

Research article

Tracking the creation of single photon emitters in AlN by implantation and annealing

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ABSTRACT

In this study, we inspect and analyse the effect of Al implantation into AlN by conducting confocal microscopy on the ion implanted regions, before and after implantation, followed by an annealing step. The independent effect of annealing is studied in an unimplanted control region, which showed that annealing alone does not produce new emitters. Through tracking individual locations in a lithographically patterned sample, we observe that point-like emitters are created in the implanted regions after annealing. The newly created quantum emitters show anti-bunching under ambient conditions and are spectrally similar to the previously discovered emitters in as-grown AlN.

1. Introduction

Quantum technologies including quantum sensing, communications and computing have emerged as a potentially paradigm-changing sector, attracting strong scientific interest and substantial investments [1]. The major reason for this interest is the quantum advantage over classical information processing paradigms for specific and demanding computational problems. The advantage relies heavily on finding platforms hosting qubits that can be reliably prepared, manipulated and measured [2]. One such platform is colour centres in wide-bandgap semiconductors. An extensively investigated system is the negatively charged nitrogen-vacancy centre in diamond (NV⁻) [3], which hosts an electron spin which can be manipulated using microwave fields and read out optically [4], even at room temperature.

Other colour centres in semiconductors are being investigated for narrow transitions, high brightness, longer coherence lifetimes or ease of photonics integration; notable examples are SiV [5] and SnV [6] in diamond, divacancy centre [7] and V_{Si} [8] in silicon carbide (SiC), T- [9] and G- [10] centre in silicon and carbon-related defects [11] in hexagonal boron nitride (hBN). Recently, quantum emitters have also been detected in III-nitrides, such as AlN [12] and GaN [13]. Owing to the unique qualities of III-nitrides such as tunable bandgap and piezoelectricity, these emitters are promising candidates for electronically- or mechanically-modulated single photon sources in ambient conditions.

AlN is hypothesised to host defect complexes that can be useful as potential qubits due to its large bandgap (6.02 eV). Varley et al. previously conducted an ab initio study on the vacancy complexes with group IVA and IVB elements, and reported that Ti-, Zr-vacancy centres can be hosted within the bandgap, with the ground state sufficiently apart from the valence band of AlN [14]. Previously conducted experiments in Ti-implanted AlN found an electron spin resonance signal in implanted samples but could not confirm optical activity [15]. Aghdaei et al. reported formation of broad peaks in Zr-implanted AlN [16], but in a comprehensive study Senichev et al. found that emission in Zr-implanted AlN is statistically not different than Kr-implanted AlN with similar ion damage [17]. In line with our previous observations in Al-implanted AlN [18], these results point to a native origin for the created single-photon emitters (SPE). While these studies demonstrated SPE creation by statistically comparing different samples, the creation event of a single distinct emitter was not microscopically assessed.

In this study, we rigorously inspect and analyse the effect of Al implantation into AlN by conducting confocal microscopy on the ion implanted regions, before and after implantation, followed by an annealing step. The independent effect of annealing is studied in an unimplanted control region, which showed that annealing alone does not produce new emitters. Through tracking individual locations in a lithographically patterned sample, we observe that point-like emitters

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are created in the implanted regions after annealing. The newly created quantum emitters show anti-bunching under ambient conditions and are spectrally similar to the previously discovered emitters in as-grown AlN.

2. Materials and methods

The studied samples consist of single crystal (0001) 1 μm MOCVD-grown AlN epilayer on 480 μm thick sapphire substrate, acquired from DOWA Electronics. SIMS measurements conducted by the manufacturer showed the main chemical impurities were H, C, O and Si, as is usual for samples grown by MOCVD. The epilayer in its as-grown form hosts single-photon emitters reported in our previous studies [12].

2.1. Confocal microscopy

The sample was inspected through confocal microscopy with a dry objective with NA = 0.9 under CW 520 nm excitation. The emission was filtered through a 550 nm long pass filter to suppress the laser line, and a 650 nm short pass filter to remove undesired background fluorescence from the sapphire substrate between 690–750 nm from the Cr³⁺ impurity. The confocal microscopy scans were conducted with below-saturation excitation power (215 μW at the back of the objective, compared to a typical saturation power of 606 μW) to increase emitter-to-background contrast. The collected emission was guided to a fibre-coupled Excelitas silicon avalanche photodiode for photon counting measurements. As the excitation laser has a well-defined polarisation, the emitters whose absorption dipoles align with the excitation polarisation were detected preferentially against the emitters that did not. To counter this inherent bias, we conducted the confocal scans under linear polarisations of 0°, 60° and 120° in the plane of the sample to maximise the probability of detecting changes in the scan maps as described in Cannon et al. [19]. The individual intensity maps were then spatially aligned each other and the individual pixels were quadrature summed across three different polarisations in MATLAB. We note that the brightness of different emitters are not uniform. To show clearly the presence of less intense emitters in the scan maps presented in the manuscript, we have consistently rescaled the colour bar to the range between 1% and 25% of the raw image range. This means the brighter emitters are above the top end of the scale, and single pixel noise in the regions between emitters is less visible.

For the Hanbury-Brown and Twiss (HBT) intensity interferometry, the filtered emission was split between two APDs using a multimode fibre beamsplitter. For each emitter, power dependent data was recorded to determine the best signal-to-noise ratio. Coincidence events from APD back-flash reflections were masked when fitting to the correlation functions. For the spectral measurements, the emission was coupled to an Andor Kymera 328i spectrometer with a grating of 150 l/mm and a silicon CCD camera through a polarisation-maintaining fibre. The emission spectrum was captured under maximum excitation power (3.7 mW) and laser polarisation aligned to the absorption dipole angle of the individual emitters.

2.2. Sample preparation

The as-grown epilayer of AlN contains a random distribution of point-like emitters which were mapped before implantation to track the creation of new colour centres. The epilayer was patterned using photolithography, a hardmask and a Cl₂/BCl₃ dry etch [20]. The hardmask was then removed and the sample cleaned in solvents. The sample was subjected to a 100 W 40 sccm O₂ ashing step to remove any surface contaminants. Two marked regions of 25 μm \times 25 μm , separated by 1 mm, were dedicated to the zero fluence (control) and 10¹⁵ ions/cm² implanted area on the same 5 mm \times 5 mm chip. The regions were subsequently inspected under scanning laser confocal microscopy.

The sample was implanted with 70 keV Al ions with a multi-element ion source at the University of Torino. The implantation region was implanted with an ion beam masked through a 400 μm \times 400 μm aperture for a fluence of 10¹⁵ ions/cm². The vacancy distribution within the epilayer was simulated with the SRIM software package [21], showing ion (vacancy) distribution peaking around 87 (57) nm below the epilayer surface [20]. The effect of initial ion implantation does not extend further than 200 nm into the material.

We conducted two annealing steps at 600 °C and 800 °C to thermally activate the implanted emitters. Both steps were conducted under nitrogen atmosphere and at atmospheric pressure with a JIPELEC JetFirst300 rapid thermal annealer in Cardiff University. The initial temperature ramp was set to 1600 °C/h and the sample stayed at the target temperature for 30 min, after which it was left to thermalise to room temperature. At all stages, the nitrogen flow was set to 300 sccm.

3. Results

3.1. Depth inspection on the native emitters

To interpret the impact of implantation and annealing in the epilayer, the depth distribution of the as-grown emitters was investigated. We took a direct approach and studied the density of emitters through repetitive material removal and confocal inspection steps. This enabled determining the depth of the emitters with accuracy determined by the etch, rather than by the Rayleigh range of the confocal mapping system.

Following the lithographic marking, the sample was inspected under scanning confocal photoluminescence microscopy. The 1 μm AlN epilayer showed an emitter density of 0.22 μm^{-2} . 18 of the bright and isolated emitters, all showing anti-bunched emission statistics, were selected for tracking in the subsequent inspections. The AlN layer was then iteratively etched and scanned, effectively capturing the depth-dependent emitter distribution within the epilayer. Between each etching step and confocal inspection, the sample was cleaned with a 100 W oxygen plasma for 5 min to remove any remaining etch products from the surface.

We observed no substantial change in the emitter distribution in the first few etch steps, as 93% of the epilayer is removed. Only one area (marked in Fig. 1a) changed in the first step. All emitters that were showing anti-bunched emission statistics (lack of coincidences between two detection events for small values of inter-path delay) were still observable at a layer thickness of 75 nm (Fig. 1c). The epilayer was then removed at a final etch step, exposing the sapphire substrate. Subsequent inspection revealed that all emitters had been removed (Fig. 1d). This suggests that the native emitters are located at the first 75 nm of AlN, close to the AlN/Sapphire surface. We note that this is the area in which the density of threading dislocations is greatest, and so emitter creation during MOCVD growth may be linked to this structural property as previously observed in GaN [22]. Alternatively, the changing Fermi level near the sapphire interface might be populating the defect levels; the inverse of this effect was observed in NV⁻¹ centres situated near Schottky barriers [23]. While the investigation of such effects is outside of the scope of this work, cathodoluminescence mapping of dislocations and electroluminescence experiments on as-grown defects can form the focal point of future studies.

3.2. Effect of annealing in the control region

Following the depth profiling of the as-grown emitters, we studied the effect of annealing through step-wise confocal microscopy. Fig. 2a shows the region prior to any annealing. The density of the native emitters was found to be 0.19 μm^{-2} . Following the annealing at 600 °C for 30 min (Fig. 2b), we observed that the brightness of the emitters detected in the first inspection was non-uniformly modified. Overall, annealing at 600 °C decreased the luminescence from the featureless

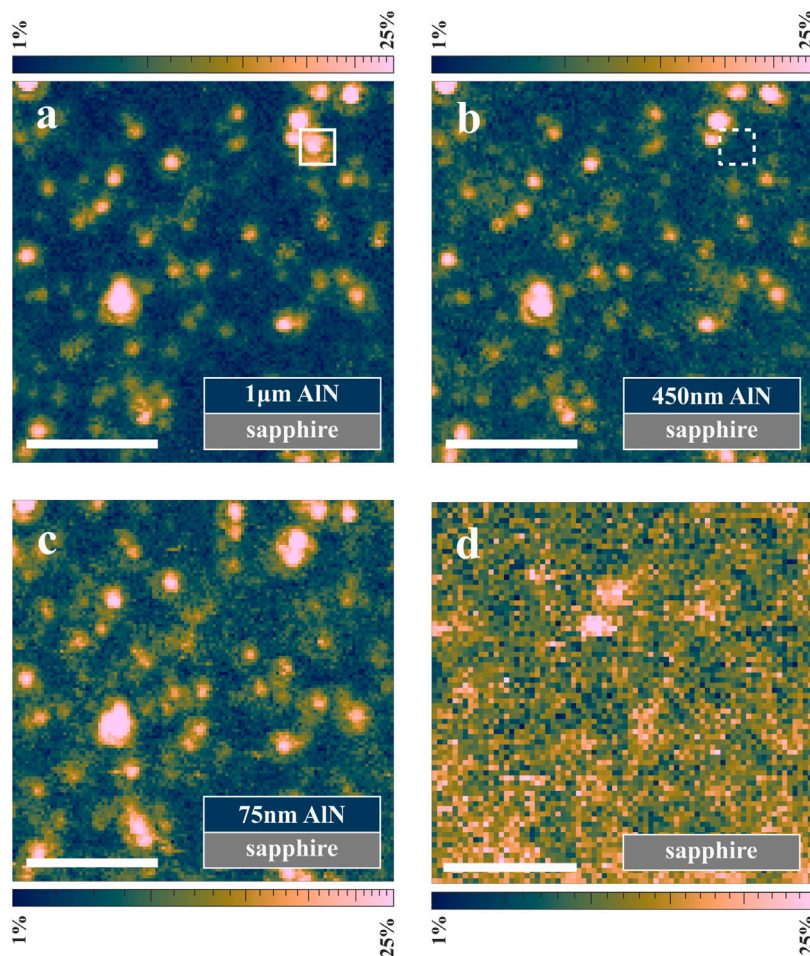


Fig. 1. Investigating the depth of quantum emitters in the as-grown material as it is etched. Scale bar is 5 μm . (a) The epilayer at its initial thickness of 1 μm . The emitter which disappears in (b) is marked. (b) The epilayer at a thickness of 450 nm. (c) The epilayer at a thickness of 75 nm. (d) The epilayer completely etched, with the substrate exposed.

“background” without removing or creating emitters. This causes the emitters in Fig. 2b to appear brighter compared to their pre-anneal scans (Fig. 2a). Similarly, annealing at 800 $^{\circ}\text{C}$ for 30 min (Fig. 2c) did not create any new emitters in the control region, but instead lowered the brightness of the native emitters. Removal of the top 500 nm of the epilayer (Fig. 2d) did not result in removal of the native emitters, which again confirms the emitters are located in the region of the epilayer closest to the sapphire interface. As a result, we concluded that annealing at 600 $^{\circ}\text{C}$ and 800 $^{\circ}\text{C}$ for 30 min did not have a significant effect on the unimplanted control region.

3.3. Effect of annealing on the implanted region

The region implanted with Al ions at 10^{15} cm^{-2} fluence can be seen after various stages of processing in Fig. 3. In panel (a), the native emitters in the implanted region before implantation are illustrated. Post-implantation damage resulted in a high “background” luminescence dominating the whole area with no spatial structure, which displayed exposure-dependent bleaching. This is common in implantation studies, and is a transient result of the ion damage [24].

After annealing at 600 $^{\circ}\text{C}$ for 30 min under nitrogen atmosphere (Fig. 3b), new isolated emitters were detected in the scan maps. Two of these emitters (E1, E2) are highlighted in Fig. 3b, which show antibunched photon statistics under ambient conditions. Both of the emitters displayed emission spectra similar to the previously-discovered emitters in AlN (Fig. 4a-b) [12], with E2 having a distinct ZPL at 611 nm. In total, 8 newly-created emitters showed up in the confocal scans over a $25 \mu\text{m} \times 25 \mu\text{m}$ area, signifying a low yield in the creation of

the emitters. Annealing at 800 $^{\circ}\text{C}$ for 30 min produced some additional effects. Firstly, 4 emitters observed at the previous step ceased to fluoresce, including E1. Secondly, some of the emitters, such as E2, while still observable, displayed changed emission spectrum (Fig. 4b, red). The second annealing step also caused diffusion of the previously-created emitters such as E4 and induced the creation of additional emitters (highlighted in Fig. 3c).

We did not observe a systematic difference between the spectral signatures of the as-grown emitters [12] and the newly-created emitters (Fig. 4). In both cases, the variability in the signatures might be originating from non-homogeneous host environments [25], occupation of different lattice sites of identical defects [26], or differences in defect species involved. The thermal stability of the as-grown emitters show that these emitters are in energetically favoured configurations within the accessible parameter space of our annealer. In contrast, the relative instability of E1-like emitters suggest that these optically active emitters are formed when the constituents are nudged into a meta-stable formation energy local minimum during limited thermal treatment, only to morph into optically inactive, but possibly energetically favoured defect configurations during prolonged heat treatment [10].

To confirm quantum-mechanical emission coming from the newly-created emitters, Hanbury-Brown and Twiss interferometry was conducted on E1 and E2 with the setup detailed in Section 2.1. The correlations were normalised and fit with an empirical equation based on the state dynamics of a three-level model with a metastable shelving state:

$$g^{(2)}(\tau) = 1 - a_1 * \exp^{-|\tau|/\lambda_1} + a_2 * \exp^{-|\tau|/\lambda_2} \quad (1)$$

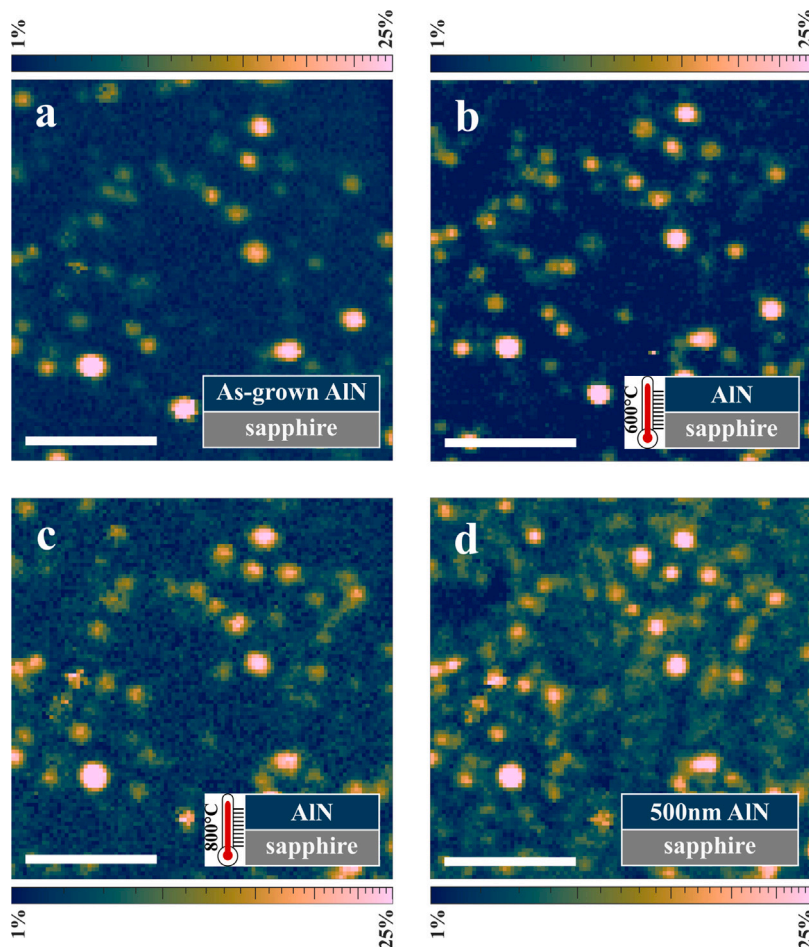


Fig. 2. Effect of annealing on the as-grown (unimplanted) control area of the AlN chip. (a) Control region before any thermal processing. (b) Control region after annealing at 600 °C for 30 min under N₂ atmosphere. (c) Control region after annealing at 800 °C for 30 min under N₂ atmosphere. (d) Control region after removing 500 nm of AlN epilayer. Scale bar is 5 μm.

where $a_{1,2}$ are the amplitudes and $\lambda_{1,2}$ are the lifetimes of antibunching and bunching contributions, respectively. We observe that both E1 and E2 display anti-bunching below the conventional limit used to define a single quantum emitter ($g_{E1}^{(2)}(0) = 0.375$, $g_{E2}^{(2)}(0) = 0.24$) with short lifetimes ($\lambda_1^{E1} = 2.56$ ns, $\lambda_1^{E2} = 4.46$ ns), albeit with significant bunching ($a_2 = 2.2$, $\lambda_2 = 10.5$ μs for E1; $a_2 = 0.63$, $\lambda_2 = 10.9$ μs for E2). Such behaviour is consistent with our previous work [27], where we report the existence of multiple long-lived metastable states in as-grown emitters in AlN on sapphire epilayers. While these metastable states are undesirable since they represent significant deviations from ideal two-level dynamics required for qubit applications and reduce the time-averaged intensity, their potentially spin-selective transitions to the ground level can be utilised to provide brightness contrast between different spin states in optically-detected magnetic resonance (ODMR) experiments.

To further inspect the depth distribution of the created emitters, the epilayer was thinned by 500 nm through dry etching. Surprisingly, the created emitters that still luminesce could still be observed after removing the upper half of the AlN epilayer, even though the initial damage caused by the implantation was confined in the first 200 nm. This suggests that either the newly-created quantum-light emitters or its constituents were able to diffuse deep into the epilayer. As the implantation results in an excess Al content in the epilayer, we hypothesise the low-yield creation of these new emitters to complexes of interstitial Al, as Al_i is harder to form compared to V_{Al} and V_N but

has similar [28] or larger [29] diffusivity owing to a self-interstitial knock-out mechanism. Experimental studies exist for diffusion of Mg into AlN, which is thought to occur through a similar mechanism to Al_i diffusion. Okumura et al. in a recent study [30] compared depth profiles of different implanted dopants into a MOCVD-grown AlN sample and observed that Mg dopants were preferentially concentrated at AlN:Sapphire interface after annealing at elevated temperatures of >1300 °C. This may be rooted in dislocation-mediated diffusion promotion, which was previously observed in electrochemical etching of GaN in room temperature [31]. The observed preference in the depth profile of the newly-created emitters in AlN points to an interstitial-related defect complex which is activated when in close proximity of the substrate. Further studies on the density of Al_i within the epilayer, as well as DFT investigations focusing on Al_i -related defect complexes are needed for confirmation of this model.

4. Conclusion

In summary, we investigated the effects of ion implantation and annealing on Al-implanted and unimplanted aluminium nitride epilayers on sapphire. Through iterative confocal scans between each process step, we documented the creation of distinctive emitters within an implanted aluminium nitride epilayer, in contrast to the unimplanted regions which were unaffected by annealing. The spectral signatures of the new emitters were not systematically different than the native

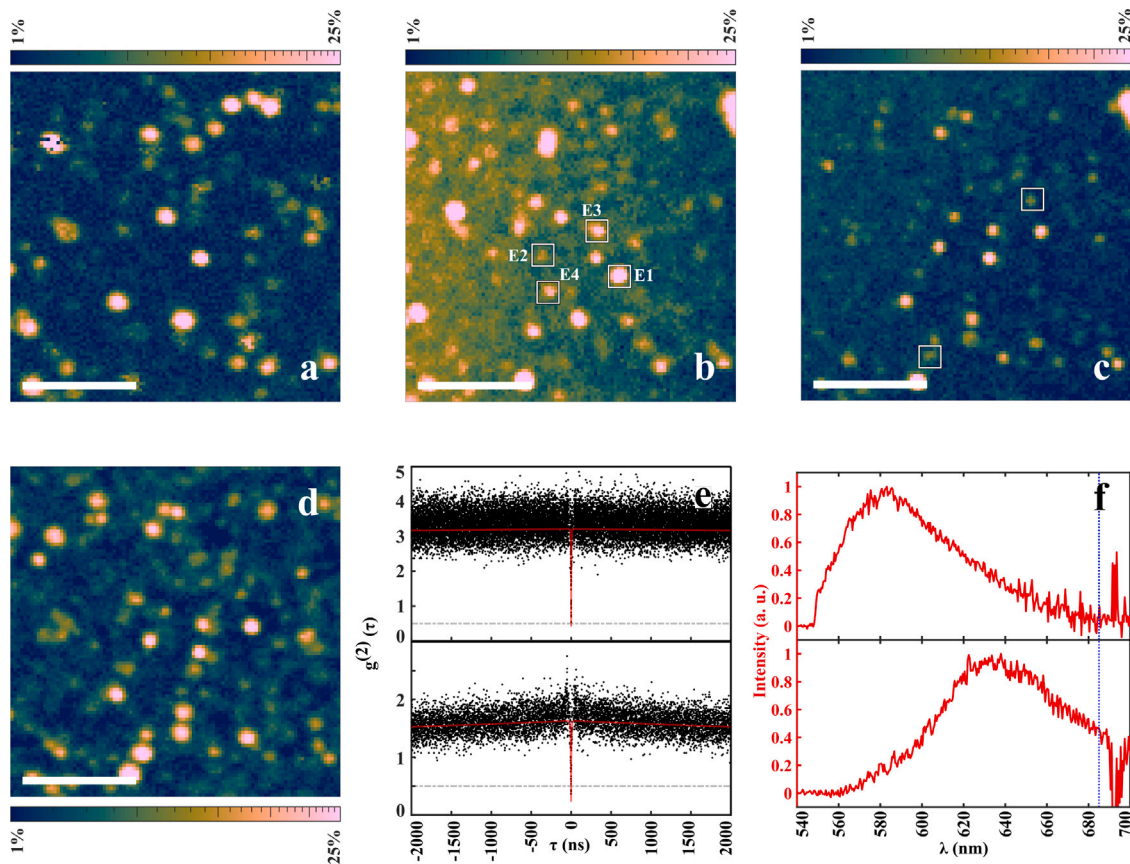


Fig. 3. Confocal scan maps of the region dedicated to the fluence of 10^{15} Al ions/cm². The scale bars are 5 μ m. (a): The region before implantation. (b): Implanted region after annealing at 600 °C for 30 min under N₂ atmosphere. Some of the newly-created emitters are highlighted. (c): Implanted region after annealing at 800 °C for 30 min under N₂ atmosphere. Some of the created emitters in (b) disappear (E1) and some new emitters appear (highlighted without label). (d): Implanted region after removing 500 nm of AlN epilayer, effectively removing the region damaged by implantation. (e) Hanbury-Brown and Twiss interferograms of two emitters created by the implantation. E1 in (b) which disappears after 800 °C anneal and E2 which does not disappear and survive the epilayer removal. (f): Spectra of E1 and E2 with the laser line removed with a 550 nm long pass filter. The sharp features align with the Cr³⁺ emission peak within the sapphire substrate. The blue dashed lines were included to represent the onset of this substrate luminescence.

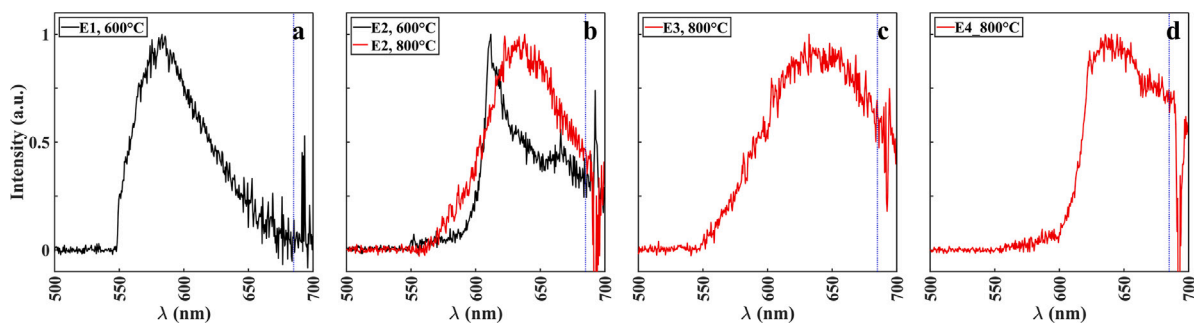


Fig. 4. Spectrograms for the newly-created emitters marked in Fig. 3. The spectra obtained following the first (second) annealing step were plotted in black (red). Blue lines represent the onset of impurity luminescence from the sapphire substrate. All spectra were taken with a 550LP filter to filter the excitation. (a) Spectra for E1 after 600 C annealing. (b) Spectrum for E2, after 600 C (black) and after 800 C (red). (c) Spectrum of E3 after 800 C annealing. (d) Spectrum of E4 at 800 C.

ones. We confirmed the created emitters displayed photon statistics consistent with single quantum emitters through Hanbury-Brown and Twiss interferometry. The created emitters displayed large and long-lived bunching contributions, consistent with previous reports. Through thinning the implanted region, we reported the created emitters migrated deep within the epilayer (>500 nm), attributed to diffusion of Al interstitials away from implantation depth. These results disclose new insight in the engineering of defect-based opto-electronic systems in a promising platform such as aluminium nitride, with appealing application in quantum communication and computing.

CRediT authorship contribution statement

H.B. Yağcı: Writing – original draft, Methodology, Investigation, Data curation, Conceptualization. **E. Nieto Hernández:** Writing – review & editing, Resources, Methodology. **J.K. Cannon:** Writing – review & editing, Methodology. **S.G. Bishop:** Writing – review & editing, Software, Methodology, Conceptualization. **E. Corte:** Writing – review & editing, Conceptualization. **J.P. Hadden:** Writing – review & editing, Supervision, Software, Methodology, Conceptualization. **P. Olivero:** Writing – review & editing, Funding acquisition, Conceptualization. **J.**

Forneris: Writing – review & editing, Funding acquisition, Conceptualization. **A.J. Bennett:** Writing – review & editing, Supervision, Project administration, Funding acquisition, Conceptualization.

Declaration of competing interest

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

Data availability

The data that support the findings of this study are openly available in the Cardiff University Research Portal at <http://doi.org/10.17035/d.2024.0324890985>.

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Appendix A. Supplementary data

Supplementary material related to this article can be found online at <https://doi.org/10.1016/j.optmat.2024.115967>.

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