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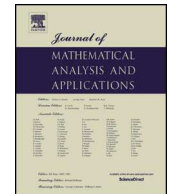




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Fractional immigration-death processes

Giacomo Ascione^a, Nikolai Leonenko^b, Enrica Pirozzi^{a,*}

^a *Dipartimento di Matematica e Applicazioni "Renato Caccioppoli", Università degli Studi di Napoli Federico II, 80126 Napoli, Italy*

^b *School of Mathematics, Cardiff University, Cardiff CF24 4AG, UK*

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ABSTRACT

In this paper we study explicit strong solutions for two difference-differential fractional equations, defined via the generator of an immigration-death process, by using spectral methods. Moreover, we give a stochastic representation of the solutions of such difference-differential equations by means of a stable time-changed immigration-death process and we use this stochastic representation to show boundedness and then uniqueness of these strong solutions. Finally, we study the limit distribution of the time-changed process.

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1. Introduction

Birth-death processes constitute an important class of continuous time Markov chain (CTMC). They are widely used, for instance, in population and evolutionary dynamics (see [35,36]), queueing theory (see [43]) and in epidemiology (see [3]). A complete classification and characterization of birth-death processes is due to Karlin and McGregor, whose papers [18,19] are the starting point of the study of family of classical orthogonal polynomials linked to such processes.

Classical orthogonal polynomials are widely used to study the solutions of Kolmogorov equations as in the case in which the state space of the process is continuous, as well as in the discrete one. In the continuous case, the families of classical orthogonal polynomials are used to give a spectral decomposition of Kolmogorov equations induced by the generators of Pearson diffusions [14]. In the discrete case, the discrete analogue of Pearson diffusions is given by a certain class of solvable birth-death processes. Moreover one can associate to any family of classical orthogonal polynomials of discrete variable another particular family, called the dual family [34]. In some cases, a family of classical orthogonal polynomials of discrete variable could be in duality with itself: in this case it is called *self-dual* family [42]. Among self-dual families, the simplest one is the family of Charlier polynomials, whose self-duality is induced by the following formula

* Corresponding author.

E-mail addresses: giacomo.ascione@unina.it (G. Ascione), leonenkon@cardiff.ac.uk (N. Leonenko), enrica.pirozzi@unina.it (E. Pirozzi).

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$$C_n(x, \alpha) = C_x(n, \alpha), \quad n, x \in \mathbb{N}_0,$$

called duality formula for Charlier polynomials (see Section 3 for the definition of Charlier polynomials). Charlier polynomials are really useful in the study of immigration-death processes (or $M/M/\infty$ queues) [42] and in their general version on 1-dimensional lattice, called Charlier processes [1]. Indeed, one can give a spectral decomposition of the strong solutions of Kolmogorov equations induced by the generator of the immigration-death processes in terms of such polynomials.

For Pearson diffusions, the classical orthogonal polynomials are powerful tools to study strong solutions of fractional Kolmogorov equations and characterize a stochastic representation of such solutions via time-changed (through the inverse of a Lévy subordinator) Markov processes [15,25–27]. In the discrete case, fractional (time-changed) processes have been widely considered via different approaches. First of all, a fractional version of the Poisson process has been introduced using Mittag-Leffler distributed inter-jump times instead of exponential ones [2,23,24,29,30] (this approach has been also applied to general counting processes [12]). Such process can be also obtained using a fractional differential-difference equations approach [7,8] and by means of a time-change [31].

With the same approach, some classes of fractional birth-death processes have been introduced and studied [37–39]: in these papers, properties of these processes are deduced from a fractional version of their Kolmogorov forward equation. Let us recall that fractional processes are shown to be interesting in different application contexts, as, for instance, queueing theory ([4,5,11]).

Here, following the approach of [25], we show the existence of strong solutions for the time-fractional counterpart of the Kolmogorov backward and forward equations of immigration-death processes with the aid of Charlier polynomials and link them to a time-changed immigration-death process.

In particular:

- in Section 2 we give some basics on birth-death processes;
- in Section 3 we give some notions on the classical immigration-death process, defining its generator and its forward operator;
- in Section 4 we show the existence of strong solutions of the time-fractional Kolmogorov backward and forward equations under suitable assumptions on the initial data;
- in Section 5 we introduce a fractional immigration-death process and show how the strong solutions of the time-fractional Kolmogorov backward and forward equations can be interpreted by using such process;
- in Section 6, we show the uniqueness of such strong solutions by using the aforementioned stochastic representation and a uniqueness criterion for uniformly bounded solutions [4], under suitable assumptions on the initial data;
- finally, in Section 7 we give the limit distribution of the constructed fractional immigration-death process and we discuss its autocovariance function.

2. Birth-death processes

Let us give some information about general birth-death processes, following the lines of [18,19]. We say that a time-homogeneous continuous time Markov chain $N(t)$ defined on $\mathbb{N}_0 = \{0, 1, 2, \dots\}$ is a birth-death process if and only if, denoting with

$$p(t, x; y) = \mathbb{P}(N(t+s) = x | N(s) = y), \quad x, y = 0, 1, 2, \dots; \quad t, s \geq 0,$$

the transition probability function and $P(t) = (p(t, x; y))_{x, y \geq 0}$ the transition probability matrix, it is solution of the following two differential equations

$$P'(t) = \mathcal{A}P(t), \quad P'(t) = P(t)\mathcal{A}, \tag{1}$$

with initial condition $P(0) = I$ and the infinite matrix $\mathcal{A} = (A(x, y))_{x, y \geq 0}$ is such that:

$$\begin{aligned} A(x, x + 1) &= B(x) \quad x \geq 0, & A(x, x) &= -(B(x) + D(x)) \quad x \geq 0, \\ A(x, x - 1) &= D(x) \quad x \geq 1, & A(x, y) &= 0 \quad |x - y| > 1, \end{aligned}$$

where $B(x) > 0$ for any $x \geq 0$, $D(x) > 0$ for any $x \geq 1$ and $D(0) \geq 0$. Equations (1) are called respectively *backward and forward Kolmogorov equation*. In order to obtain $P(t)$ we need to impose other two properties:

$$P_{i,j}(t) \geq 0, \quad \sum_{j=0}^{+\infty} P_{i,j}(t) \leq 1.$$

In particular it is possible to show that $N(t)$ is a birth-death process if and only if its generator is given by:

$$\begin{aligned} \mathcal{G}f(x) &= (B(x) - D(x))\nabla^+ f(x) + D(x)\Delta f(x) \\ &= (B(x) - D(x))\nabla^- f(x) + B(x)\Delta f(x), \end{aligned}$$

for $x = 0, 1, 2, \dots$ and $f(-1) = 0$, where the difference-type operators ∇^\pm and Δ are defined as

$$\begin{aligned} \nabla^+ f(x) &= f(x + 1) - f(x) \quad \forall x \in \mathbb{N}_0 \\ \nabla^- f(x) &= f(x) - f(x - 1) \quad \forall x \in \mathbb{N}_0 \\ \Delta f(x) &= f(x + 1) - 2f(x) + f(x - 1) \quad \forall x \in \mathbb{N}_0, \end{aligned}$$

and we consider $(P(t))_{t \geq 0}$ as a c_0 -semigroup acting on a suitable Banach sequence space $(\mathbf{b}, \|\cdot\|)$. In particular the generator \mathcal{G} can be represented in terms of the infinite matrix \mathcal{A} , whenever $f \in \text{Dom}(\mathcal{G})$, $x \in \mathbb{N}_0$ and $f(x)$ is considered as a column vector, as $\mathcal{G}f(x) = (\mathcal{A}f)(x)$, where $(\mathcal{A}f)(x)$ is the x -th term of the sequence $\mathcal{A}f$.

The following discrete versions of the Leibnitz rule will be useful

$$\nabla^+(fg)(x) = f(x + 1)\nabla^+g(x) + g(x)\nabla^+f(x) \tag{2}$$

$$\nabla^-(fg)(x) = f(x)\nabla^-g(x) + g(x - 1)\nabla^-f(x) \tag{3}$$

$$\Delta(fg)(x) = f(x + 1)\nabla^+g(x) - f(x - 1)\nabla^-g(x) + g(x)\Delta f(x). \tag{4}$$

The backward Kolmogorov equation becomes, for fixed $x \in \mathbb{N}_0$

$$\begin{cases} p'(t, x; y) = \mathcal{G}p(t, x; y) \\ p(0, x; y) = \delta_{x,y}, \end{cases}$$

where \mathcal{G} works on y and

$$\delta_{x,y} = \begin{cases} 1 & x = y \\ 0 & \text{otherwise,} \end{cases}$$

is Kronecker symbol.

Moreover we can find a *forward* operator

$$\begin{aligned}\mathcal{L} f(x) &= -\nabla^-((B(\cdot) - D(\cdot))f)(x) + \Delta(D(\cdot)f)(x) \\ &= -\nabla^+((B(\cdot) - D(\cdot))f)(x) + \Delta(B(\cdot)f)(x),\end{aligned}$$

so that for fixed $y \in \mathbb{N}_0$ the forward Kolmogorov equation becomes

$$\begin{cases} p'(t, x; y) = \mathcal{L} p(t, x; y) \\ p(0, x; y) = \delta_{x,y}, \end{cases}$$

where \mathcal{L} works on x . As for \mathcal{G} , if we consider $f \in \text{Dom}(\mathcal{L})$, $y \in \mathbb{N}_0$ and $f(y)$ as a row vector, we can represent \mathcal{L} in terms of the infinite matrix \mathcal{A} as $\mathcal{L} f(y) = (f \mathcal{A})(y)$ where $(f \mathcal{A})(y)$ is the y -th term of the sequence $f \mathcal{A}$.

We will focus on the case in which the generator is in the form:

$$\mathcal{G} = p_1(x)\nabla^+ + p_2(x)\Delta,$$

where $p_1(x)$ and $p_2(x)$ are polynomials such that $\deg p_1(x) \leq 1$ and $\deg p_2(x) \leq 2$. Then we can find the *classical orthogonal polynomials of discrete variable* as solution of the equation

$$\mathcal{G} f(x) = -\lambda f(x),$$

for some λ , which is an hypergeometric type difference equation. The values that these polynomials assume on a lattice $\{D_1, D_1 + 1, \dots, D_2\}$ for some D_1, D_2 fully characterize the transition probability and the solutions of the backward and forward Kolmogorov equations. Moreover, these polynomials respect an orthogonality relation in $\ell^2(\mathbf{m})$ for some measure \mathbf{m} called the spectral measure, which is an atomic measure on the lattice. In this case, the spectral measure coincides with the invariant measure of the process $N(t)$ and its mass function $m(x) = \mathbf{m}(\{x\})$ is solution of a discrete analogue of the Pearson equation

$$\nabla^+(p_2(x) \mathbf{m}(x)) = p_1(x) \mathbf{m}(x).$$

For $p_1(x) = a - bx$, we can recognize the following three class of solvable birth-death processes:

- For $p_2(x) = bx$ we have the Immigration-Death process;
- For $p_2(x) = \frac{1}{2}\sigma^2 x$ where $\frac{1}{2}\sigma^2 \neq b$ we have a negative binomial process;
- For $p_2(x) = \frac{1}{2}\sigma^2 x(A - x)$ we have a hypergeometric process.

However, we will focus only on the first case for the choice of the polynomials p_1 and p_2 .

3. Immigration-death processes

Fix $a, b > 0$ the operator

$$\mathcal{G} = (a - bx)\nabla^- + a\Delta;$$

which is a discrete version of the Ornstein-Uhlenbeck generator on \mathbb{N}_0 .

A continuous time Markov chain $N(t)$ defined on \mathbb{N}_0 that admits \mathcal{G} as generator will be called *immigration-death process* (or also $M/M/\infty$ queue: see, for instance, [42]). This process can be generalized to a particular birth-death process with values on a 1-dimensional lattice called *Charlier process* (see [1]), but we will focus on the \mathbb{N}_0 -valued one. For such process, the backward Kolmogorov equations are in the form

$$\frac{du}{dt}(t, x) = \mathcal{G} u(t, x).$$

Moreover, from \mathcal{G} we can recognize the birth and death parameters as

$$B(x) = a, \quad D(x) = bx,$$

and thus the forward operator as

$$\mathcal{L} f(x) = -\nabla^+((a - bz)f(z))(x) + a\Delta f(x),$$

where with $\nabla^+((a - bz)f(z))(x)$ we intend the operator ∇^+ applied to the function $z \mapsto (a - bz)f(z)$ and then evaluated in x .

The operators \mathcal{G} and \mathcal{L} can be represented as infinite matrices. In particular we have $\mathcal{G} = (G(x, y))_{x, y \geq 0}$ where, for $x > 0$

$$\begin{aligned} G(x, x-1) &= bx & G(x, x) &= -(a + bx) & G(x, x+1) &= a \\ G(0, 0) &= -a & G(0, 1) &= a \end{aligned}$$

and $\mathcal{L} = (L(x, y))_{x, y \geq 0}$ where, for $x > 0$

$$\begin{aligned} L(x, x-1) &= a & L(x, x) &= -(a + bx) & L(x, x+1) &= b(x+1) \\ L(0, 0) &= -a & L(0, 1) &= b. \end{aligned}$$

The stationary measure of the process $N(t)$ is the Poisson distribution of parameter $\alpha = \frac{a}{b}$, given by:

$$\mathbf{m}(\{x\}) = e^{-\alpha} \frac{\alpha^x}{x!}, \quad x = 0, 1, 2, \dots$$

Now let us introduce the main Banach sequence spaces we will use through this paper:

- Let us denote with ℓ^∞ the Banach space of bounded functions $f : \mathbb{N}_0 \rightarrow \mathbb{R}$ equipped with the norm

$$\|f\|_{\ell^\infty} = \sup_{x \in \mathbb{N}_0} |f(x)|;$$

- Let us denote with c_0 the subspace of ℓ^∞ of bounded functions $f : \mathbb{N}_0 \rightarrow \mathbb{R}$ such that $\lim_{x \rightarrow +\infty} f(x) = 0$;
- Let us denote with ℓ^1 the Banach space of the functions $f : \mathbb{N}_0 \rightarrow \mathbb{R}$ such that

$$\|f\|_{\ell^1} = \sum_{x=0}^{+\infty} |f(x)| < +\infty.$$

- Let us denote with ℓ^2 the Hilbert space of functions $f : \mathbb{N}_0 \rightarrow \mathbb{R}$ such that

$$\|f\|_{\ell^2}^2 := \sum_{x=0}^{+\infty} f^2(x) < +\infty$$

equipped with the scalar product

$$\langle f, g \rangle_{\ell^2} = \sum_{x=0}^{+\infty} f(x)g(x)$$

- Let us denote with $\ell^2(\mathbf{m})$ the Hilbert space of functions $f : \mathbb{N}_0 \rightarrow \mathbb{R}$ such that

$$\|f\|_{\ell^2(\mathbf{m})}^2 := \sum_{x=0}^{+\infty} \mathbf{m}(\{x\}) f^2(x) < +\infty$$

equipped with the scalar product

$$\langle f, g \rangle_{\ell^2(\mathbf{m})} = \sum_{x=0}^{+\infty} \mathbf{m}(\{x\}) f(x)g(x).$$

Remark 3.1. Let us observe that ℓ^2 is continuously included in $\ell^2(\mathbf{m})$. Consider a function $f \in \ell^2$. Then

$$\sum_{x=0}^{+\infty} m(x) f^2(x) = e^{-\alpha} \sum_{x=0}^{+\infty} \frac{\alpha^x}{x!} f^2(x).$$

Now, let us observe that the sequence $x \mapsto \frac{\alpha^x}{x!}$ converges to 0 as $x \rightarrow +\infty$, hence there exists a constant $C(\alpha)$ such that $\frac{\alpha^x}{x!} \leq C(\alpha)$. Thus

$$e^{-\alpha} \sum_{x=0}^{+\infty} \frac{\alpha^x}{x!} f^2(x) \leq e^{-\alpha} C(\alpha) \|f\|_{\ell^2}^2.$$

Moreover, since ℓ^1 is continuously included in ℓ^2 (see [45]), we have that ℓ^1 is also continuously included in $\ell^2(\mathbf{m})$. Finally, let us observe that, being \mathbf{m} a probability measure, also ℓ^∞ (and then c_0) is continuously embedded in $\ell^2(\mathbf{m})$, with embedding of norm 1.

Concerning the semigroup $(P(t))_{t \geq 0}$, we consider it acting on $\ell^2(\mathbf{m})$ and then $\text{Dom}(\mathcal{G}) = \text{Dom}(\mathcal{L}) = \ell^2(\mathbf{m})$. From the matrix representation of the generator \mathcal{G} and the forward operator \mathcal{L} one can prove the following Lemma.

Lemma 3.2. *The operators $\mathcal{G} : \ell^2(\mathbf{m}) \mapsto \ell^2(\mathbf{m})$ and $\mathcal{L} : \ell^2(\mathbf{m}) \mapsto \ell^2(\mathbf{m})$ are continuous.*

Proof. The proof is a straightforward consequence of Schur's test (see [17]). \square

Moreover, another interesting property that follows from the matrix representation of \mathcal{G} is given by the following Lemma.

Lemma 3.3. *The process $N(t)$ is a Feller process, i.e. the semigroup $(P(t))_{t \geq 0}$ is strongly continuous, contractive and positive on c_0 and $P(t)\mathbf{1} = \mathbf{1}$ for any $t \geq 0$, where $\mathbf{1}(x) = 1$ for any $x \in \mathbb{N}_0$.*

Proof. The proof is a straightforward consequence of [13, Corollary 3.2, Chapter 8]. \square

Let us also observe that the spectrum of \mathcal{G} is given by the sequence $\lambda_n = -bn$, while the eigenfunctions are defined as $x \mapsto C_n(x, \alpha)$ where $\alpha = \frac{a}{b}$ and C_n are the Charlier polynomials (see [34,42]), which are defined by the generating function

$$\sum_{n=0}^{+\infty} C_n(x, \alpha) \frac{t^n}{n!} = e^{-t} \left(1 + \frac{t}{\alpha}\right)^x, \quad t \in \mathbb{R}$$

or via the three terms recurrence relations:

$$-xC_n(x, \alpha) = \alpha C_{n+1}(x, \alpha) - (n + \alpha)C_n(x, \alpha) + nC_{n-1}(x, \alpha), \quad n \geq 0,$$

where $C_0(x, \alpha) \equiv 1$ and $C_{-1}(x, \alpha) \equiv 0$, or

$$C_{n+1}(x, \alpha) = \frac{1}{\alpha} [xC_n(x-1, \alpha) - C_n(x, \alpha)].$$

The first few Charlier polynomials are

$$C_0(x, \alpha) = 1, \quad C_1(x, \alpha) = \frac{x}{\alpha} - 1, \quad C_2(x, \alpha) = \frac{x(x-1)}{\alpha^2} - 2\frac{x}{\alpha} + 1, \quad \dots$$

The orthogonality relation between the polynomials C_n is given by

$$\sum_{x=0}^{+\infty} C_n(x, \alpha) C_m(x, \alpha) \mathbf{m}(\{x\}) = \frac{n!}{\alpha^n} \delta_{n,m},$$

where $\delta_{n,m}$ is the Kronecker delta symbol. Thus, posing $d_n^2 = \frac{n!}{\alpha^n}$, we have that

$$\|C_n(\cdot, \alpha)\|_{\ell^2(\mathbf{m})} = d_n.$$

Let us then define an orthonormal system of polynomials given by

$$Q_n(x) = \frac{C_n(x, \alpha)}{d_n}. \quad (5)$$

Let us also recall that we can exploit the decomposition of a function $g \in \ell^2(\mathbf{m})$ by means of the orthonormal basis $\{Q_n\}_{n \in \mathbb{N}_0}$. Indeed for any $g \in \ell^2(\mathbf{m})$, given the decomposition $g(x) = \sum_{n=0}^{+\infty} g_n Q_n(x)$ where $g_n = \langle g, Q_n \rangle_{\ell^2(\mathbf{m})}$, the sequence $\{g_n\}_{n \in \mathbb{N}_0} \in \ell^2$.

By using such orthonormal system of polynomials, it is well known (see [18,19]) that the transition probability function of the immigration-death process is given by

$$p(t, x_1; x_0) = m(x_1) \sum_{n=0}^{\infty} e^{-bnt} Q_n(x_0) Q_n(x_1),$$

where $m(x) = \mathbf{m}(\{x\})$ and is the fundamental solution of the backward Kolmogorov equation, that is to say that the Cauchy problems

$$\begin{cases} \frac{du}{dt}(t, x) = \mathcal{G} u(t, x) \\ u(0, x) = g(x), \end{cases}$$

and

$$\begin{cases} \frac{dv}{dt}(t, x) = \mathcal{L} v(t, x) \\ v(0, x) = f(x), \end{cases}$$

with $g, f/m \in \ell^2(\mathbf{m})$ admit strong solutions v given by

$$u(t, x) = \sum_{y=0}^{+\infty} p(t, y; x) g(y) = \sum_{n=0}^{+\infty} g_n e^{-bnt} Q_n(x), \quad (6)$$

1 and

$$v(t, x) = \sum_{y=0}^{+\infty} p(t, x; y) f(y) = m(x) \sum_{n=0}^{+\infty} f_n e^{-bnt} Q_n(x),$$

6 where $g(x) = \sum_{n=0}^{+\infty} g_n Q_n(x)$ and $f(x)/m(x) = \sum_{n=0}^{+\infty} f_n Q_n(x)$ and the convergence is uniform. With strong
7 solutions, we intend here that both the functions $t \in [0, +\infty) \mapsto u(t, \cdot) \in \ell^2(\mathbf{m})$ and $t \in [0, +\infty) \mapsto v(t, \cdot) \in$
8 $\ell^2(\mathbf{m})$ belong to $C([0, +\infty); \ell^2(\mathbf{m})) \cap C^1((0, +\infty); \ell^2(\mathbf{m}))$ and the equations hold pointwise.

9 In particular from (6) one easily obtains that

$$\mathbb{E}[N(t)|N(0) = x] = x e^{-bt} + \alpha(1 - e^{-bt}). \quad (7)$$

13 4. Strong solutions in the fractional case

15 Let us introduce the fractional derivative operator (see [28]). Fix $\nu \in (0, 1)$ and consider the Caputo
16 fractional derivative given by

$$\frac{\partial^\nu u}{\partial t^\nu}(t, x) = \frac{1}{\Gamma(1-\nu)} \left[\frac{\partial}{\partial t} \int_0^t (t-\tau)^{-\nu} u(\tau, x) d\tau - \frac{u(0, x)}{t^\nu} \right], \quad (8)$$

21 that, if u is differentiable in t , can be written also as

$$\frac{\partial^\nu u}{\partial t^\nu}(t, x) = \frac{1}{\Gamma(1-\nu)} \int_0^t (t-\tau)^{-\nu} \frac{\partial u}{\partial t}(\tau, x) d\tau,$$

27 and set, for $\nu = 1$, $\frac{\partial^\nu u}{\partial t^\nu} = \frac{\partial u}{\partial t}$. Note that the classes of functions for which the Caputo fractional derivative
28 is well defined are discussed in [32, Section 2.2 and 2.3] (in particular one can use the class of absolutely
29 continuous functions).

30 Denote with

$$\tilde{u}(s, x) = \int_0^{+\infty} e^{-st} u(t, x) dt, \quad s > 0$$

35 the one-sided Laplace transform of u with respect to t . Thus we have that the Laplace transform of $\frac{\partial^\nu u}{\partial t^\nu}$ is
36 given by

$$s^\nu \tilde{u}(s, x) - s^{\nu-1} u(0^+, x).$$

40 We want to find strong solutions for fractional Cauchy problems in the form:

$$\begin{cases} \frac{\partial^\nu u}{\partial t^\nu}(t, x) = \mathcal{G} u(t, x); \\ u(0, x) = g(x), \end{cases} \quad (9)$$

45 for $g \in \ell^2(\mathbf{m})$ with the decomposition $g(x) = \sum_{n=0}^{+\infty} g_n Q_n(x)$. A strong solution of the fractional Cauchy
46 problem (9) will be a function $u : [0, +\infty) \times \mathbb{N}_0 \rightarrow \mathbb{R}$ such that $t \mapsto u(t, \cdot)$ belongs to $C([0, +\infty), \ell^2(\mathbf{m}))$,
47 $\frac{\partial^\nu u}{\partial t^\nu}(t, x)$ exists for any $t > 0$ and $x \in \mathbb{N}_0$, $t \mapsto \frac{\partial^\nu u}{\partial t^\nu}(t, \cdot)$ belongs to $C((0, +\infty); \ell^2(\mathbf{m}))$ and the equation
48 holds pointwise.

The main idea is to find a solution via separation of variables. Indeed, we can suppose that $u(t, x) = T(t)\varphi(x)$ and then observing that, if u is solution of the first equation of (9), then

$$\varphi(x) \frac{d^\nu T}{dt^\nu}(t) = T(t) \mathcal{G} \varphi(x),$$

that leads, if φ and T do not vanish, to the two coupled equations:

$$\begin{cases} \mathcal{G} \varphi(x) = -\lambda \varphi(x), \\ \frac{d^\nu T}{dt^\nu}(t) = -\lambda T(t), \end{cases}$$

which are two eigenvalue problems. In particular we have observed that the first one admits a non zero solution if and only if $\lambda = -bn$ for some $n \in \mathbb{N}_0$ and in that case we can consider $\varphi(x) = Q_n(x)$. Moreover, the second problem admits a solution in the form

$$T(t) = E_\nu(-\lambda t^\nu),$$

where E_ν is the Mittag-Leffler function defined as

$$E_\nu(z) = \sum_{j=0}^{+\infty} \frac{z^j}{\Gamma(1 + \nu j)}, \quad z \in \mathbb{C} \quad (10)$$

(see, for instance, [21]). Thus the idea is to find a solution in the form

$$u(t, x) = \sum_{n=0}^{+\infty} u_n E_\nu(-bnt^\nu) Q_n(x).$$

Moreover, the initial condition suggests that

$$\sum_{n=0}^{+\infty} u_n Q_n(x) = \sum_{n=0}^{+\infty} g_n Q_n(x),$$

so we have $u_n = g_n$ and then we expect the solution to be

$$u(t, x) = \sum_{n=0}^{+\infty} g_n E_\nu(-bnt^\nu) Q_n(x). \quad (11)$$

These heuristic arguments have shown us how should the solution look like, hence we have to prove that such function u is the solution we are searching for.

With the following Lemma, we will first exhibit the fundamental solution of the fractional Cauchy problem in Eq. (9).

Lemma 4.1. *Consider the series*

$$p_\nu(x, t; y) = m(x) \sum_{n=0}^{+\infty} E_\nu(-bnt^\nu) Q_n(x) Q_n(y), \quad (12)$$

where Q_n and E_ν are the functions defined in Equations (5) and (10) and $m(x) = \mathbf{m}(\{x\})$. Then such series converges for fixed $t > 0$ and $x, y \in \mathbb{N}_0$.

Proof. To show the convergence of $p_\nu(x, t; y)$, we need the following *self-duality* property of the Charlier polynomials (see [34, Equation 2.7.10a]):

$$C_n(x, \alpha) = C_x(n, \alpha), \quad \forall n, x \in \mathbb{N}_0. \quad (13)$$

From this relation we have

$$\begin{aligned} p_\nu(x, t; y) &= m(x) \sum_{n=0}^{+\infty} E_\nu(-bnt^\nu) Q_n(x) Q_n(y) \\ &= m(x) \sum_{n=0}^{+\infty} \frac{1}{d_n^2} E_\nu(-bnt^\nu) C_n(x, \alpha) C_n(y, \alpha) \\ &= m(x) \sum_{n=0}^{+\infty} \frac{1}{d_n^2} E_\nu(-bnt^\nu) C_x(n, \alpha) C_y(n, \alpha), \end{aligned}$$

hence we need to show the convergence of the series

$$\sum_{n=0}^{+\infty} \frac{1}{d_n^2} E_\nu(-bnt^\nu) C_x(n, \alpha) C_y(n, \alpha).$$

Now, let us observe that equation (13) made us fix the degrees of the polynomials involved in the series. Thus, let us denote with z_x and z_y the last real zeroes of $C_x(\cdot, \alpha)$ and $C_y(\cdot, \alpha)$ and then let us consider $n_0 > \max\{z_x, z_y\}$.

We will equivalently prove that the series

$$\sum_{n=n_0}^{+\infty} \frac{1}{d_n^2} E_\nu(-bnt^\nu) C_x(n, \alpha) C_y(n, \alpha) \quad (14)$$

converges. To do this, we need to recall another property of the Charlier polynomials. In particular it is known (see [34, Table 2.3]) that the director coefficient of $C_n(\cdot, \alpha)$ is given by

$$c_n = \frac{1}{(-\alpha)^n}.$$

In particular, recalling that $\alpha = \frac{a}{b}$, $\alpha > 0$ since $a, b > 0$ and then $c_n > 0$ if n is even and $c_n < 0$ if n is odd. By using this observation, we can distinguish two cases:

- i If $x + y$ is even, then, since $c_x c_y > 0$, for any $n \geq n_0$ $C_x(n, \alpha) C_y(n, \alpha) > 0$ and then the series (14) admits only positive summands. Recalling that $E_\nu(-bnt^\nu) \leq 1$ we obtain

$$\sum_{n=n_0}^{+\infty} \frac{1}{d_n^2} E_\nu(-bnt^\nu) C_x(n, \alpha) C_y(n, \alpha) \leq \sum_{n=n_0}^{+\infty} \frac{1}{d_n^2} C_x(n, \alpha) C_y(n, \alpha)$$

where the RHS series converges since

$$\sum_{n=0}^{+\infty} \frac{1}{d_n^2} C_x(n, \alpha) C_y(n, \alpha) = e^\alpha \sum_{n=0}^{+\infty} \frac{\alpha^n}{n!} e^{-\alpha} C_x(n, \alpha) C_y(n, \alpha) = e^\alpha d_x^2 \delta_{x,y}. \quad (15)$$

ii If $x + y$ is odd, then, since $c_x c_y < 0$, for any $n \geq n_0$ $C_x(n, \alpha)C_y(n, \alpha) < 0$ and then the series (14) admits only negative summands. As before, we obtain

$$\sum_{n=n_0}^{+\infty} \frac{1}{d_n^2} E_\nu(-bnt^\nu) C_x(n, \alpha) C_y(n, \alpha) \geq \sum_{n=n_0}^{+\infty} \frac{1}{d_n^2} C_x(n, \alpha) C_y(n, \alpha)$$

where the RHS series converges for equation (15). \square

With Lemma 4.1, we have exploited the fundamental solution of the equation in (9). Now we have to show that a function in the form (11) is a solution for such fractional Cauchy problem. To do this, let us first show a technical lemma.

Lemma 4.2. For any $t_0 > 0$, there exists a constant $K(t_0, \nu)$ such that

$$bnE_\nu(-bnt^\nu) \leq K(t_0, \nu), \quad t \in [t_0, +\infty).$$

Proof. Let us use the uniform estimate for the Mittag-Leffler function given in [44, Theorem 4]:

$$bnE_\nu(-bnt^\nu) \leq \frac{bn}{1 + \frac{bnt^\nu}{\Gamma(1+\nu)}}.$$

Consider the function

$$f(x) = \frac{x}{1 + Cx}, \quad C = \frac{t^\nu}{\Gamma(1 + \nu)}.$$

Thus we have

$$f'(x) = \frac{1}{(1 + Cx)^2} > 0$$

hence the function f is strictly increasing. So we have

$$f(x) \leq \lim_{x \rightarrow +\infty} f(x) = \frac{1}{C} = \frac{\Gamma(1 + \nu)}{t^\nu},$$

and then

$$bnE_\nu(-bnt^\nu) \leq \frac{bn}{1 + \frac{bnt^\nu}{\Gamma(1+\nu)}} \leq \frac{\Gamma(1 + \nu)}{t^\nu} \leq \frac{\Gamma(1 + \nu)}{t_0^\nu} =: K(t_0, \nu). \quad \square$$

Now let us exhibit a strong solution for our fractional Cauchy problem.

Theorem 4.3. Let $g \in \ell^2(\mathbf{m})$ with decomposition $g(x) = \sum_{n=0}^{+\infty} g_n Q_n(x)$. Then the fractional difference-differential Cauchy problem

$$\begin{cases} \frac{\partial^\nu u}{\partial t^\nu}(t, x) = \mathcal{G} u(t, x) \\ u(0, x) = g(x), \end{cases} \quad (16)$$

admits a strong solution u in the form

$$u(t, x) = \sum_{y=0}^{+\infty} p_\nu(t, y; x) g(y) = \sum_{n=0}^{\infty} E_\nu(-bnt^\nu) Q_n(x) g_n. \quad (17)$$

Proof. First let us observe that obviously if u is in the form (17), then $u(0, x) = g(x)$.

Now, let us notice that

$$\mathcal{G} E_\nu(-bnt^\nu)Q_n(x)g_n = E_\nu(-bnt^\nu)g_n \mathcal{G} Q_n(x) = -bnE_\nu(-bnt^\nu)g_nQ_n(x) = g_nQ_n(x)\frac{d^\nu E_\nu(-bnt^\nu)}{dt^\nu}.$$

Hence we need to show that the series in (17) is convergent at least uniformly in t and that we can change the series with the operators.

Starting from the convergence of the series, by using Cauchy-Schwartz inequality we have

$$\begin{aligned} \sum_{n=0}^{+\infty} |E_\nu(-bnt^\nu)Q_n(x)g_n| &\leq \sum_{n=0}^{+\infty} |Q_n(x)g_n| \\ &\leq \left(\sum_{n=0}^{+\infty} \frac{\alpha^n}{n!} C_n^2(x, \alpha)\right)^{\frac{1}{2}} \left(\sum_{n=0}^{+\infty} g_n^2\right)^{\frac{1}{2}} \\ &= \|g\|_{\ell^2(\mathbf{m})} \left(\sum_{n=0}^{+\infty} \frac{\alpha^n}{n!} C_x^2(n, \alpha)\right)^{\frac{1}{2}} \\ &= \|g\|_{\ell^2(\mathbf{m})} e^{\frac{\alpha}{2} dx}, \end{aligned} \tag{18}$$

hence the series in (17) totally converges.

To show that the series converges in $\ell^2(\mathbf{m})$ for any $t > 0$, let us recall that $g_n \in \ell^2$ by definition. Let us consider $N \in \mathbb{N}$ and define

$$u_N(t, x) = \sum_{n=0}^N E_\nu(-bnt^\nu)Q_n(x)g_n.$$

Then we have, for any $N, M \in \mathbb{N}$ with $N < M$, by also using Lemma 4.2

$$\begin{aligned} \|u_N(t, \cdot) - u_M(t, \cdot)\|_{\ell^2(\mathbf{m})}^2 &= \left\| \sum_{n=N}^M E_\nu(-bnt^\nu)Q_n(\cdot)g_n \right\|_{\ell^2(\mathbf{m})}^2 \\ &= \sum_{n=N}^M E_\nu^2(-bnt^\nu)g_n^2 \\ &\leq \sum_{n=N}^M g_n^2, \end{aligned}$$

that implies the convergence in $\ell^2(\mathbf{m})$ of the sequence $u_N(t, \cdot)$ by Cauchy's criterion.

Now we need to show that one can exchange the operators with the series. To do that, let us first observe that

$$\int_0^t (t - \tau)^{-\nu} u(\tau) d\tau = \int_0^t \frac{u(\tau)}{\nu - 1} d(t - \tau)^{1-\nu},$$

and since $(t - \tau)^{1-\nu}$ is strictly decreasing in $[0, t]$ we can use [40, Theorem 7.16] with the total convergence of the series (17) to obtain

$$\int_0^t (t-\tau)^{-\nu} \sum_{n=0}^{+\infty} E_{\nu}(-bn\tau^{\nu}) Q_n(x) g_n d\tau = \sum_{n=0}^{+\infty} \int_0^t (t-\tau)^{-\nu} E_{\nu}(-bn\tau^{\nu}) Q_n(x) g_n d\tau.$$

Now we want to use the following relation:

$$\begin{aligned} \frac{d}{dt} \int_0^t (t-\tau)^{-\nu} \sum_{n=0}^{+\infty} E_{\nu}(-bn\tau^{\nu}) Q_n(x) g_n d\tau &= \frac{d}{dt} \sum_{n=0}^{+\infty} \int_0^t (t-\tau)^{-\nu} E_{\nu}(-bn\tau^{\nu}) Q_n(x) g_n d\tau \\ &= \sum_{n=0}^{+\infty} \frac{d}{dt} \int_0^t (t-\tau)^{-\nu} E_{\nu}(-bn\tau^{\nu}) Q_n(x) g_n d\tau, \end{aligned}$$

but to do this, by using [40, Theorem 7.17], we need to show the uniform convergence of

$$\sum_{n=0}^{+\infty} \frac{d}{dt} \int_0^t (t-\tau)^{-\nu} E_{\nu}(-bn\tau^{\nu}) Q_n(x) g_n d\tau,$$

in any compact interval included in $(0, +\infty)$. Hence, by definition of Caputo fractional derivative, as given in (8), we actually need to show the uniform convergence of

$$\sum_{n=0}^{+\infty} \frac{d^{\nu}}{dt^{\nu}} E_{\nu}(-bnt^{\nu}) Q_n(x) g_n \quad (19)$$

in any interval of the form $[t_0, +\infty)$. To do this, let us recall that $\frac{d^{\nu}}{dt^{\nu}} E_{\nu}(-bnt^{\nu}) = -bnE_{\nu}(-bnt^{\nu})$ and thus we need to show the uniform convergence of

$$\sum_{n=0}^{+\infty} -bnE_{\nu}(-bnt^{\nu}) Q_n(x) g_n.$$

Thus, fix $t_0 > 0$ and observe that

$$\begin{aligned} \sum_{n=0}^{+\infty} |bnE_{\nu}(-bnt^{\nu}) Q_n(x) g_n| &\leq K(t_0, \nu) \sum_{n=0}^{+\infty} |Q_n(x) g_n| \\ &\leq K(t_0, \nu) \|g\|_{\ell^2(\mathbf{m})} e^{\frac{\sigma}{2} dx}, \quad t \in [t_0, +\infty), \end{aligned}$$

where the first inequality follows from Lemma 4.2 and the second inequality from Cauchy-Schwartz inequality as done before in (18). Hence we have shown the total convergence of (19) in any interval of the form $[t_0, +\infty)$.

We have already shown that $\sum_{n=0}^{+\infty} E_{\nu}(-bnt^{\nu}) Q_n(x) g_n$ totally converges with respect to t : in the same way we have that also $\sum_{n=0}^{+\infty} E_{\nu}(-bnt^{\nu}) Q_n(x-1) g_n$ and $\sum_{n=0}^{+\infty} E_{\nu}(-bnt^{\nu}) Q_n(x+1) g_n$ totally converge with respect to t .

Now, observe that

$$\begin{aligned}
 \nabla^- \sum_{n=0}^{+\infty} E_\nu(-bnt^\nu) Q_n(x) g_n &= \sum_{n=0}^{+\infty} E_\nu(-bnt^\nu) Q_n(x) g_n - \sum_{n=0}^{+\infty} E_\nu(-bnt^\nu) Q_n(x-1) g_n \\
 &= \lim_{N \rightarrow +\infty} \sum_{n=0}^N E_\nu(-bnt^\nu) Q_n(x) g_n - \lim_{N \rightarrow +\infty} \sum_{n=0}^N E_\nu(-bnt^\nu) Q_n(x-1) g_n \\
 &= \lim_{N \rightarrow +\infty} \left(\sum_{n=0}^N E_\nu(-bnt^\nu) Q_n(x) g_n - \sum_{n=0}^N E_\nu(-bnt^\nu) Q_n(x-1) g_n \right) \\
 &= \lim_{N \rightarrow +\infty} \sum_{n=0}^N E_\nu(-bnt^\nu) \nabla^- Q_n(x) g_n \\
 &= \sum_{n=0}^{+\infty} E_\nu(-bnt^\nu) \nabla^- Q_n(x) g_n,
 \end{aligned} \tag{20}$$

and in the same way one can show that

$$\Delta \sum_{n=0}^{+\infty} E_\nu(-bnt^\nu) Q_n(x) g_n = \sum_{n=0}^{+\infty} E_\nu(-bnt^\nu) \Delta Q_n(x) g_n.$$

By using these last two relations, it is easy to show that

$$\mathcal{G} \sum_{n=0}^{+\infty} E_\nu(-bnt^\nu) Q_n(x) g_n = \sum_{n=0}^{+\infty} E_\nu(-bnt^\nu) \mathcal{G} Q_n(x) g_n.$$

Finally we have that

$$\begin{aligned}
 \frac{d^\nu}{dt^\nu} \sum_{n=0}^{+\infty} E_\nu(-bnt^\nu) Q_n(x) g_n &= \sum_{n=0}^{+\infty} \frac{d^\nu}{dt^\nu} E_\nu(-bnt^\nu) Q_n(x) g_n \\
 &= \sum_{n=0}^{+\infty} \mathcal{G} E_\nu(-bnt^\nu) Q_n(x) g_n \\
 &= \mathcal{G} \sum_{n=0}^{+\infty} E_\nu(-bnt^\nu) Q_n(x) g_n,
 \end{aligned}$$

and we have shown the pointwise relations.

Moreover, $\frac{\partial^\nu u}{\partial t^\nu}(t, x)$ belongs to $\ell^2(\mathbf{m})$ by definition of \mathcal{G} . By continuity of the operator \mathcal{G} , if we show that $(t \mapsto u(t, \cdot)) \in C([0, +\infty); \ell^2(\mathbf{m}))$, then also $(t \mapsto \frac{\partial^\nu u}{\partial t^\nu}(t, \cdot)) \in C((0, +\infty); \ell^2(\mathbf{m}))$. Let us show the continuity of $t \mapsto u(t, \cdot)$ at 0, since for any point $t \in (0, +\infty)$ the proof is analogous. To do this, let us observe that, since all the series involved are uniformly convergent

$$\begin{aligned}
 \|u(t, \cdot) - g(\cdot)\|_{\ell^2(\mathbf{m})} &= \left\| \sum_{n=1}^{+\infty} (E_\nu(-bnt^\nu) - 1) Q_n(\cdot) g_n \right\|_{\ell^2(\mathbf{m})} \\
 &= \sum_{n=1}^{+\infty} (1 - E_\nu(-bnt^\nu))^2 g_n^2.
 \end{aligned}$$

Now let us fix $\varepsilon > 0$ and consider $n_\varepsilon \geq 1$ such that $\sum_{n=n_\varepsilon}^{+\infty} g_n^2 \leq \varepsilon$. Then we have

$$\|u(t, \cdot) - g(\cdot)\|_{\ell^2(\mathbf{m})} \leq (1 - E_\nu(-bn_\varepsilon t^\nu))^2 \|g\|_{\ell^2(\mathbf{m})} + \varepsilon,$$

thus, sending $t \rightarrow 0^+$ and $\varepsilon \rightarrow 0^+$ we conclude the proof. \square

The same strategy can be used to exhibit a strong solution to the fractional forward Kolmogorov equation.

Theorem 4.4. *Let f be a function such that $f/m \in \ell^2(\mathbf{m})$ with decomposition $f(x)/m(x) = \sum_{n=0}^{+\infty} f_n Q_n(x)$. Then the fractional difference-differential Cauchy problem*

$$\begin{cases} \frac{\partial^\nu u}{\partial t^\nu}(t, x) = \mathcal{L} u(t, x) \\ u(0, x) = f(x), \end{cases} \tag{21}$$

admits a strong solution $u = u(t, x)$ given by

$$u(t, x) = \sum_{y=0}^{+\infty} p_\nu(t, x; y) f(y) = m(x) \sum_{n=0}^{+\infty} E_\nu(-bnt^\nu) Q_n(x) f_n.$$

Proof. Since $\{f_n\}_{n \in \mathbb{N}} \in \ell^2$, then, from the previous theorem, we already know that we can exchange operators and series. We only need to prove that the single summand of the series is a solution of the equation and that $u(0, x) = f(x)$.

Let us first notice that

$$u(0, x) = m(x) \sum_{n=0}^{+\infty} Q_n(x) f_n = m(x) \frac{f(x)}{m(x)} = f(x),$$

thus the function u satisfies the given initial condition.

To show that the single summand is solution of the equation, let us write \mathcal{L} as

$$\mathcal{L} h(x) = -\nabla^-((a - bz)h(z))(x) + \Delta(bzh(z))(x),$$

for a generic function h .

Moreover, let us observe that

$$\begin{aligned} \mathcal{G} h(x) &= (a - bx)\nabla^- h(x) + a\Delta h(x) \\ &= ah(x) - ah(x - 1) - bx\nabla^- h(x) + ah(x + 1) - 2ah(x) + ah(x - 1) \\ &= ah(x + 1) - ah(x) - bx\nabla^- h(x) \\ &= a\nabla^+ h(x) - bx\nabla^- h(x). \end{aligned}$$

Let us also recall that m solves a discrete Pearson equation:

$$\nabla^+(b \cdot m(z))(x) = (a - bx)m(x). \tag{22}$$

Now, let us observe that

$$\mathcal{L}(m(z)Q_n(z)E_\nu(-bnt^\nu)f_n)(x) = f_n E_\nu(-bnt^\nu) \mathcal{L}(m(z)Q_n(z))(x),$$

hence we will only study $\mathcal{L}(m(\cdot)Q_n(\cdot))$. In particular we have

$$\mathcal{L}(m(z)Q_n(z))(x) = -\nabla^-((a - bz)m(z)Q_n(z))(x) + \Delta(bzm(z)Q_n(z))(x),$$

1 hence, by using the Discrete Leibnitz Rule ((3) and (4)), we obtain

$$\begin{aligned}
 \mathcal{L}(m(z)Q_n(z))(x) &= -[Q_n(x)\nabla^-((a-bz)m(z))(x) + (a-b(x-1))m(x-1)\nabla^-Q_n(x)] \\
 &\quad + Q_n(x)\Delta(bzm(z))(x) + b(x+1)m(x+1)\nabla^+Q(x) + \\
 &\quad - b(x-1)m(x-1)\nabla^-Q_n(x) \\
 &= Q_n(x)[- \nabla^-((a-bz)m(z))(x) + \Delta(bzm(z))(x)] + \\
 &\quad - am(x-1)\nabla^-Q_n(x) + b(x+1)m(x+1)\nabla^+Q_n(x).
 \end{aligned}$$

10 First let us observe that $\Delta = \nabla^- \nabla^+$, then

$$\begin{aligned}
 -\nabla^-((a-bz)m(z)) + \Delta(bzm(z)) &= -\nabla^-((a-bz)m(z)) + \nabla^- \nabla^+(bzm(z)) \\
 &= \nabla^- (\nabla^+(bzm(z)))(x) - (a-bx)m(x) = 0,
 \end{aligned}$$

15 since m satisfies equation (22). Moreover

$$am(x-1) = a \frac{\alpha^{x-1}}{(x-1)!} e^{-\alpha} = \frac{a}{\alpha} xm(x) = bxm(x),$$

19 while

$$b(x+1)m(x+1) = b(x+1) \frac{\alpha^{x+1}}{(x+1)!} e^{-\alpha} = b\alpha m(x) = am(x),$$

24 thus

$$\begin{aligned}
 \mathcal{L}(m(z)Q_n(z))(x) &= -bxm(x)\nabla^-Q_n(x) + am(x)\nabla^+Q_n(x) \\
 &= m(x)[\nabla^+Q_n(x) - bx\nabla^-Q_n(x)] \\
 &= m(x)\mathcal{G}Q_n(x).
 \end{aligned}$$

30 Finally, we obtain:

$$\begin{aligned}
 \mathcal{L}(m(z)Q_n(z)E_\nu(-bnt^\nu)f_n) &= f_n E_\nu(-bnt^\nu) \mathcal{L}(m(z)Q_n(z)) \\
 &= f_n E_\nu(-bnt^\nu)m(x)\mathcal{G}Q_n(x) \\
 &= -bnf_n E_\nu(-bnt^\nu)m(x)Q_n(x) \\
 &= f_n m(x)Q_n(x) \frac{d^\nu E_\nu(-bnt^\nu)}{dt^\nu} \\
 &= \frac{d^\nu}{dt^\nu}(f_n m(x)Q_n(x)E_\nu(-bnt^\nu)). \quad \square
 \end{aligned}$$

40 **Remark 4.5.** It is easy to see that $p_\nu(t, x; y)$ is strong solution of the fractional backward equation

$$\begin{cases} \frac{d^\nu p_\nu}{dt^\nu}(t, x; y) = \mathcal{G} p_\nu(t, x; y) \\ p_\nu(0, x; y) = \delta_{x,y}, \end{cases} \tag{23}$$

45 where \mathcal{G} operates on y , and is also strong solution of the fractional forward equation

$$\begin{cases} \frac{d^\nu p_\nu}{dt^\nu}(t, x; y) = \mathcal{L} p_\nu(t, x; y) \\ p_\nu(0, x; y) = \delta_{x,y}, \end{cases} \tag{24}$$

where \mathcal{L} operates on x . In particular, as shown by Theorems 4.3 and 4.4, it is the fundamental solution of such equations.

5. Stochastic representation of the solutions

Now we want to exhibit a process whose “transition probability” is the fundamental solution $p_\nu(t, x; y)$ we have described previously.

To do this, let us consider a classical immigration-death process $N_1(t)$ (as defined before). Let us also consider a ν -stable subordinator $\sigma_\nu(t)$ with Laplace transform

$$\mathbb{E}[e^{-s\sigma_\nu(t)}] = e^{-ts^\nu}, \quad s > 0, \quad \nu \in (0, 1)$$

and its inverse process (or first passage time process) $L_\nu(t)$ defined as

$$L_\nu(t) := \inf\{s > 0 : \sigma_\nu(s) > t\}.$$

The latter admits density (see [2,33])

$$\mathbb{P}(L_\nu(t) \in dy) = f_\nu(y, t)dy = \frac{t}{\nu} \frac{1}{y^{1+\frac{1}{\nu}}} g_\nu\left(\frac{t}{y^{\frac{1}{\nu}}}\right) dy, \quad y \geq 0, \quad t > 0,$$

where $g_\nu(x)$ is the density of $\sigma_\nu(1)$ given by

$$g_\nu(x) = \frac{1}{\pi} \sum_{k=1}^{+\infty} (-1)^{k+1} \frac{\Gamma(\nu k + 1)}{k!} \frac{1}{x^{\nu k + 1}}, \quad x \geq 0.$$

Alternatives for $f_\nu(y, t)$ are given in [20,22].

Thus, let us define the *fractional immigration-death process* as $N_\nu(t) := N_1(L_\nu(t))$. This is a semi-Markov process as defined in [16]. However, we say that such process admits a transition probability mass $p_\nu(t, x; y)$ if for any $B \subseteq \mathbb{N}_0$:

$$\mathbb{P}(N_\nu(t) \in B \mid N_\nu(0) = y) = \sum_{x \in B} p_\nu(t, x; y).$$

Hence, we can use such process to characterize the fundamental solution we found in the previous section.

Theorem 5.1. *The process $N_\nu(t)$ admits a transition probability mass $p_\nu(t, x; y)$ in the form (12).*

Proof. Let us first recall that (see, for instance, [33]) the process $L_\nu(t)$ admits a density $f_t(\tau) = \mathbb{P}(L_\nu(t) \in d\tau)$. Moreover, let us recall (see [9]) that

$$\int_0^{+\infty} e^{-s\tau} f_t(\tau) d\tau = E_\nu(-st^\nu), \quad s > 0.$$

Now, observe that for any $B \subseteq \mathbb{N}_0$, since $N_1(t)$ admits a transition probability mass, we have

$$\begin{aligned} \mathbb{P}(N_\nu(t) \in B | N_\nu(0) = y) &= \int_0^{+\infty} \mathbb{P}(N_1(\tau) \in B | N_1(0) = y) f_t(\tau) d\tau \\ &= \int_0^{+\infty} \sum_{x \in B} p_1(\tau, x; y) f_t(\tau) d\tau. \end{aligned}$$

Now, if B is a finite set, we have

$$\int_0^{+\infty} \sum_{x \in B} p_1(\tau, x; y) f_t(\tau) d\tau = \sum_{x \in B} \int_0^{+\infty} p_1(\tau, x; y) f_t(\tau) d\tau.$$

If B is infinite, let us consider the sets $I_m := \{x \in \mathbb{N}_0 : x \leq m\}$ and $B_m := B \cap I_m$. Thus, B_m is finite and then

$$\int_0^{+\infty} \sum_{x \in B_m} p_1(\tau, x; y) f_t(\tau) d\tau = \sum_{x \in B_m} \int_0^{+\infty} p_1(\tau, x; y) f_t(\tau) d\tau.$$

Since $p_1(\tau, x; y) f_t(\tau)$ is non-negative, we can use the monotone convergence theorem to obtain, taking the limit as $m \rightarrow +\infty$

$$\int_0^{+\infty} \sum_{x \in B} p_1(\tau, x; y) f_t(\tau) d\tau = \sum_{x \in B} \int_0^{+\infty} p_1(\tau, x; y) f_t(\tau) d\tau.$$

Now let us only consider

$$\int_0^{+\infty} p_1(\tau, x; y) f_t(\tau) d\tau,$$

and recall that (see [18,19])

$$p_1(\tau, x; y) = m(x) \sum_{n=0}^{+\infty} e^{-bn\tau} Q_n(x) Q_n(y).$$

Hence we have

$$\int_0^{+\infty} p_1(\tau, x; y) f_t(\tau) d\tau = m(x) \int_0^{+\infty} \sum_{n=0}^{+\infty} e^{-bn\tau} Q_n(x) Q_n(y) f_t(\tau) d\tau.$$

Now we have to show that we can exchange integral and series. To do this, let us first observe that

$$\begin{aligned} \sum_{n=0}^{+\infty} e^{-bn\tau} Q_n(x) Q_n(y) f_t(\tau) &= \sum_{n=0}^{+\infty} \frac{\alpha^n}{n!} e^{-bn\tau} C_n(x, \alpha) C_n(y, \alpha) f_t(\tau) \\ &= \sum_{n=0}^{+\infty} \frac{\alpha^n}{n!} e^{-bn\tau} C_x(n, \alpha) C_y(n, \alpha) f_t(\tau). \end{aligned}$$

Let us consider z_x and z_y the last real zeros of $C_x(n, \alpha)$ and $C_y(n, \alpha)$ and consider a $n_0 \in \mathbb{N}$ such that $n_0 > \max\{z_x, z_y\}$. Thus we have

$$\begin{aligned} \sum_{n=0}^{+\infty} \frac{\alpha^n}{n!} e^{-bn\tau} C_x(n, \alpha) C_y(n, \alpha) f_t(\tau) &= \sum_{n=0}^{n_0} \frac{\alpha^n}{n!} e^{-bn\tau} C_x(n, \alpha) C_y(n, \alpha) f_t(\tau) \\ &+ \sum_{n=n_0+1}^{+\infty} \frac{\alpha^n}{n!} e^{-bn\tau} C_x(n, \alpha) C_y(n, \alpha) f_t(\tau), \end{aligned}$$

and

$$\begin{aligned} \int_0^{+\infty} \sum_{n=0}^{+\infty} \frac{\alpha^n}{n!} e^{-bn\tau} C_x(n, \alpha) C_y(n, \alpha) f_t(\tau) d\tau &= \int_0^{+\infty} \sum_{n=0}^{n_0} \frac{\alpha^n}{n!} e^{-bn\tau} C_x(n, \alpha) C_y(n, \alpha) f_t(\tau) d\tau \\ &+ \int_0^{+\infty} \sum_{n=n_0+1}^{+\infty} \frac{\alpha^n}{n!} e^{-bn\tau} C_x(n, \alpha) C_y(n, \alpha) f_t(\tau) d\tau \\ &= \sum_{n=0}^{n_0} \frac{\alpha^n}{n!} C_x(n, \alpha) C_y(n, \alpha) \int_0^{+\infty} e^{-bn\tau} f_t(\tau) d\tau \\ &+ \int_0^{+\infty} \sum_{n=n_0+1}^{+\infty} \frac{\alpha^n}{n!} e^{-bn\tau} C_x(n, \alpha) C_y(n, \alpha) f_t(\tau) d\tau. \end{aligned}$$

Now, fix $\tau_0 > 0$ and observe that for $\tau > \tau_0$ and $n \geq n_0 + 1$ the function

$$(\tau, n) \mapsto \frac{\alpha^n}{n!} e^{-bn\tau} C_x(n, \alpha) C_y(n, \alpha) f_t(\tau)$$

does not change sign, by Fubini's theorem (see [41, Theorem 8.8]) we have that

$$\int_{\tau_0}^{+\infty} \sum_{n=n_0+1}^{+\infty} \frac{\alpha^n}{n!} e^{-bn\tau} C_x(n, \alpha) C_y(n, \alpha) f_t(\tau) d\tau = \sum_{n=n_0+1}^{+\infty} \frac{\alpha^n}{n!} C_x(n, \alpha) C_y(n, \alpha) \int_{\tau_0}^{+\infty} e^{-bn\tau} f_t(\tau) d\tau.$$

Now we have to pass to the limit as $\tau_0 \rightarrow 0$. To do this, let us observe that

$$\int_{\tau_0}^{+\infty} e^{-bn\tau} f_t(\tau) d\tau \leq \int_0^{+\infty} e^{-bn\tau} f_t(\tau) d\tau = E_\nu(-bnt^\nu),$$

and let us distinguish two cases.

i) If $x + y$ is even,

$$\frac{\alpha^n}{n!} C_x(n, \alpha) C_y(n, \alpha) \int_{\tau_0}^{+\infty} e^{-bn\tau} f_t(\tau) d\tau \geq 0,$$

and in particular we have

$$\begin{aligned} \frac{\alpha^n}{n!} C_x(n, \alpha) C_y(n, \alpha) \int_{\tau_0}^{+\infty} e^{-bn\tau} f_t(\tau) d\tau &\leq \frac{\alpha^n}{n!} C_x(n, \alpha) C_y(n, \alpha) E_\nu(-bnt^\nu) \\ &\leq \frac{\alpha^n}{n!} C_x(n, \alpha) C_y(n, \alpha), \end{aligned}$$

where

$$\sum_{n=n_0+1}^{+\infty} \frac{\alpha^n}{n!} C_x(n, \alpha) C_y(n, \alpha) < +\infty,$$

as we observed before. Then we can use dominated convergence theorem to take the limit as $\tau_0 \rightarrow 0$ and obtain

$$\int_0^{+\infty} \sum_{n=n_0+1}^{+\infty} C_x(n, \alpha) C_y(n, \alpha) e^{-bn\tau} f_t(\tau) d\tau = \sum_{n=n_0+1}^{+\infty} C_x(n, \alpha) C_y(n, \alpha) \int_0^{+\infty} e^{-bn\tau} f_t(\tau) d\tau.$$

ii) If $x + y$ is odd, then

$$\frac{\alpha^n}{n!} C_x(n, \alpha) C_y(n, \alpha) \int_{\tau_0}^{+\infty} e^{-bn\tau} f_t(\tau) d\tau \leq 0,$$

and in particular we have

$$\begin{aligned} -\frac{\alpha^n}{n!} C_x(n, \alpha) C_y(n, \alpha) \int_{\tau_0}^{+\infty} e^{-bn\tau} f_t(\tau) d\tau &\leq -\frac{\alpha^n}{n!} C_x(n, \alpha) C_y(n, \alpha) E_\nu(-bnt^\nu) \\ &\leq -\frac{\alpha^n}{n!} C_x(n, \alpha) C_y(n, \alpha), \end{aligned}$$

where

$$-\sum_{n=n_0+1}^{+\infty} \frac{\alpha^n}{n!} C_x(n, \alpha) C_y(n, \alpha) < +\infty,$$

as we observed before. Then we can use dominated convergence theorem to take the limit as $\tau_0 \rightarrow 0$ and obtain

$$\int_0^{+\infty} \sum_{n=n_0+1}^{+\infty} C_x(n, \alpha) C_y(n, \alpha) e^{-bn\tau} f_t(\tau) d\tau = \sum_{n=n_0+1}^{+\infty} C_x(n, \alpha) C_y(n, \alpha) \int_0^{+\infty} e^{-bn\tau} f_t(\tau) d\tau.$$

Hence in general we have for any $x, y \in \mathbb{N}_0$

$$\int_0^{+\infty} \sum_{n=n_0+1}^{+\infty} C_x(n, \alpha) C_y(n, \alpha) e^{-bn\tau} f_t(\tau) d\tau = \sum_{n=n_0+1}^{+\infty} C_x(n, \alpha) C_y(n, \alpha) \int_0^{+\infty} e^{-bn\tau} f_t(\tau) d\tau$$

and then

$$\begin{aligned}
\int_0^{+\infty} p_1(\tau, x; y) f_t(\tau) d\tau &= \int_0^{+\infty} m(x) \sum_{n=0}^{+\infty} \frac{\alpha^n}{n!} e^{-bn\tau} C_x(n, \alpha) C_y(n, \alpha) f_t(\tau) d\tau \\
&= m(x) \sum_{n=0}^{n_0} \frac{\alpha^n}{n!} C_x(n, \alpha) C_y(n, \alpha) \int_0^{+\infty} e^{-bn\tau} f_t(\tau) d\tau \\
&\quad + m(x) \int_0^{+\infty} \sum_{n=n_0+1}^{+\infty} \frac{\alpha^n}{n!} e^{-bn\tau} C_x(n, \alpha) C_y(n, \alpha) f_t(\tau) d\tau \\
&= m(x) \sum_{n=0}^{n_0} \frac{\alpha^n}{n!} C_x(n, \alpha) C_y(n, \alpha) \int_0^{+\infty} e^{-bn\tau} f_t(\tau) d\tau \\
&\quad + m(x) \sum_{n=n_0+1}^{+\infty} \frac{\alpha^n}{n!} C_x(n, \alpha) C_y(n, \alpha) \int_0^{+\infty} e^{-bn\tau} f_t(\tau) d\tau \\
&= m(x) \sum_{n=0}^{+\infty} \frac{\alpha^n}{n!} C_x(n, \alpha) C_y(n, \alpha) \int_0^{+\infty} e^{-bn\tau} f_t(\tau) d\tau \\
&= m(x) \sum_{n=0}^{+\infty} \frac{\alpha^n}{n!} C_x(n, \alpha) C_y(n, \alpha) E_\nu(-bnt^\nu).
\end{aligned}$$

Finally we have

$$\begin{aligned}
\mathbb{P}(N_\nu(t) \in B | N_\nu(0) = y) &= \int_0^{+\infty} \sum_{x \in B} p_1(\tau, x; y) f_t(\tau) d\tau \\
&= \sum_{x \in B} \int_0^{+\infty} p_1(\tau, x; y) f_t(\tau) d\tau \\
&= \sum_{x \in B} m(x) \sum_{n=0}^{+\infty} \frac{\alpha^n}{n!} C_x(n, \alpha) C_y(n, \alpha) E_\nu(-bnt^\nu).
\end{aligned}$$

Thus $p_\nu(t, x; y)$ exists and

$$p_\nu(t, x; y) = m(x) \sum_{n=0}^{+\infty} \frac{\alpha^n}{n!} C_x(n, \alpha) C_y(n, \alpha) E_\nu(-bnt^\nu). \quad \square$$

Now we are ready to prove the following Theorem.

Theorem 5.2. *Let $g \in c_0$ such that $g(x) = \sum_{n=0}^{+\infty} g_n Q_n(x)$. Then the function*

$$u(t; x) = \mathbb{E}[g(N_\nu(t)) | N_\nu(0) = x]$$

is a solution of (16).

Proof. First of all observe that

$$\sum_{x=0}^{+\infty} m(x)g^2(x) \leq \|g\|_{\ell^\infty}^2,$$

hence $g \in \ell^2(\mathbf{m})$ and we are under the hypotheses of Theorem 4.3.

Consider a generic $f \in c_0$ and define the family of operators

$$T_t f(x) = \mathbb{E}[f(N_1(t)) | N_1(0) = x] = \sum_{y \in \mathbb{N}_0} p_1(t, y; x) f(y).$$

In particular, by Lemma 3.3, we know that $N_1(t)$ is a Feller process, hence $(T_t)_{t \geq 0}$ is a Feller semigroup. Moreover, strong continuity of $(T_t)_{t \geq 0}$ follows from [10, Lemma 1.4]. Then, by using [6, Theorem 3.1], we know that, since \mathcal{G} is the generator of T_t , the function

$$u(t; x) := \int_0^{+\infty} T_{(\frac{t}{s})^\nu} g(x) g_\nu(s) ds,$$

where $g_\nu(s)$ is the density of $\sigma_\nu(1)$, is a solution of (16). But if we use the change of variables $\tau = (\frac{t}{s})^\nu$, and the fact that $f_t(\tau) = \frac{t}{\nu} \tau^{-1-\frac{1}{\nu}} g_\nu(t\tau^{-\frac{1}{\nu}})$ for $\tau \geq 0$ (see, for instance, [33]), we obtain

$$\begin{aligned} u(t; x) &= \frac{t}{\nu} \int_0^\infty T_\tau g(x) \tau^{-1-\frac{1}{\nu}} g_\nu(t\tau^{-\frac{1}{\nu}}) d\tau \\ &= \int_0^{+\infty} T_\tau g(x) f_t(\tau) d\tau \\ &= \int_0^{+\infty} \mathbb{E}[g(N_1(\tau)) | N_1(0) = x] f_t(\tau) d\tau \\ &= \mathbb{E}[g(N_1(L_\nu(t))) | N_1(0) = x] = \mathbb{E}[g(N_\nu(t)) | N_\nu(0) = x], \quad t \geq 0, x \in \mathbb{N}. \quad \square \end{aligned}$$

Finally, we can provide the stochastic representation of solutions of (21).

Corollary 5.3. *Let $p_\nu(t, x; y)$ be the transition density of $N_\nu(t)$. Then, for any f such that $\frac{f}{m} \in \ell^2(\mathbf{m})$ with decomposition $f(x)/m(x) = \sum_{n=0}^{+\infty} f_n Q_n(x)$. Thus*

$$u(t, x) = \sum_{y \in \mathbb{N}_0} p_\nu(t, x; y) f(y)$$

is a solution of (21).

Proof. This easily follows from Theorems 5.1 and 4.4. \square

We can use the last Corollary to exploit the asymptotic behaviour of the density of the process $N_\nu(t)$.

6. Uniqueness of strong solutions

In this section, we aim to show that the strong solutions of (16) and (21) are unique under some hypotheses.

Proposition 6.1. *The function $p_\nu(t, x; y)$ given in (12) is the unique global solution of (23) (for fixed x) and (24) (for fixed y).*

Proof. Let us first notice that, by Theorem 5.1 we know that $0 \leq p_\nu(t, x; y) \leq 1$ and

$$\sum_{x=0}^{+\infty} p_\nu(t, x; y) = 1, \quad \sum_{y=0}^{+\infty} p_\nu(t, x; y) \leq 1,$$

hence $\|p_\nu(t, \cdot; y)\|_{\ell^1} = 1$ and $\|p_\nu(t, x; \cdot)\|_{\ell^1} \leq 1$. Thus $p_\nu(t, x; y)$ is bounded in ℓ^1 uniformly with respect to $t \geq 0$ for fixed x or y . Since ℓ^1 is continuously embedded in $\ell^2(\mathbf{m})$ (see Remark 3.1), then $p_\nu(t, x; y)$ is uniformly bounded also in $\ell^2(\mathbf{m})$. Moreover, we have shown in Lemma 3.2 that the operators \mathcal{G} and \mathcal{L} are continuous. Hence by [4, Corollary 2] we can conclude that $p_\nu(t, x; y)$ is the unique global solution of (23) and (24). \square

Now let us show the uniqueness of the solutions of the backward equation (16).

Proposition 6.2. *Let $g \in \ell^\infty$ such that $g(x) = \sum_{n=0}^{+\infty} g_n Q_n(x)$. Then the strong solution $u(t, x)$ of (16) is in ℓ^∞ , hence also in $\ell^2(\mathbf{m})$, for any $t \geq 0$ and it is the unique global solution in $\ell^2(\mathbf{m})$.*

Proof. First of all, let us observe that if $g \in \ell^\infty$, then $g \in \ell^2(\mathbf{m})$ too, so we are under the hypotheses of Theorem 4.3. Moreover we have, by using Theorem 5.1 and Jensen inequality:

$$\begin{aligned} u^2(t, x) &= \left(\sum_{y=0}^{+\infty} p_\nu(t, y; x) g(y) \right)^2 \\ &\leq \sum_{y=0}^{+\infty} p_\nu(t, y; x) g^2(y) \leq \|g\|_{\ell^\infty}^2, \end{aligned}$$

and then

$$\sum_{x=0}^{+\infty} m(x) u^2(t, x) \leq \|g\|_{\ell^\infty}^2,$$

obtaining the uniform bound for $x \mapsto u(t, x)$. Hence $u(t, \cdot) \in \ell^2(\mathbf{m})$ for any $t \geq 0$. Moreover, since \mathcal{G} is a continuous operator, by [4, Corollary 2], it is the unique global solution of (16). \square

Remark 6.3. We can actually show uniqueness as $g \in \ell^2(\mathbf{m})$. Indeed, since $u(t, x)$ is a strong solution of (16), then it belongs to $\ell^2(\mathbf{m})$ and

$$\|u(t, \cdot)\|_{\ell^2(\mathbf{m})}^2 = \sum_{n=0}^{+\infty} E_\nu^2(-bnt^\nu) g_n^2 \leq \|g\|_{\ell^2(\mathbf{m})}^2$$

for any $t \geq 0$, hence it is uniformly bounded in $\ell^2(\mathbf{m})$, concluding uniqueness by [4, Corollary 2].

We can also obtain the uniqueness of solutions of (21).

Proposition 6.4. Let $f : \mathbb{N}_0 \rightarrow \mathbb{R}$ be a function such that $\frac{f}{m} \in \ell^2(\mathbf{m})$ and $\frac{f(x)}{m(x)} = \sum_{n=0}^{+\infty} f_n Q_n(x)$. Then the strong solution $u(t, x)$ of (21) is in ℓ^∞ , hence in $\ell^2(\mathbf{m})$, and it is the unique global solution in $\ell^2(\mathbf{m})$.

Proof. Let us observe that

$$u^2(t, x) = \left(\sum_{y=0}^{+\infty} p_\nu(t, x; y) f(y) \right)^2 = \left(\sum_{y=0}^{+\infty} m(y) p_\nu(t, x; y) \frac{f(y)}{m(y)} \right)^2.$$

Thus, by Jensen inequality, we have

$$u^2(t, x) \leq \sum_{y=0}^{+\infty} m(y) p_\nu^2(t, x; y) \frac{f^2(y)}{m^2(y)}.$$

Now, from Theorem 5.1, we know that $p_\nu(t, x; y) \leq 1$, then

$$u^2(t, x) \leq \sum_{y=0}^{+\infty} m(y) \frac{f^2(y)}{m^2(y)} = \|f/m\|_{\ell^2(\mathbf{m})}^2.$$

Finally, we have

$$\sum_{x=0}^{+\infty} m(x) u^2(t, x) \leq \|f/m\|_{\ell^2(\mathbf{m})}^2 \sum_{x=0}^{+\infty} m(x) = \|f/m\|_{\ell^2(\mathbf{m})}^2$$

thus, since \mathcal{L} is a continuous operator, from [4, Corollary 2] we have that $u(t, x)$ is the unique global solution of (21). \square

Remark 6.5. The condition $f/m \in \ell^2(\mathbf{m})$ is stronger than $f \in \ell^2$ for any probability measure \mathbf{m} on \mathbb{N}_0 . Indeed we can show that the following two properties

- a) $f \in \ell^2$;
- b) $f/\sqrt{m} \in \ell^2(\mathbf{m})$;

are equivalent: this can be done simply observing that

$$\sum_{x=0}^{+\infty} f^2(x) = \sum_{x=0}^{+\infty} m(x) \left(\frac{f(x)}{\sqrt{m(x)}} \right)^2.$$

Moreover, if we consider the property

- c) $f/m \in \ell^2(\mathbf{m})$;

we can see that c) \Rightarrow a). Indeed we have, since $m(x) \leq 1$

$$\sum_{x=0}^{+\infty} f^2(x) = \sum_{x=0}^{+\infty} m^2(x) \left(\frac{f(x)}{m(x)} \right)^2 \leq \sum_{x=0}^{+\infty} m(x) \left(\frac{f(x)}{m(x)} \right)^2.$$

However, if we consider $f(x) = \sqrt{m(x)}$, it is easy to verify that $f \in \ell^2$ but $f/m \notin \ell^2(\mathbf{m})$.

7. Limit distribution of $N_\nu(t)$

In this section we want to give some results on the limit distribution of $N_\nu(t)$. In particular we have

Theorem 7.1. *Let $p_\nu(t, x; y)$ be the transition probability mass of $N_\nu(t)$. Then, given any initial probability mass f such that $f/m \in \ell^2(\mathbf{m})$ and $f(x)/m(x) = \sum_{n \in \mathbb{N}_0} f_n Q_n(x)$, the probability mass of $N_\nu(t)$ asymptotically converges towards a Poisson measure, that is to say if $p_\nu(t, x) = \sum_{y \in \mathbb{N}_0} p_\nu(t, x; y) f(y)$, then*

$$\lim_{t \rightarrow +\infty} p_\nu(t, x) = m(x).$$

Proof. By Corollary 5.3 we know that $p_\nu(t, x)$ is solution of (21). Moreover, since $f/m \in \ell^2(\mathbf{m})$, we know that this solution is unique from Proposition 6.4. Finally, from Theorem 4.4, we have that

$$p_\nu(t, x) = m(x) \sum_{n \in \mathbb{N}_0} E_\nu(-bnt^\nu) f_n Q_n(x).$$

In particular we have

$$p_\nu(t, x) = m(x) f_0 Q_0(x) + m(x) \sum_{n \in \mathbb{N}} E_\nu(-bnt^\nu) f_n Q_n(x).$$

But $Q_0(x) = 1$ and $f_0 = \sum_{x \in \mathbb{N}_0} f(x) = 1$ since f is a probability mass. Thus we have

$$p_\nu(t, x) = m(x) + m(x) \sum_{n \in \mathbb{N}} E_\nu(-bnt^\nu) f_n Q_n(x).$$

Now, by Cauchy-Schwartz inequality, the fact that $E_\nu(-bnt^\nu) \leq 1$ and the duality formula for $C_n(x, \alpha)$, we obtain

$$\begin{aligned} \sum_{n \in \mathbb{N}} |E_\nu(-bnt^\nu) f_n Q_n(x)| &\leq \sum_{n \in \mathbb{N}} |f_n Q_n(x)| \\ &\leq \|f/m\|_{\ell^2(\mathbf{m})} \left(\sum_{n \in \mathbb{N}} \frac{\alpha^n}{n!} C_n^2(x, \alpha) \right)^{\frac{1}{2}} \\ &\leq e^{\frac{\alpha}{2}} \|f/m\|_{\ell^2(\mathbf{m})} d_x^2 \end{aligned}$$

hence the second series totally converges. Thus we can take the limit inside the series and, since $\lim_{t \rightarrow +\infty} E_\nu(-bnt^\nu) = 0$, we have

$$\lim_{t \rightarrow +\infty} p_\nu(t, x) = m(x) + m(x) \sum_{n \in \mathbb{N}} \lim_{t \rightarrow +\infty} E_\nu(-bnt^\nu) f_n Q_n(x) = m(x). \quad \square$$

From Theorem 7.1 we know that whatever is the distribution of $N_\nu(0)$, the limit distribution of $N_\nu(t)$ is always \mathbf{m} . Moreover, we can show that \mathbf{m} is an invariant one-dimensional distribution for $N_\nu(t)$, that is to say that if $N_\nu(0)$ has distribution \mathbf{m} , then $N_\nu(t)$ admits \mathbf{m} as distribution for any $t > 0$.

Proposition 7.2. *Suppose $N_\nu(0)$ has distribution \mathbf{m} . Then for any $t > 0$, $N_\nu(t)$ has distribution \mathbf{m} .*

Proof. Let us observe that the density of $N_\nu(t)$ is given by

$$p_\nu(t, x) = \sum_{y \geq 0} p_\nu(t, x; y) m(y).$$

From Theorem 5.1 we have that

$$p_\nu(t, x; y) = m(x) \sum_{n=0}^{+\infty} E_\nu(-bnt^\nu) Q_n(x) Q_n(y)$$

then we have

$$\begin{aligned} p_\nu(t, x) &= \sum_{y \geq 0} \left(m(x) \sum_{n=0}^{+\infty} E_\nu(-bnt^\nu) Q_n(x) Q_n(y) \right) m(y) \\ &= m(x) \sum_{n=0}^{+\infty} E_\nu(-bnt^\nu) Q_n(x) \left(\sum_{y=0}^{+\infty} Q_n(y) m(y) \right). \end{aligned}$$

Recalling that $Q_0(y) \equiv 1$ we have that

$$\sum_{y=0}^{+\infty} Q_n(y) m(y) = \sum_{y=0}^{+\infty} Q_0(y) Q_n(y) m(y) = \delta_{n,0}$$

hence

$$p_\nu(t, x) = m(x) \sum_{n=0}^{+\infty} E_\nu(-bnt^\nu) Q_n(x) \delta_{n,0} = m(x). \quad \square$$

However, since $N_\nu(t)$ is not Markovian, this Proposition does not guarantee the stationarity of the process when $N_\nu(0)$ admits \mathbf{m} as distribution. However, it is still possible to compute the autocovariance function of the process $N_\nu(t)$.

Proposition 7.3. *Suppose $N_\nu(t)$ admits \mathbf{m} as initial distribution. Then, for any $t \geq s > 0$, it holds*

$$\text{Cov}(N_\nu(t), N_\nu(s)) = \alpha \left(E_\nu(-bt^\nu) + \frac{bvt^\nu}{\Gamma(1+\nu)} \int_0^{\frac{s}{t}} \frac{E_\nu(-bt^\nu(1-z)^\nu)}{z^{1-\nu}} dz \right). \quad (25)$$

We omit the proof of this Proposition since it is identical to the one in [26], after observing that if $N_1(t)$ admits \mathbf{m} as initial distribution, then $N_1(t)$ is stationary and, from (7),

$$\text{Cov}(N_1(t), N_1(0)) = \alpha e^{-bt}.$$

Remark 3.2 and 3.3 of [26] easily apply also to our process $N_\nu(t)$. Indeed, since in this case $N_\nu(t)$ is distributed as $N_\nu(0)$, then the variance $\mathbb{D}[N_\nu(t)] = \mathbb{D}[N_\nu(0)] = \alpha$, which can be obtained from (25) when $t = s$ with the same calculations as in [26, Remark 3.2]. Moreover, $N_1(t)$ exhibits short-range dependence, while, with the same calculations of [26, Remark 3.3], one can show that $\text{Cov}(N_\nu(t), N_\nu(s))$ decays as a power of t , hence it exhibits long-range dependence.

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