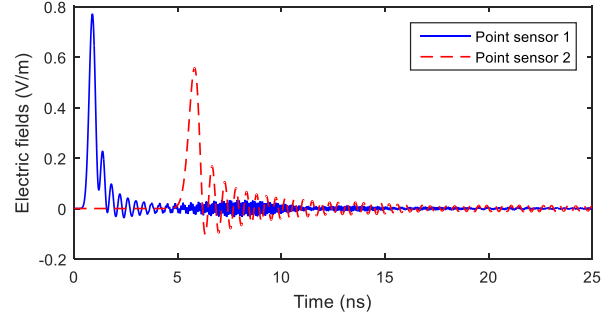


Fig. 4. Electric fields recorded by point sensors 1 and 2. The simulation was carried out with a cell size $(\Delta x, \Delta y, \Delta z)$ of $(0.20, 20, 0.20)$ mm and took 1 hour, 48 minutes to complete.

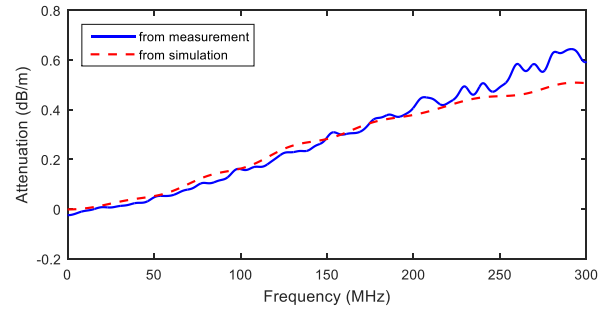


Fig. 5. Cable attenuation represented by the FDTD model of the 11 kV XLPE cable.

direction component of electric fields recorded by the point sensors. In both electric field traces, the first pulse was followed by a few smaller pulses, which was caused by impedance mismatch between the TDR input module and the cable within the model. There was some kind of “ringing” noise starting from about 5 ns in the trace recorded by point sensor 1. This was a result of frequency components of the “broadband” pulse interacting with the FDTD cells. FFT analysis of the electric field suggested that the “ringing” noise frequency was centered at 10 GHz, which was well above the frequency range of interest here and therefore was not of concern.

With the electric fields recorded by point sensors 1 and 2 and using (2), the cable propagation constant represented by the FDTD model was calculated. Fig. 5 shows the obtained cable attenuation, i.e., the real part of the calculated propagation constant, against that obtained using a calibrated time-domain reflection measurement method [27]. General agreement between the prediction of the FDTD model and the

measurement result was achieved within 200 MHz higher than which the difference between the simulation and measurement results became more and more evident. The fact that the skin effect in the cable conductors and the dielectric loss in the cable insulation cannot be accounted for by the FDTD model will contribute to the difference. However, it is believed that the difference was mostly caused by that the Debye function can only predict the semicon properties with acceptable accuracy at a limited range of frequencies. Simulation results of another FDTD model created for an 11 kV EPR cable showed that agreement in the cable attenuation between the FDTD model and the corresponding measurement was valid only up to 50 MHz. This was attributed to that good estimations of the cable's semicon properties were not obtained and moreover the EPR insulation had a dielectric loss tangent of 0.004 which was ten times of that of the XLPE insulation and was not considered in the FDTD model.

V. MEASUREMENT OF EXCITATION SOURCE WAVEFORM

To obtain "original" PD pulses that can be used as the source waveform to excite a PD in the FDTD model, "direct measurement" of PD [28] in a cable was performed. The measurement setup is plotted in Fig. 6a. The actual connection is shown in Fig. 6b. The experiment was conducted on a 1.5 m 11 kV EPR cable sample manufactured in 1983 according to BS 6622. Proper terminations were installed on both ends of the cable sample to avoid termination corona discharges. A potential protrusion defect was introduced into the cable insulation by using a hypodermic needle. The center of a 50 Ω coaxial cable was connected with the needle and the other end of the coaxial cable was connected to the 50 Ω input of an oscilloscope. Since it was a direct connection between the energized sample and the oscilloscope, there was a potential risk of damage in the event of flashover. Therefore, energizing at lower voltages like 5 kV and 6 kV was performed. PD current pulses resulting from the needle were recorded. Fig. 7 shows one of the recorded traces consisting of one single pulse. The single pulse trace can be used to define the waveform of the excitation source in the FDTD model.

VI. MODELING OF HFCT-BASED PD DETECTION IN CABLES

A. Hybrid Modeling Approach

A hybrid approach to modeling of HFCT-based PD testing of HV cables was developed. The approach is based on using FDTD modeling to investigate coupling of PD source currents onto the cable conductors and transfer function theory to work out HFCTs' voltage output while being input by the PD-induced currents on the cable conductors. Fig. 8 shows this approach as five steps:

Step 1) is the measurement procedure detailed in Section V.

Step 2) - An FDTD model of a cable under investigation is created and one of the recorded PD waveforms from Section V is used to define the excitation current as a data file within the model. Simulations are then carried out to characterize the effects caused by coupling of PD currents from the PD source location onto the cable outer conductor.

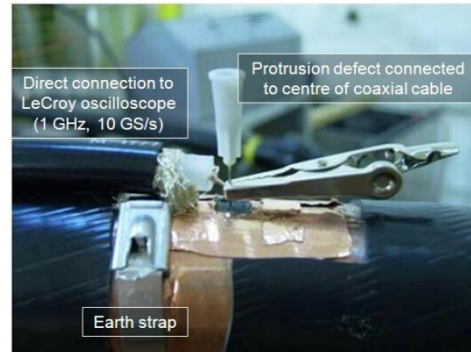
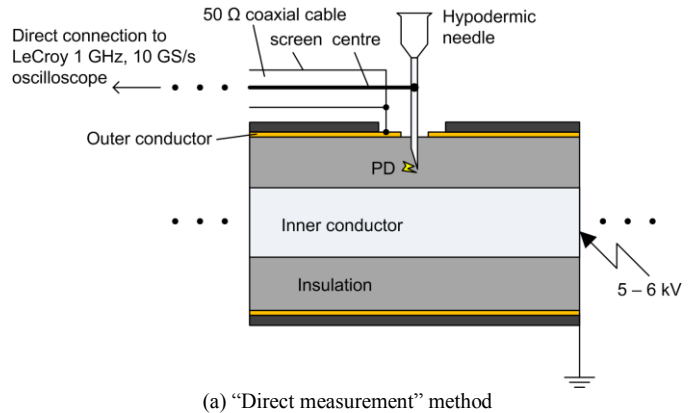


Fig. 6. "Direct measurement" method to obtain "original" PD pulses.

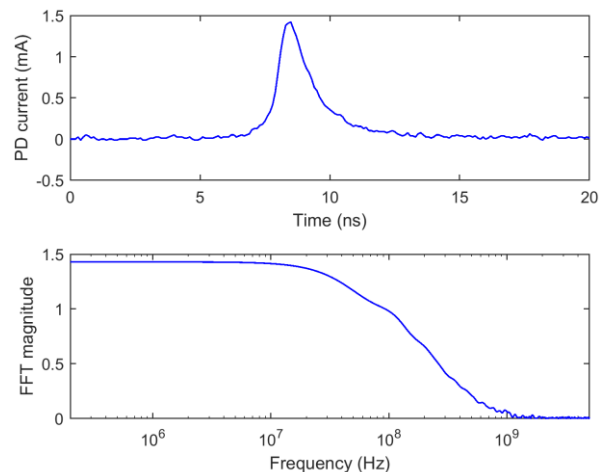


Fig. 7. A single positive pulse recorded using the "direct measurement" method and FFT of the pulse. 28 pulses with similar magnitude and shape were recorded. Average rise-time of the pulses is 0.94 ns. Average full width at half magnitude (FWHM) is 1.3 ns. Average -3 dB bandwidth is 96.4 MHz.

Step 3) - A method of measuring currents flowing on conductors within FDTD models was developed. The method requires that a sensor loop is created out of point sensors around the cable outer conductor within the FDTD model. The point sensor is used to record magnetic fields at a certain point of the model.

Step 4) - The recorded magnetic fields from these point sensors will be used to derive currents flowing on the cable outer conductor by using Ampere's law. It should be noted that the correct vector component of \mathbf{H} is used at each point so that the numerical integration does sum the contributions to \mathbf{H}

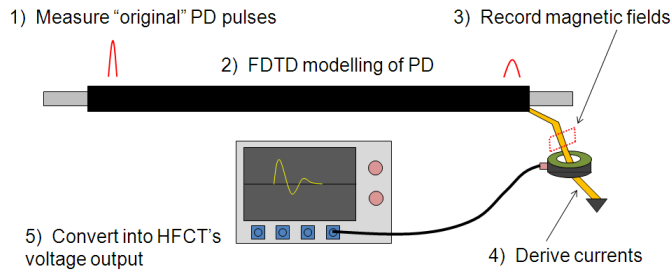


Fig. 8. Modeling of HFCT-based PD detection in high voltage cables.

around a closed loop.

Step 5) - The complex transfer function of a HFCT can be obtained by simultaneously measuring the HFCT's input and output and then dividing FFT of the output by FFT of the input. The step and pulse responses of the HFCT are measured to construct the low and high frequency parts of the transfer function respectively [29]. Through multiplying the obtained HFCT transfer function by FFT of the currents derived in step 4) and using IFFT, the HFCT's time-domain PD pulse output voltage can be determined.

B. Illustration of the Modeling Approach

As an example, the proposed modeling approach was implemented based on the FDTD model of the 11 kV XLPE cable evaluated in Section IV. To simulate PD detection in the cable, some changes were made, and the updated model is shown in Fig. 9. Both ends of the cable within the model were terminated with 7 PMLs to simulate matched loads. The cable outer conductor was curtailed by 0.2 m at one end and instead an extended outer conductor was made for the remaining 0.2 m to simulate an earth strap on which a HFCT is clamped. According to steps 3) and 4) of the modeling approach, a closed loop consisting of point sensors was created on the extended outer conductor to record magnetic fields. A PD

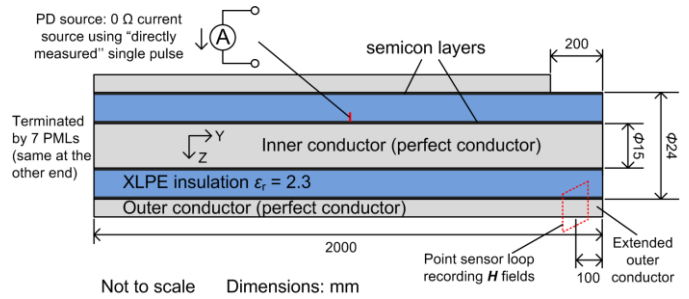


Fig. 9. Axial cross-section of the FDTD model of the 11 kV XLPE cable for demonstrating the modeling approach. The cell size used in the FDTD simulation was $(\Delta x, \Delta y, \Delta z) = (0.20, 20, 0.20)$ mm.

source was created as a current filament represented using a current source with 0 Ω source impedance. The PD source was defined in the FDTD model on the inner conductor in the middle of the cable. The single pulse trace shown in Fig. 7 was used as the PD source current waveform. After the simulation, magnetic fields recorded by the point sensor loop were integrated to obtain the current flowing on the extended outer conductor. The derived current in combination with a measured HFCT transfer function, converting to the HFCT's voltage output was carried out by following step 5).

Piecing together the PD source current, the outer conductor current, the HFCT transfer function and the HFCT's voltage output, an illustration of the different stages of the modeling approach was made in Fig. 10. For steps 3) and 4), it shows that the derived current pulse on the extended outer conductor had a negative polarity while a positive current pulse was supposed to be flowing on the inner conductor. This was a result of (refer to Fig. 9) the positive direction of the current source flowing into the inner conductor. The obtained HFCT's voltage output had some oscillation after the first pulse, which was caused by the peak at about 50 MHz in the magnitude

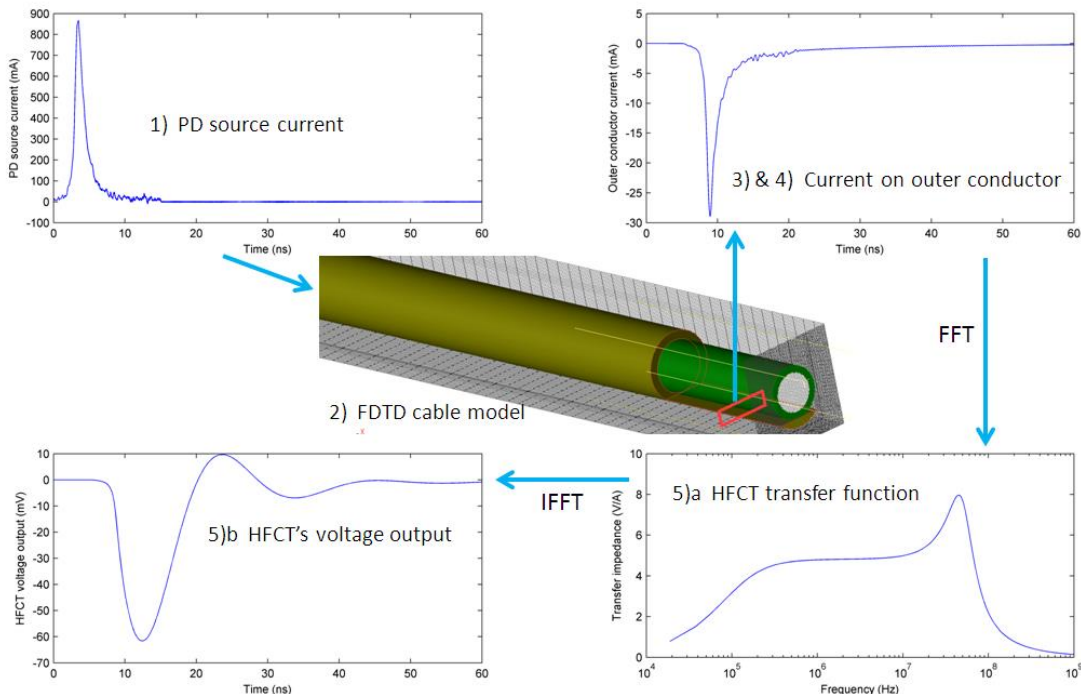


Fig. 10. Illustration of the whole modeling methodology.

response of the HFCT transfer function. The first pulse was widened compared with the current pulse on the outer conductor because of the limited bandwidth of the HFCT.

VII. RELEVANT EXPERIMENTAL RESULTS

A pulse-injecting test was carried out to investigate attenuation and distortion of PD pulses in a 35 m 11 kV EPR cable sample. Fig. 11 shows the measurement setup. The pulse was injected from one end of the cable and measured at the other end. The injected pulse was first of all measured when the 22.3 m test cable was disconnected from the EPR cable and connected to channel 2 of the oscilloscope, and then the measurement was performed with the connection shown in Fig. 11. Measurement results are shown in Fig. 12. The pulse measured at the other end of the cable had magnitude less than one fourth of the injected pulse and the rise-time increased by about seven times. While the increasing rise-time was due to propagating through the EPR cable, the magnitude attenuation was caused not only by propagation effects but also by impedance mismatching at the interfaces between the test cables and the EPR cable. This test did not include measurements using the HFCT but another test [29] to characterize the transfer function of the HFCT made use of the same avalanche pulse generator to excite the HFCT and the HFCT output (as shown in Fig. 13) shows clear difference from the input pulse. The experimental results presented here are not meant to be a direct comparison with the modeling detailed in Section VI but some similarities between the experimental and modeling results, such as attenuation of pulse magnitude and distortion of pulse shape, are indeed observed. The experimental results show some typical effects that contribute to the change of PD signals when HFCT-based PD testing of cables is conducted. It has been envisaged that these effects can be further investigated using the proposed modeling approach.

VIII. DISCUSSION

A. Potential Advantages of the Modeling Methodology

Although the illustration of the modeling approach in Section VI is basic, the purpose is to show how the approach will work on an “original” PD current pulse and prove effectiveness of the model. Some potential advantages of the modeling approach are summarized as follows:

1. By creating 3D models of the cables, both the cable cross-sectional and axial details can be characterized. Therefore, modeling earth connections, terminations, joints and various defects are all possible.
2. By exciting a PD pulse anywhere in the created models, the PD-induced currents on the cable inner and outer conductors some distance away from the PD launching site can be predicted. This permits variations in physical parameters of the PD that would not be possible in practical experiments.
3. Instead of physically modeling HFCTs in FDTD models, implementing the effect of HFCTs using the transfer function method is more practical and makes it easier to compare different HFCTs.

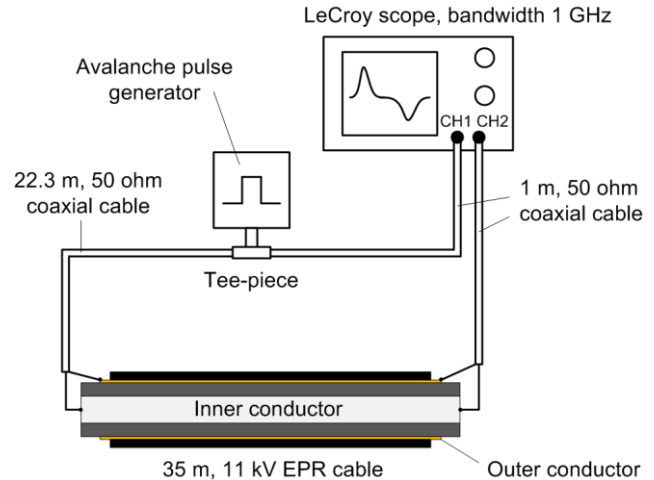
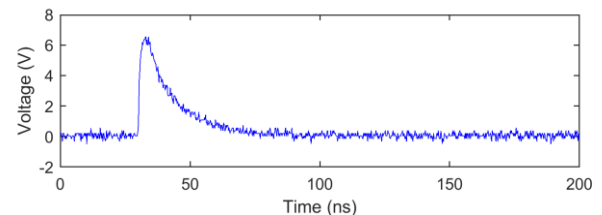
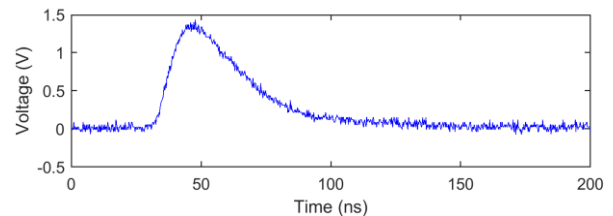


Fig. 11. Measurement setup for injecting pulses into the cable. The avalanche pulse generator can generate pulses with sub-nanosecond risetime.



(a) The injected pulse with a 1.5 ns rise-time and 6.53 V magnitude.



(b) The 1st pulse propagating through the EPR cable. The rise-time and magnitude of the pulse are 10.8 ns and 1.43 V, respectively.

Fig. 12. Measurement results of the pulse-injecting test.

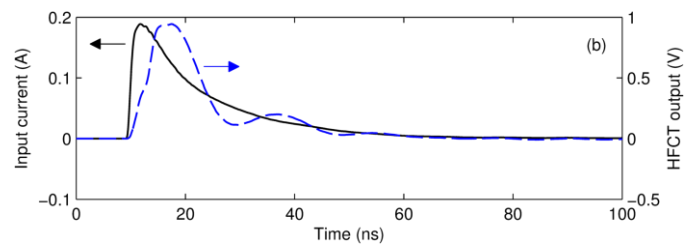


Fig. 13. The current input and the voltage output of the HFCT. The current pulse was generated with the same avalanche pulse generator as that used in the pulse-injecting test.

B. Implementation of the Approach

Implementing the hybrid modeling technique generally requires “direct measurement” of “original” PD pulses, creating FDTD models of HV cables and characterization of HFCT transfer functions. Regarding polymeric-insulated HV cables, one of the challenges faced is that the material properties of the cable inner and outer semicon layers must be known. To measure them can involve sophisticated test setup and devices, and furthermore it is difficult to remove suitable semicon samples from the cables if the semicon layers are bonded to the cable insulation, which is quite normal for

modern polymeric-insulated cables. Instead of practical measurement, a possible solution of using published semicon properties and an optimization method can be adopted. This solution was found effective in the modeling of the 11 kV XLPE cable where the cable attenuation predicted by the FDTD model agreed appropriately with the measurement result. However, it is worth noting that the FDTD technique being a time-domain method has a relatively limited capability in dealing with frequency-dependent materials like the semicon layers of polymeric-insulated cables. For modeling the semicon layers, the most appropriate material definition provided in the software package XFDTD is based on the Debye function which can only fit the semicon properties at limited frequencies. As a result, the frequency range where the FDTD models can predict the cable attenuation will be limited, e.g., within 200 MHz for the 11 kV XLPE cable. While the modeling work presented in this paper is based on a XLPE cable the principles are readily applicable to other cables, such as EPR cables and paper-insulated cables. However, it must be emphasized that the dielectric losses in the insulation of these cables are substantially different to those of XLPE cables where the cable attenuation is predominantly caused by the losses in the semicon layers. Therefore, for modeling of EPR cables and paper-insulated cables, it would be necessary to construct the model appropriately to represent the predominant loss mechanisms so that sufficient modeling accuracy can be achieved.

IX. CONCLUSIONS

An FDTD model of an 11 kV XLPE cable was created in this study. The semicon layers of the cable were created out of the Debye medium in the FDTD model and the parameters of the corresponding Debye function were obtained based on published data of cable semicon properties. The simulation results suggest that the cable attenuation represented by the FDTD model can achieve appropriate agreement with the measurement result within a frequency range of 200 MHz. A hybrid modeling approach based on combining FDTD modeling and transfer function theory was developed to investigate the PD coupling and detection processes involved in HFCT-based PD testing of HV cables. Implementation of the method has been detailed and illustrated. Furthermore, a “direct measurement” method to obtain “original” PD pulses necessary for the modeling has been described. The resulting modeling approach can be used to investigate the relationship between measured PD signals and those excited at PD sources, which is expected to improve our understanding of PD measurement for cables and interpretations of the measurement results. Future work will involve modeling of cable joints and terminations with more complex structures. Moreover, the possibility of applying the inverse transfer function to infer the PD source current pulse based on measurements at a HFCT sensor will be explored, which would improve diagnostic capabilities for measurements in the field.

X. REFERENCES

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