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The effect of *para*-substitution upon the structural and electronic properties of a series of *m*-terphenyl lithium complexes $[R-Ar^{\#}-Li]_2$ ($R = t\text{-Bu}$ **1**, SiMe_3 **2**, H **3**, Cl **4**, CF_3 **5**; where $R-Ar^{\#} = 2,6\text{-}\{2,6\text{-Xyl}\}_2\text{-}4\text{-}R\text{-C}_6\text{H}_2$ and $2,6\text{-Xyl} = 2,6\text{-Me}_2\text{C}_6\text{H}_3$) has been investigated. X-ray crystallography reveals the complexes to be structurally similar, with little variation in C–M–C bond lengths and angles across the series. However, in-depth NMR spectroscopic studies reveal notable electronic differences, showing a linear correlation between the $^7\text{Li}\{^1\text{H}\}$ NMR chemical shifts of the *para*-substituted complexes and their Hammett constants. The flanking methyl protons exhibit a similar electronic shift in the ^1H NMR spectra, which has been rationalised by the presence of through-space Li...H interactions, as evidenced by two-dimensional $^7\text{Li}\text{--}^1\text{H}$ heteronuclear Overhauser spectroscopy (HOESY). In both cases, electron-withdrawing substituents are found to cause an upfield peak shift. A computational analysis is employed to account for these trends.

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1. Introduction

The development of sterically demanding ligand systems has become a central research theme for the stabilisation of low-coordinate metal complexes. One important example is the *m*-terphenyl ligand, which has been repeatedly utilised in the isolation of low-coordinate main group and transition metal complexes.^{1–6} Although a plethora of *m*-terphenyl frameworks have been designed, most studies have been structurally focused, aimed at investigating the effects of steric bulk on the geometries, bonding modes, and reactivities of the resulting compounds.^{7–18}

For instance, the solid state structures observed for a series of *m*-terphenyl lithium complexes vary depending on the steric demands of the flanking *ortho*-aryl substituents.^{19,20}

Increasing the steric bulk of $2,6\text{-Ar}_2\text{C}_6\text{H}_3\text{Li}$ alters its aggregation state from a dimer,¹⁹ to a more crowded dimer featuring η^6 -arene coordination of the flanking aryl groups,²⁰ to a monomer stabilised by a coordinated molecule of benzene,²⁰ for $\text{Ar} = \text{Mes}$ ($2,4,6\text{-Me}_3\text{C}_6\text{H}_2$), Dipp ($2,6\text{-}i\text{-Pr}_2\text{C}_6\text{H}_3$) and Tripp ($2,4,6\text{-}i\text{-Pr}_3\text{C}_6\text{H}_2$) respectively.^{5,6}

However, less work has been reported on varying the electronic effects of these *m*-terphenyl systems, with studies limited mainly to a handful of main group complexes^{21–25} and the quintuply-bonded Cr–Cr dimer $[(2,6\text{-Dipp}_2\text{-}4\text{-}R\text{-C}_6\text{H}_2)\text{Cr}]_2$ ($R = \text{H}, \text{SiMe}_3, \text{OMe}, \text{F}$).²⁶ Even their lithium precursors, while structurally characterised, have not been studied from an electronic viewpoint. Therefore, it is the objective of this research to develop a toolbox of *m*-terphenyl ligands featuring a range of *para*-substituents (R) to investigate the electronic effects of substitution upon their metal complexes.

To this end, a series of *para*-substituted *m*-terphenyl lithium complexes $[R-Ar^{\#}-Li]_2$ ($R-Ar^{\#} = 2,6\text{-}\{2,6\text{-Xyl}\}_2\text{-}4\text{-}R\text{-C}_6\text{H}_2$; $R = t\text{-Bu}, \text{SiMe}_3, \text{Cl}, \text{CF}_3$; $2,6\text{-Xyl} = 2,6\text{-Me}_2\text{C}_6\text{H}_3$) are reported, and discussed alongside the previously published unsubstituted analogue $[\text{H}-Ar^{\#}-Li]_2$.²⁷ The geometric and electronic structures of these compounds are elucidated through X-ray crystallographic and NMR spectroscopic studies, respectively. Specifically, ^7Li NMR spectroscopy is employed to assess the electronic effects directly at the metal centre, with two-dimensional $^7\text{Li}\text{--}^1\text{H}$ heteronuclear Overhauser spectroscopy (HOESY) measurements demonstrating the presence of through-space $^7\text{Li}\cdots^1\text{H}$ interactions. Computational modelling with density

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functional theory (DFT) is also employed to help rationalise the observed trends in NMR parameters across the series.

2. Results and discussion

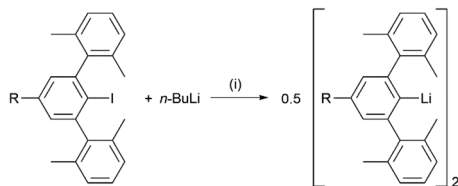
2.1. Synthesis

A series of *m*-terphenyl iodides R-Ar[#]-I (R = *t*-Bu, SiMe₃, Cl, CF₃)²⁸ was synthesised *via* similar procedures to other terphenyl compounds,^{22,25–27,29–33} and H-Ar[#]-I was obtained following a previously reported method^{15,32} (see ESI, section S2†). Lithium-halogen exchange reactions performed on these iodide compounds with excess *n*-butyllithium in *isohexane* at 0 °C afforded the *m*-terphenyl lithium complexes [R-Ar[#]-Li]₂ (R = *t*-Bu **1**, SiMe₃ **2**, H **3**, Cl **4**, CF₃ **5**) as white powders, according to Scheme 1. A range of yields was obtained (38%, 60%, 99%, quantitative and 95% for **1–5** respectively); the low yield of **1** can be ascribed to the greater solubility of the *t*-Bu substituted complex in hexane. Recrystallisation of **1**, **2**, **4** and **5** from a –30 °C *isohexane* solution produced colourless crystals suitable for X-ray diffraction analysis. The crystal structure of **3** has previously been reported.²⁷ Complexes **1–5** have been characterised by multinuclear (¹H, ¹³C{¹H}, ¹⁹F, ²⁹Si, ⁷Li{¹H}) NMR spectroscopies, diffusion-ordered spectroscopy (DOSY), and ⁷Li-¹H HOESY experiments.

2.2. Crystallographic characterisation

Single crystal X-ray diffraction data were collected for all new iodide ligand precursors, as well as the new lithium complexes **1**, **2**, **4**, and **5**. The crystal structure of **1** is presented in Fig. 1 and full crystallographic data are provided in the ESI (section S4, Fig. S30–S32†). The crystallographic data confirms that all complexes adopt a dimeric structure featuring two anionic *m*-terphenyl units linked by two lithium cations to form a planar Li₂C₂ core (see Fig. S31, S32 and Table S1 in the ESI†). Similar Li₂C₂ cores are observed across the series, with Li(1)⋯Li(2) distances of 2.304(3)–2.332(4) Å, C_{ipso}⋯C_{ipso} distances between 3.6893(18)–3.702(3) Å, and C_{ipso}-Li-C_{ipso} angles in the range of 114.59(13)–116.60(10)° [Σ_{Li-C-Li-C} = 360°]. Furthermore, all C_{ipso}-Li bond lengths occur within a narrow range [2.158(3)–2.1985(18) Å] and are comparable to the unsubstituted analogue **3** [2.143(5)–2.187(6) Å] and other *m*-terphenyl lithium dimers.^{15,19,27}

Weaker intramolecular interactions are also observed between the lithium ions and the xylyl flanking groups, with



Scheme 1 Synthesis of lithium complexes [R-Ar[#]-Li]₂ (R = *t*-Bu **1**, SiMe₃ **2**, H **3**, Cl **4**, CF₃ **5**). Reaction conditions: (i) *isohexane*, 0 °C to room temperature, 16 h, –*n*-BuLi.

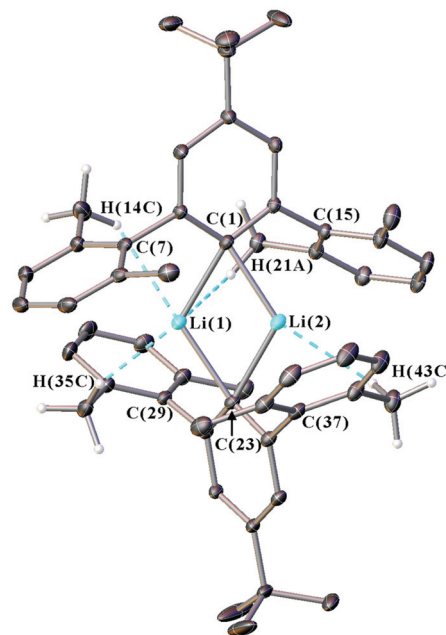


Fig. 1 Crystal structure of the *m*-terphenyl lithium complex **1**. Dashed lines indicate the short Li⋯H-C anagostic contacts. Ellipsoids set at 40% probability, and the hydrogen atoms are placed in idealised positions for structure refinement. All non-anagostic hydrogen atoms and residual solvent molecules are omitted for clarity.

secondary contacts that are shorter than the sum of the van der Waals radii for lithium and hydrogen (3.02 Å) or lithium and carbon (3.52 Å).^{34,35} A summary of these contacts is given in Table S2 of the ESI.† For complexes **1**, **2**, **4** and **5**, the lithium ions interact with the *ipso*-carbons of the flanking aryls [2.4244(15)–2.729(4) Å], and also form multiple Li⋯H-C anagostic interactions³⁶ with the hydrogens of the flanking methyl groups [2.338(2)–2.920(3) Å].

In summary, the crystal structures of **1**, **2**, **4** and **5** show little structural variation as the *para*-substituent is changed, which suggests that the geometries of *m*-terphenyl organolithiums are dominated by steric rather than electronic factors.^{5,6,19,20} The structures do, however, show weak Li⋯H-C anagostic contacts between the lithium ions and the flanking methyl groups. These interactions have been explored further in our NMR spectroscopic investigations, as discussed below.

2.3. NMR spectroscopic characterisation

The electronic properties of **1–5** were studied by ¹H, ¹³C{¹H} and ⁷Li{¹H} NMR spectroscopies in *d*₆-benzene. Due to their poor solubilities the solutions were saturated, except for **1** whose *t*-Bu group aided dissolution. To determine the aggregation state of the complexes in solution, DOSY experiments were performed.^{37,38} Whilst the unsubstituted iodide H-Ar[#]-I gave a hydrodynamic radius (*r*_H) of 3.0 Å, the equivalent lithium complex **3** gave a larger *r*_H of 4.8 Å, which correlates with its approximate crystallographic dimensions (9.4 × 9.9 Å). The other *para*-substituted complexes yield similarly large radii that, except for **5**, increase with bulkier substituents. In



all cases, only one set of NMR spectroscopic resonances is observed, suggesting that only one aggregation state is present in solution for each of the complexes. It should be noted that the r_H values for **1**, **2**, **4**, and **5** deviate more from the crystallographic dimensions than **3**, which is likely due to the more ellipsoidal shape of the molecules (see Fig. S33, ESI†).³⁹ However, the overall increase of hydrodynamic radii for **1–5** compared to the monomeric iodide $H-Ar^{\#}-I$ suggests that all these lithium complexes are dimeric in d_6 -benzene solution. The results are summarised in Table 1.

The key 1H , $^{13}C\{^1H\}$, and $^7Li\{^1H\}$ NMR spectroscopic signals of **1–5** are summarised in Table 2, with the numbering scheme presented in Fig. 2. The 1H NMR spectra show four characteristic peaks, although in some cases the signals for the aromatic xylyl protons (H-7 and H-8) overlap. Relative to the iodides ($R-Ar^{\#}-I$), all resonances are shifted upfield due to the greater shielding provided by the anionic ligands. A comparison of the 1H NMR spectra across the series reveals three key features. Firstly, the *meta*-hydrogens on the central rings exhibit small variations in chemical shift (H-3 = 6.85, 7.04, 6.77, 6.78, 6.99 ppm for **1–5**, respectively) with no overall trend, see Table 2. Secondly, the xylyl aryl protons, H-7 (6.86–7.01 ppm) and H-8 (6.94–7.01 ppm), remain essentially

unshifted, suggesting that the substituent in the *para*-position has minimal electronic communication with the flanking aryl rings. Thirdly, the xylyl methyl protons shift upfield (H-9 = 1.83, 1.81, 1.80, 1.61, 1.55 ppm for **1–5**, respectively) with increased electron-withdrawing strength of the *para*-substituent. A graph of these chemical shifts (δ) against the corresponding Hammett constants (σ_{para})⁴⁰ reveals a linear correlation, see Fig. 4 (blue line; $R^2 = 0.92$). This trend is not observed in the respective iodide compounds (H-9 = 2.06, 2.05, 2.06, 2.07 and 2.02 ppm for $R-Ar^{\#}-I$, where R = *t*-Bu, SiMe₃, H, Cl and CF₃, respectively) and proceeds in the opposite direction to what one might expect, with electron-withdrawing groups causing an apparent shielding effect. This suggests that the *para*-substituent is exerting a direct electronic influence on the environment of the H-9 methyl protons.

The $^{13}C\{^1H\}$ NMR spectra of **1–5** show nine peaks for the carbons of the ligand framework, with additional peaks for the *t*-Bu, SiMe₃ and CF₃ groups in **1**, **2** and **5**, respectively. Compared to the iodides ($R-Ar^{\#}-I$), the largest change is seen in the *ipso*-carbons, which are shifted downfield (C-1 = 106.9 vs. 172.7 ppm for $H-Ar^{\#}-I$ and **3**, respectively) owing to a large deshielding effect in the plane perpendicular to the C–Li bond.^{19,41–44} In all cases, the NMR resonances for the *ipso*-carbons of **1–5** were of low intensity and therefore assigned from $^1H-^{13}C$ HMBC spectra. For **3**, a longer $^{13}C\{^1H\}$ NMR experiment with d_1 of 10 s was performed to achieve better resolution of the *ipso*-carbon peak (Fig. 3). This revealed a seven-line splitting pattern (blue) centred at 172.7 ppm with a 1 : 2 : 3 : 4 : 3 : 2 : 1 intensity ratio ($^1J_{^{13}C^7Li} = 23.3$ Hz) character-

Table 1 Diffusion coefficients, D , and hydrodynamic radii, r_H , obtained from the DOSY measurements of the monomeric $H-Ar^{\#}-I$ and the dimeric $[R-Ar^{\#}-Li]_2$ (**1–5**) species. Their crystallographic diameters, d_c , are provided for comparison

Compound ^a	Diffusion coefficient, D (10^{-10} m ² s ⁻¹)	Hydrodynamic radius, r_H (Å)	Crystallographic diameters, d_c (Å)	
			Length ^b	Width ^c
$H-Ar^{\#}-I$	12.30	3.0	4.9	9.9
$[H-Ar^{\#}-Li]_2$ (3)	7.54	4.8	9.4	9.9
$[Cl-Ar^{\#}-Li]_2$ (4)	6.72	5.4	12.8	9.9
$[F_3C-Ar^{\#}-Li]_2$ (5)	7.59	4.8	13.4	9.9
$[t-Bu-Ar^{\#}-Li]_2$ (1)	6.59	5.5	13.6	9.9
$[Me_3Si-Ar^{\#}-Li]_2$ (2)	6.27	5.8	14.4	9.9

^a All samples were prepared as 24 mM solutions in d_6 -benzene.

^b Length along the compound measured between $R\cdots I$ ($H-Ar^{\#}-I$) and $R\cdots R$ (**1–5**), where R is the outermost nucleus of the *para*-substituent.

^c Width across the compound measured between C \cdots C for the outermost carbon atoms of the flanking aryl rings on each *m*-terphenyl ligand. All widths averaged at 9.9 Å.

Table 2 Relevant 1H , $^{13}C\{^1H\}$ and $^7Li\{^1H\}$ NMR chemical shifts, δ , for the *m*-terphenyl lithium complexes $[R-Ar^{\#}-Li]_2$ (**1–5**). The flanking aryl atoms remain unshifted so have been omitted

$[R-Ar^{\#}-Li]_2$	R group	1H , $^{13}C\{^1H\}$ and $^7Li\{^1H\}$ NMR Chemical Shifts, δ (ppm)							
		H-3	H-9	C-1	C-2	C-3	C-4	C-9	Li
1	<i>t</i> -Bu	6.85	1.83	168.1	152.0	120.4	148.7	21.8	1.60
2	SiMe ₃	7.04	1.81	174.2	151.6	128.1	136.5	21.9	1.47
3	H	6.77	1.80	172.7	152.0	123.6	126.3	21.8	1.46
4	Cl	6.78	1.61	170.7	153.6	123.6	132.9	21.5	1.10
5	CF ₃	6.99	1.55	180.2	152.4	119.5	— ^a	21.5	0.93

^a Peak C-4 for **5** not observed (obscured by d_6 -benzene solvent peak).

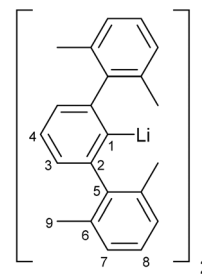


Fig. 2 NMR numbering scheme for the *m*-terphenyl lithium complexes **1–5**.



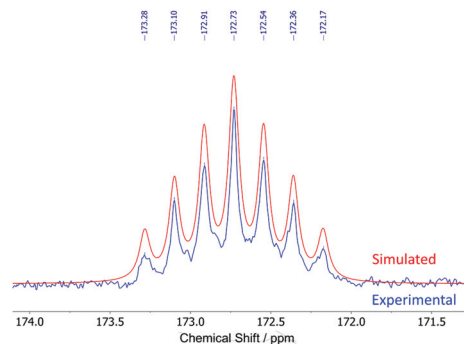


Fig. 3 Seven-line splitting pattern of the *ipso*-carbons in the experimental (blue) and simulated (red) $^{13}\text{C}\{^1\text{H}\}$ NMR spectra of $[\text{H}-\text{Ar}^\#-\text{Li}]_2$ (**3**).

istic of quadrupolar coupling to two ^7Li nuclei ($I = 3/2$, natural abundance: 92.41%), providing further evidence that the lithium complexes are dimeric in d_6 -benzene solution.^{45–48} The experimental spectrum was well simulated (red, see ESI, section S1†). Multiplets arising from ^6Li ($I = 1$, natural abundance = 7.59%) were too low intensity to resolve experimentally.

A comparison of the $^{13}\text{C}\{^1\text{H}\}$ NMR spectra of **1–5** reveals that all ^{13}C signals of the xylyl flanking groups remain essentially unshifted across the series. The shift observed for the xylyl methyl groups in the ^1H NMR spectrum (H-9) is not replicated in the $^{13}\text{C}\{^1\text{H}\}$ NMR spectrum (C-9). The peaks for the central ring carbons vary considerably across the series (except for C-2, see Table 2): larger variations are noted for C-4 (126.3–148.7 ppm) than for C-1 (168.1–180.2 ppm) and C-3 (119.5–128.1 ppm). Though no reliable trends can be identified, a general downfield shift is observed for C-1 with increasing σ_{para} .

The $^7\text{Li}\{^1\text{H}\}$ NMR spectra of **1–5** also provided valuable information on the complexes. Although ^7Li NMR spectroscopy is typically sensitive to the analyte concentration, aggregation, and solvent effects, these factors appear to be negligible here.^{49–51} This was demonstrated by the spectra for the more soluble *t*-Bu complex (**1**) whose $^7\text{Li}\{^1\text{H}\}$ NMR resonance remains unshifted over a wide concentration range (1.60 ppm at 24, 48 and 72 mM, respectively). This may be due to the steric bulk of the ligands shielding the metal centre from the surrounding solvent. For complexes **2–5**, saturated samples were used for NMR measurements due to their poor solubilities in d_6 -benzene. The steric influence of the xylyl flanking groups on the lithium ions is assumed to remain consistent throughout the series and, therefore, not to contribute to changes in the $^7\text{Li}\{^1\text{H}\}$ NMR chemical shifts. This is in contrast to previous work on *m*-terphenyl lithium complexes, where changing the steric bulk of the flanking groups varies the $^7\text{Li}\{^1\text{H}\}$ NMR signal.¹⁵

In all cases, the $^7\text{Li}\{^1\text{H}\}$ NMR spectra show a single environment for the two lithium ions of the dimer. The position of these peaks shifts upfield (1.60, 1.47, 1.46, 1.10 and 0.93 ppm for **1–5**, respectively) as the electron-withdrawing strength of

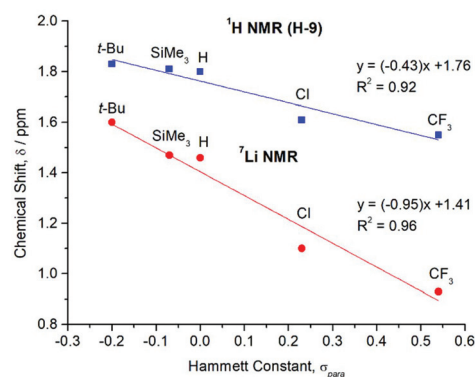


Fig. 4 Plot of the ^1H (for flanking methyl protons, H-9) and $^7\text{Li}\{^1\text{H}\}$ NMR chemical shifts (δ) for the *para*-substituted lithium complexes $[\text{R}-\text{Ar}^\#-\text{Li}]_2$ (**1–5**) versus their corresponding literature Hammett constants (σ_{para}).⁴⁰

the *para*-substituent is increased (Table 2). Once again, plotting a graph of these chemical shifts (δ) against the Hammett constants (σ_{para}) reveals a linear correlation, see Fig. 4 (red line; $R^2 = 0.96$),⁴⁰ indicating a direct influence on the electronic properties of the lithium ions by the *para*-substituent. As with the H-9 protons, the trend is counterintuitive, since one might expect electron-withdrawing groups to deshield the nuclei and cause a downfield shift. However, a similar observation has been reported for the series of *para*-substituted aryl-lithium complexes $[\text{4-R-C}_6\text{H}_4]_2\text{Li}$ ($\text{R} = \text{OMe}, \text{Me}, \text{H}, \text{F}, \text{Cl}, \text{CF}_3$).^{52–54}

These $^7\text{Li}\{^1\text{H}\}$ NMR spectroscopic results help to rationalise the trend observed for the flanking methyl protons (H-9) (Fig. 4). Since all other atoms on the xylyl rings remain unshifted, this suggests that electronic communication occurs *via* a through-space interaction between the methyl protons and lithium atoms. Some evidence of this is seen in the crystal structures, with weak $\text{Li}\cdots\text{H}-\text{C}$ anagostic interactions observed in the solid state (Fig. S31 and S32†). In solution, this is further corroborated by two-dimensional $^7\text{Li}-^1\text{H}$ HOESY experiments,^{55,56} which reveal short $^7\text{Li}\cdots^1\text{H}$ contacts in **1–5** as evidenced by the intense cross-peak between the lithium signals and the H-9 protons (Fig. 5 and ESI, section 3†).

Similar $\text{Li}\cdots\text{H}$ interactions have been reported for the monomeric $[\text{Mes}^*-\text{Li}]$ species, whose $^6\text{Li}-^1\text{H}$ HOESY spectrum gave a cross-peak for the *ortho tert*-butyl protons, but not for the *para tert*-butyl protons.⁴² Moreover, the *ortho* groups (1.61 ppm) occurred at lower field than the *para* groups (1.56 ppm), owing to the electric field produced by lithium that deshields the protons within close proximity.⁴² Comparable findings have been described for other aryl-lithium complexes $[\text{Ar}-\text{Li}]$ ($\text{Ar} = \text{Naph}, 2-\{t\text{-BuS}\}_2\text{C}_6\text{H}_4$).^{57,58}

From these results we suggest that the electronic influence of the *para*-substituents in complexes **1–5** is relayed through the central *m*-terphenyl ring, onto the lithium atoms, then through-space onto the nearby methyl protons. Evidence of similar through-space interactions within a *m*-terphenyl framework has already been reported for the *m*-terphenyl carboxylic



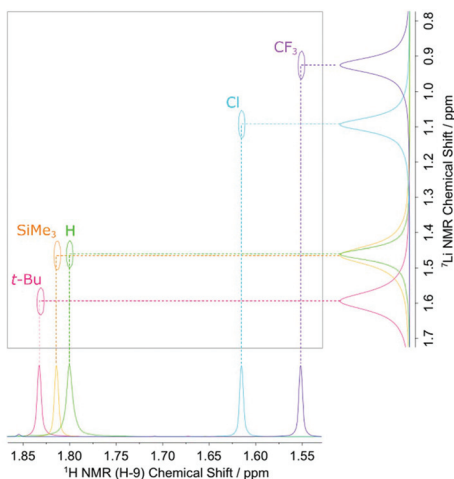


Fig. 5 Superimposed 2D ${}^7\text{Li}$ - ${}^1\text{H}$ HOESY spectra of the *m*-terphenyl lithium complexes $[\text{R}-\text{Ar}^\#-\text{Li}]_2$ (1–5) showing strong cross-peaks between the flanking methyl protons, H-9, and the ${}^7\text{Li}$ nuclei.

acids (2,6- $\text{Ar}_2\text{C}_6\text{H}_3$)COOH ($\text{Ar} = 4\text{-R-C}_6\text{H}_4$; $\text{R} = \text{OMe, Me, H, F, Cl, Br, C}\{\text{O}\}\text{Me}$) where *para*-substitution of the flanking aryl rings influences the $\text{p}K_a$ value, owing to through-space effects.⁵⁹ In addition, the *m*-terphenyl silane (2,6- $\text{Ar}_2\text{C}_6\text{H}_3$) SiMe_2H ($\text{Ar} = 2,6\text{-F}_2\text{C}_6\text{H}_3$) possesses flanking aryl fluorine atoms that couple through-space to the Si-H proton.⁶⁰ To further investigate the origin of this effect we turned to computational analysis and DFT studies.

2.4. Computational studies

The calculation of NMR chemical shift parameters was carried out using the ReSpect program^{61–66} on structures obtained directly from the crystal structures, with optimisation of only the H atom locations at the PBE0 level;^{67,68} see ESI, section S5† for full computational details. NMR shielding constants were calculated using the KT2 density-functional approximation,⁶⁹ in the pcS- n basis sets⁷⁰ ($n = 1, 2$). This functional was specifically designed for the calculation of NMR shielding constants, as were the pcS- n basis sets. Comparisons of the results for a range of other commonly used functionals (BLYP, BP86, PBE, PBE0, PP86)^{67,68,71–74} may be found in the ESI.†

Tables summarising the calculated shifts for H-9 and Li are given in the ESI, Tables S6 and S7.† Because of the relatively narrow chemical shift range, it is difficult to accurately model the trends shown in Fig. 3. Furthermore, both the H-9 and Li shifts are composed from absolute shielding constants containing large paramagnetic components (ESI, Table S5†). It is known that when the paramagnetic components are dominant, as in these cases, then density-functional methods often fail to achieve high accuracy. It is also known that density-functionals tend to be more accurate for ${}^1\text{H}$ and ${}^{13}\text{C}\{{}^1\text{H}\}$ NMR resonances than those for other nuclei.⁷⁵ In line with these observations the calculated values for the ${}^7\text{Li}\{{}^1\text{H}\}$ NMR spectrum do not accurately capture the trend in Fig. 3. However, the trend for the ${}^1\text{H}$ NMR of the flanking H-9 nuclei is more

adequately reproduced, particularly for the KT2 functional in the larger pcS-2 basis. As shown in Fig. S34 of the ESI† the general trend for an upfield shift with increasing σ_{para} is reproduced, see also Table S7.† Linear regression of the calculated ${}^1\text{H}$ NMR shifts for the H-9 nuclei as a function of σ_{para} yields a fit with $R^2 = 0.92$, a somewhat steeper slope of -0.62 and a y -intercept of $+1.82$ ppm. Given the challenges associated with the relatively narrow range of chemical shifts and large paramagnetic contributions for these nuclei, these results are in reasonable agreement with the experimental observations. However, based on these results, it seems likely that the counterintuitive trends observed in both the ${}^1\text{H}$ and ${}^7\text{Li}\{{}^1\text{H}\}$ NMR resonances is due to a large paramagnetic contribution to the shielding constant. Finally, we note that the more challenging ${}^7\text{Li}\{{}^1\text{H}\}$ NMR spectroscopy may provide a useful test case for the development of more refined density-functionals for NMR studies.

3. Conclusions

A series of *para*-substituted *m*-terphenyl lithium complexes $[\text{R}-\text{Ar}^\#-\text{Li}]_2$ ($\text{R} = t\text{-Bu}$ 1, SiMe_3 2, H 3, Cl 4, CF_3 5) have been reported. Whilst crystallographically similar, and all displaying $\text{Li}\cdots\text{H}-\text{C}$ anagostic contacts in the solid state, NMR spectroscopic studies reveal significant electronic differences at the metal centre. A linear correlation was observed between the Hammett constants of the *para*-substituent and the ${}^7\text{Li}$ and ${}^1\text{H}$ NMR shifts of the lithium atoms and flanking methyl protons (H-9). Electronic interaction between these atoms was confirmed by ${}^7\text{Li}$ - ${}^1\text{H}$ HOESY measurements. In both cases, increasing the electron-withdrawing power of the *para*-substituent results in a counterintuitive upfield peak shift. Computational analysis suggests this effect is due to an electric field generated about the lithium atoms, which is influenced by the *para*-substituent, and results in large paramagnetic and diamagnetic contributions to chemical shifts. While the observed trend in the ${}^1\text{H}$ NMR spectra of the H-9 protons was reproduced relatively well, it was difficult to accurately model the ${}^7\text{Li}$ NMR spectroscopic trend by DFT.

Conflicts of interest

There are no conflicts to declare.

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