

CARDIFF UNIVERSITY SCHOOL OF ENGINEERING

FLASHBACK ANALYSIS AND AVOIDANCE IN SWIRL BURNERS

By FARES AMER HATEM

Supervisors: Prof Philip J Bowen Dr Agustin Valera-Medina

A Thesis submitted to Cardiff University For the Degree of Doctor of Philosophy In Mechanical Engineering, 2017

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ABSTRACT

his study is aimed to investigate and demonstrate the feasibility and validity of various flame flashback resistance techniques for premixed fuel conditions. It presents a series of experiments to determine the impact of different configurations on flame flashback mechanisms. Experiments were performed using a 150 kW tangential swirl burner working on premixed mode with various swirl numbers; the flow field characteristics were measured by 1D LDA.

The first part of the project targeted the effect of central fuel injector geometries on flame flashback mechanisms, especially combustion induced vortex breakdown (CIVB). It was found that changing the central fuel injector outside diameter can significantly alter the flame flashback mechanism. Large injector diameters result in boundary layer flashback (BLF), contrary, the use of small injectors diameter led to CIVB. Thus a dimensionless number (χ) which represent the ratio between the injector outside diameter and the nozzle inside diameter was introduced. Using this dimensionless number the critical value of transition from CIVB to BLF has been defined, the value being χ = 0.280 for Sg=1.12 and χ = 0.320 for Sg= 0.9.

The second part was about the effect of using axial air injection instead of central fuel injectors. It was found that axial air jets have a considerable potential for flame stability requirements, they producing wider stability operation than that of central injectors. Moreover, the stability limits increase regarding both equivalence ratio and inlet tangential velocity. It appeared that using such air jets could reduce the combustor maintenance cost that arises due to a continuous harsh environment.

However, it was found that axial air jets could enforce flame propagation during flashback via wall boundary layer. Thus, the third part of the study was about the validity of using micromeshes to improve BLF resistance in addition to axial air injection. It was found that using both techniques produced high flashback resistance for both mechanisms, i.e. CIVB and BLF.

Publications

- [1] F. A. Hatem, N. Syred, A. Valera-Medina, R. Marsh, and P. J. Bowen, "Experimental Investigation of the Effects of Fuel Diffusive Injectors on Premixed Swirling Flames," presented at the 53rd AIAA Aerospace Sciences Meeting, Kissimmee, Florida, USA, 2015.
- [2] F. A. Hatem, A. S. Alsaegh, N. Syred, A. Valera-Medina, and R. Marsh, "Experimental investigation of the Effect of Air Diffusive injection on premixing swirl flames," in *55th AIAA Aerospace Sciences Meeting*, ed: American Institute of Aeronautics and Astronautics, 2017.
- [3] F. A. Hatem, A. S. Alsaegh, A. Valera-Medina, and R. Marsh, "Experimental and numerical investigation of the effect of diffusive air injection on turbulence generation and flashback propensity in swirl combustors " presented at the The 8th European Combustion Meeting (ECM 2017), Dubrovnik, Croatia, 2017.
- [4] F. A. Hatem, A. S. Alsaegh, M. Al-Faham, and A. Valera-Medina, "Enhancement flame flashback resistance against CIVB and BLF in swirl burners," presented at the 9th International Conference on Applied Energy, ICAE2017, Cardiff, UK, 2017
- [5] F. A. Hatem, A. S. Alsaegh, A. Valera-Medina, N. Syred, and C. T. Chong, "Investigation of the Effect of Air Diffusive Injection on Premixed Swirling Flames," *Journal of Propulsion and Power*, 2017.
- [6] M. M. Hossain, G. Lu, F. A. Hatem, A. Valera-Medina, R. Marsh, and Y. Yan, "Temperature Measurement of Gas Turbine Swirling Flames using Tomographic Imaging Techniques," presented at the International Conference on Imaging Systems and Techniques (IST), Macau, China, 2015.
- [7] M. Al-Fahham, F. A. Hatem, A. S. Alsaegh, A. V. Medina, S. Bigot, and R. Marsh, "Experimental study to enhance resistancefor boundary layer flashback in swirl burners using microsurfaces.," in *Proceedings of ASME Turbo Expo 2017: Turbomachinery Technical Conference and Expansion, GT2017*, Charlotte, NC, USA, 2017.
- [8] M. Al-Fahham, F. A. Hatem, A. Valera-Medina, and S. Bigot, "Investigation of Boundary Layer Flashback enhancement in Swirl Burners Using Woven Wire Steel Mesh," presented at the The 8th European Combustion Meeting (ECM 2017), Dubrovnik, Croatia, 2017.
- [9] A. S. Alsaegh, F. A. Hatem, A. Valera-Medina, and R. Marsh, "CFD Simulation and Validation of Hydrodynamic Instabilities Onset in Swirl Combustors " presented at the The 8th European Combustion Meeting (ECM 2017), Dubrovnik, Croatia, 2017.
- [10] M. Al-Fahham, F. A. Hatem, and Z. Al-Dulami, "Experimental study to enhance swirl burner against boundary layer flashback," presented at the 9th International Conference on Applied Energy, ICAE2017, Cardiff, UK, 2017.
- [11] A. S. Alsaegh, F. A. Hatem, and A. Valera-Medina, "Visualisation of Turbulent Flow in Swirl Burner under the Effects of Axial Air Jet," presented at the 9th International Conference on Applied Energy, ICAE2017, Cardiff, UK, 2017.

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Nomenclature

Alphabetic Symbols

Ao	The nozzle exit area	[m ²]
At	Total area of tangential inlets	[m ²]
Cquench	quench parameter	-
D	Exhaust diameter of swirl burner	[m]
d	characteristic length (recirculation zone length)	[m]
Da	Damköhler number	-
df	The fringe spacing	[m]
Dinj	Central fuel injector outside diameter	[m]
dq	quenching distance	[m]
dx, dy , dz	LDA control volume dimensions	[m]
f	Oscillation frequency (vortex shading)	[1/s]
FARact	Actual Fuel Air Ratio	-
FAR _{stoich}	Stoichiometry Fuel Air Ratio	-
f_d	Doppler frequency	[1/s]
f_i	The incident frequency	[1/s]
f_o	The observed frequency	[1/s]
f_s	Shift frequency	[1/s]
g_c	The critical velocity gradient (wall or parallel plate)	[1/s]
g f	The flow velocity gradient	[1/s]
g_t	The critical velocity gradient (pipe)	[1/s]
G _x	Axial flux of axial momentum	
G_{θ}	Axial flux of tangential momentum	
K	Flame stretch rate	[1/s]
Kext	Extinction strain rate	[1/s]
ke	turbulent kinetic energy	(m^2/s^2)
Le	Lewis number	-
Р	The pressure	[bar]
Pe	Peclet number	-
Qta	Tangential mass flow rate	$[m^{3}/s]$
Qto	Total mass flow rate	$[m^3/s]$
r	Nozzle radius	m
r	The radial coordinate	[m]
R_0	Nozzle exit radius	[m]
r _{eff}	Effective swirl radius	[m]
S	Swirl number	-
Scr	Critical swirl number	-
Sf	Flame speed	[m/s]
Sg	Geometric swirl number	-
Sg, comb	Combustion swirl number	-
S _{g, iso}	Swirl number at isothermal conditions	-
SL	Laminar flame speed	[m/s]
Sr	Strouhal number	-
ST	Turbulent flame speed	[m/s]
ti	the ignition time delay	[s]

Tinlet	the inlet temperature of the fluid	[k]
Tm	the initial mixture temperature	[K]
T _{outlet}	the outlet temperature of the fluid	[k]
Tu	Turbulent intensity	%
u	Swirl axial velocity	[m/s]
u'	Swirl axial velocity fluctuation	[m/s]
ū	The mean axial velocity	[m/s]
Ub	Average bulk exit axial velocity	[m/s]
U _{ref}	Characteristic velocity scale	[m/s]
u _{rms}	root mean square of axial velocity measured by LDA	[m/s]
V	Particle velocity	[m/s]
Vs	Shift speed of fringes in measurement volume	[m/s]
W	Swirl tangential velocity	m/s
w'	Swirl tangential velocity fluctuation	m/s
Win	Inlet tangential velocity	[m/s]
X	Air injector position	[m]
<i>x</i> ₀	the reference point of the divergent flow or the	[m]
	stagnation point at the tip of the recirculation bubble	
x_f	Leading edge position	[m]
у	The wall - normal coordinate	[m]
Y	LDA measurement level downstream nozzle	[m]

Greek symbols

α	Thermal diffusivity	$[m^{2}/s]$
α	Angle of orientation of the observer	deg
β	Angle of particle movement direction	deg
δ _R	Thickness of the reaction zone	[m]
θ	The angle between two incident beams	deg
λ	Air excess ratio	-
λ	Wavelength	[m]
λ_{lo}	The unstarched laminar flame thickness	[m]
ρ	density	$[kg/m^3]$
τc	The required time for chemical reaction	[s]
τυ	The period that reactants remain in the reaction zone	[s]
υ	the kinematic velocity of the combustible mixture	[m2/s]
Φ	Equivalence ratio	-
$\Phi_{ m FB}$	Equivalence ratio close to flashback conditions	-
Φ_{stable}	Equivalence ratio of stable operation	-
χ	The ratio of the outside injector diameter to the inside	-
	nozzle diameter	

List of abbreviations

AL3O2	Aluminium oxide
BLF	Boundary Layer Flashback
C ₃ H ₈	Propane
CCS	Carbon Capture and Storage
CFCS	Chlorofluorocarbons
CFD	Computational fluid dynamics
CH ₄	Methane
CIVB	Combustion Induced Vortex Breakdown
СО	Carbon monoxide
CO ₂	Carbon dioxide
CRZ	Central Recirculation Zone
CU-CI	Thermochemical water splitting with Copper-Chlorine
DME	Dimethyl Ether
EUETS	European Union Emission Trading Scheme
FBC	Fluidized-Bed Combustion Systems
H2	Hydrogen
HFCS	Hydrofluorocarbons
HHV	Higher Heating Value
HMFR	High Momentum Flow Region
HOQ	Head-on quenching
IRZ	Internal Recirculation Zones
ISL	Inner Shear Layer
LDA	Laser Doppler Anemometry
LDV	Laser Doppler Velocimetry
LHV	Lower Heating Value
LNG	Liquid natural gas
LPM	Litre per minute
LPP	Lean Premixed Prevaporised
NH3	Ammonia
NO _x	Nitrogen oxides
OEMs	Original Equipment Manufacturer
OH-PLIF	Hydroxyl Planar Laser Induced Fluorescence
ORZS	Outer Recirculation Zones
OSL	Outer Shear Layer
PMT	Photomultiplier Tubes
ppm	parts per million
PVC	Precessing Vortex Core
SF6	Sulphur hexafluoride
SL	Shear Layer
SO _x	Sulphur Oxides
SWQ	Side-wall quenching
UHC	unburned hydrocarbons
UNFCCC	United Nations Framework Convention on Climate Change
VB	Vortex breakdown



Introduction

Chapter1

All ideas in science were born in the dramatic conflict between reality and our attempt to understand it.

Albert Einstein

1.1 Climate Change and Global Warming

The demand for power production has increased significantly since the beginning of the industrial revolution to meet requirements of human life in such way that energy becomes the actual measurement of development in any society. In addition to that, the vast majority of technological breakthroughs partially or totally depend on the development of energy production techniques.

However, excessive consumption of energy has brought a huge issue represented by Global warming. Consequently, unpredicted climate change can play an important role in creating disasters such as tsunamis or floodings that always have the potential risk of producing epidemics. Moreover, these also have a direct effect on agriculture that can cause serious starvation problems.

Greenhouse gases emitted due to the burning of fossil fuels are the dominance factor that causes an increase in the global temperature. These gases absorb the reflected solar heat from the Earth surface and re-emitted it in all directions, causing considerable temperature rise of the Earth, figure 1.1 [1]



Figure 1-1 Global warming [1].

Recently, many studies have warned about potential dangers of this serious threat. There are many reports about the unprecedented rise in temperature over land, over oceans and the sea surface. Other important observations are presented by sea level rising, humidity and ocean heat content. Those ecosystem changes are accompanied by a reduction in snow cover, glaciers and sea ice. Figure 1.2 illustrates all these important climate changes [2].

Greenhouses are gases that can absorb and emit longwave (infrared) radiation in the atmosphere. This process is based on the gases' molecular nature, and consequently, the process of energy emission is multidirectional. Thus, about 50% of the longwave emission is re-radiated back to the Earth and turned back into heat energy [3]. Main types of these gases are either naturally found in the atmosphere such as carbon dioxide (CO_2), methane (CH_4), nitrous oxides (NO_X) and water vapour, or synthetic gases such as chlorofluorocarbons (CFCs), Hydrofluorocarbons (HFCS) as well as Sulphur Hexafluoride (SF_6) can also be found.

Carbon dioxide has special importance amongst other greenhouse gases because it represents their highest ratio and it is also one gas that is highly emitted by humans. For example, it accounted for 81% of all US greenhouse gases in 2014 [4]. The burning of fossil fuels has contributed to a significant increase in CO₂ levels. About 65% is

emitted to the atmosphere as a result of their combustion [5]. According to recent studies, the level of CO_2 will continue to increase, as shown in figure 1.3 [6].



Figure 1-2 .Climate indicators [2].



Figure 1-3 Recent Global Monthly Mean CO₂ [6].

Another greenhouse gas which can also act negatively on the climate is Nitrous oxides (NO_X). It represents about 6% of climate forcing, and it also results from the burning of fossil fuels. Similar to CO_2 it is increasing annually as can be seen from figure 1.4 [2].



Figure 1-4 the rise in the atmospheric nitrous oxide [2].

1.2 Fuel Flexibility

Since the beginning of the industrial revolution in the 18th-century fossil fuels such as coal, crude oil and natural gas have represented the dominant energy source. Steam locomotives used coal as fuel, later the discovery of oil and natural gas contributed significantly to the development of the energy sector. Those fuels have the main share amongst other fuel types for power production. Thus, world demand for those fuels is expected to increase for the next few decades as shown in figure 1.5 [7].



Figure 1-5 World energy consumption by fuel [7].

Therefore, burning of these fuels has been considered a serious threat to the climate by emitting a considerable amount of pollutants [8, 9]. Moreover, fossil sources have gone a rapid depletion leading to limitation of resources [10-12]. Thus, there has been a strong trend for investment in ecologically friendly alternatives named renewable energy resources such as the wind, tidal, hydro, biomass, solar and geothermal energy. Nevertheless, since the 2nd half of the twentieth century, those sources are judged by many limitations like their variation in time and location [13]. Also, cost and low carbon emissions are important features when differentiating between energy types, as can be seen from figure 1.6 [14]. Hence, it is important to rely on some alternative fuels while keeping developing research on renewable energy. More efforts are required in order to have flexible fuels that can meet clean energy requirements, in addition, to cover the predicted shortage in traditional fuels.



Figure 1-6 comparisons between fossil and renewable energy sources (production cost and CO₂ emissions) [14].

Introducing new fuel types must satisfy many regulations such as sustainable development via renewable production. Presently, an example is biofuels that have demonstrated a potential to be a renewable fuel source in addition to their flexibility on engine operation. Nevertheless, the reliable and safe transition from using fossil fuels to other sustainable sources is still a major topic of research [15]. In general, many features can address the fuel flexibility for power generation;

1.2.1 Sustainability

Long-term sustainability represents the availability of enough feedstock to allow future generations to use it without affecting the ecosystems or climate through cost effective implementation.

1.2.2 Global Climate Impacts

The significant increase in the awareness about the high level of pollution and how to

minimise carbon footprint, and the regulations of the international agreements have urgent priority.

1.2.3 Fuel Quality and Operability Boundaries

Fuel properties have the direct impact on the working efficiency. Hence the type of combustion system required specific designs, as can be seen in figure 1.7[16]. In addition, fuel operability limits are crucial from an operational stability requirements point of view.



Figure 1-7 GE Gas turbines for syngas applications [16].

1.2.4 Work and Safety Requirements Environment

Type of duty and working environment play an important role in choosing the fuel type since some fuels cannot work properly at high altitudes (aviation conditions) because some properties like density or flow characteristics may change due to considerable temperature reduction [17].

1.2.5 Fuel Prices and Cost Limitations

Competition inside energy markets is one of the main players when differentiating between fuel types.

1.2.6 The Ability of Integration with other Fuel

Recently, hybrid combustion regimes became one of the possible techniques to achieve environmentally friendly regulations with a significant increment in overall system efficiency. However, such systems require fuels that have the ability to work with different fuel types inside the same plant cycle [18].

These are only some of the vast numbers of issues raised when fuel is selected for power generation or production.

1.3 Alternative Fuels

Recently, alternative fuels have been introduced to many combustion systems in order to allow consumption reduction of fossil fuels as well as to meet clean energy requirements [11, 19, 20]. Alternative fuels such as biofuels are multifunctional energy carers and well geographically distributed resources [21]. Historically the needs for using new fuel resources have changed continuously based on availability, sustainability, economic feasibility, progress in technology and presently, climate concerns. In early human ages, the society was entirely dependent on wood as a source of heat. Nevertheless, excessive use of wood as an energy resource in Europe, for example, caused dramatic deforestation in the early medieval period before starting to use coal as an alternative fuel in the sixteenth century [22].

Coal mining acquired particular importance during the medieval centuries; importance has significantly risen with the beginning of the industrial revolution which was mainly based on coal availability to power steam engines. Coal became the heartbeat of this technological breakthrough. Even though the discovery of oil in the 19th century followed enormous development in power sector technology, coal still has a special importance as a desirable fossil fuel for energy production and still represents a sizable share of global energy and industrial applications as can be seen in figure 1.8 [23]. There has been a strong trend to find more reliable and more environmentally friendly energy resources. Thus renewable energy which comes from natural resources such as sunlight, wind, tides, rain, waves and geothermal heat are an interesting alternative energy source. However, this contribution in the overall global energy sector is still low and comprise about 17% of the total world energy consumption, as can be seen in figure 1.9 [24].



Figure 1-8. Global primary energy consumption share by fuel type 2012 [23].

Nevertheless, increasing this share seems to be difficult and faces a lot of obstacles represented by their intermittent nature, high cost, and availability of suitable technology as well as the limitation of those resources. Thus the importance of using alternative fuels instead of traditional fossil fuels has increased significantly in the last decade to reduce high pollutant levels and to lower energy cost [25].



Figure 1-9 Renewable energy share in the world energy sector 2010 [24].

Some common alternative fuels for power generation are as follows:

1.3.1 Biofuels

Fuel made from renewable biological materials sources such as plants and animals. It is derived from organic matter directly obtained from plants or indirectly from agricultural, commercial, domestic and industrial wastes. Biofuels have the potential to be significantly less expensive than other fuels. Liquid and gaseous fuels from biomass can achieve low emission requirements, consequently, could be used as potential fuels [26]. Common biofuels are biodiesel, ethanol, biogas, syngas, solid biofuels, vegetable oil, bio-ethers, biofuel gasoline, green diesel and other bio-alcohols [27].

1.3.2 Hydrogen

It is considered as the fuel of the future. It can be produced from a number of sources such as water, hydrocarbon fuels, biomass, hydrogen sulphide, and chemical elements with hydrogen. Because hydrogen is not available as a separate element, it is required to be separated from the sources mentioned above. The process of separation requires energy to do the disassociation, like water decomposition into O_2 and H_2 by passing a direct current which drives electrochemical reactions. Hydrogen can be blended with natural gas and other hydrocarbons in order to have less carbon dioxide, carbon monoxide and unburned hydrocarbons emissions [28].

1.3.3 Liquefied Natural Gas (LNG)

It consists of natural gas (mainly methane, CH₄) and it is produced by liquefying natural gas taken from gas fields after removing impurities. This process leads to condensation of the gas into liquid its volume 1/600th that of the natural gas. Consequently it has much more energy density than receiving natural gas [29]. Although it comes from fossil fuels, an increase of energy density has a direct impact on cost.

1.3.4 Ammonia NH₃

Considered as a safe hydrogen holder, it has a high hydrogen content per unit volume, which makes it a promising green energy storage. Furthermore, it can be produced commercially by using synthesis methods. Its most advantageous feature is its small CO₂ production and high octane number [30].

1.3.5 P-Series

It is a family of renewable, non-petroleum liquid fuels that can be used instead of gasoline. They are mixtures of natural gas, methyl tetrahydrofuran pentanes-plus and butane. They are high octane number alternative fuels used in flexible fuel plants [31].

1.3.6 Dimethyl Ether (DME)

Also known as methoxymethane, it can be produced from natural gas, coal or biomass by converting the syngas to methanol followed by methanol dehydration to dimethyl ether. It is considered as a future energy option. It has been demonstrated that produces lower NO_x and SO_x emissions than conventional diesel. Furthermore, it has demonstrated an increase in efficiency with decreased in NO_x and CO emissions compared to methane when used in gas turbines [32].

Although the above mentioned alternative fuels seem to be promising for clean energy requirements, switching gas turbines systems to be able to work on different fuels instead of being limited to natural gas is challenging task. Many parameters must be taken into consideration in ordered to make gas turbines capable of fuel-flexi operation, consequently burning different fuel types safely and efficiently. For example, the compressor is designed to work on certain margin of mass flow rate, thus upon introducing different fuel means different density, hence the compressor required to be able to work on various level of mass flow rate, higher or lower according to the degree of dilution [33].

Moreover, controlling and balancing the amount of pollutants is also a difficult task, high hydrogen blends, from one hand, can significantly reduce carbon-containing pollutants. On the other hand, NO_x emissions might be increased as result of working at high temperature [34], hence controlling the combustion temperature is required. Furthermore, most of the combustors can be subjected to different operability issues; those issues are strongly affected by fuel properties [35]. The most critical combustion issues are blowoff (when the flame detached from its stable operation and physically blown out), flashback (the upstream flame propagation from the combustion zone to the premixing zone inside the combustor), combustion instability, and autoignition. These undesirable operational issues are so complicated, especially in swirl combustors.

Flame flashback has particular importance because it can cause considerable damages to the system. Thus when introducing alternative fuel to gas turbines instead of natural gas, many factors and operational parameters must be considered, some alternative fuels have a considerably high flame speed such as syngas or high hydrogen blends, thus they are more likely to initiate earlier flashback than natural gas, consequently the operation stability map when using those alternative fuels is different than that of natural gas. Therefore to achieve safe fuel switching process, a reasonable margin of overlapping between two stable operational regions is crucial.

In conclusion, despite the high importance of using alternative fuels to operate gas turbines to achieve low harmful emissions, many unpredicted issues can arise. Thus many modifications and enhancements of gas turbines still in demand. One of the crucial and urgent obstacle that facing such developments is the flame flashback. For this reason, more additional studies and investigations in this area are still highly valuable.

1.4 Power Generation Gas Turbines

Large scale power generation plants are prepared and developed to produce power, especially electricity for public consumption. The efficiency of operation of such plants depends mainly on the integration of the existing technology and infrastructures.

Gas turbines played an important role in power generation sector since the second half of the previous century and still considered as a very important part of the energy market regarding large-scale power generation. The gas turbine is an internal combustion engine in which the expanding gases emerging from one or more combustion chambers drive a turbine. A rotary compressor driven by the turbine compresses the air which is used for combustion. The power is taken either as torque from the turbine or thrust from the expanding gases, figure 1.10.



Figure 1-10. Gas turbine schematic [36].

In a gas turbine, the air is drawn into the intake by a compressor which is attached to the main shaft. Air pressure and temperature will increase as air passes through the compressor stage. After that, the air leaves the compressor and passes through the diffuser which converts the kinetic energy to pressure. Then the air passes into the combustor. Each combustor can be equipped with swirl burners where the fuel pressurised by fuel pumps is then mixed with incoming primary air to form a fine combustible mixture. Combustible mixtures are then ignited.

Upon completion of the igniting cycle, the igniters are shut off, leaving continuous combustion at constant pressure. Inside the combustion chamber, the gases rise their temperature and increase in volume. After leaving the combustors, the hot gases enter the nozzle guide vanes which direct them at increased velocity into the turbine blades causing the turbine discs to rotate. The output power of the turbine powers the compressor and any remaining power is used to drive a shaft or from a high-speed jet.

1.4.1 Cycles and Materials

Thermodynamically gas turbine engines work according to the Brayton cycle. Efficiency is function of pressure ratio over this cycle, the higher pressure ratio, the higher overall system efficiency. Therefore, to achieve high operational efficiency both compressor and combustor should be designed to work at very high temperatures and pressures.

1.4.1.1 The Compressor

The compressor is responsible for raising the inlet air pressure to its maximum possible value before feeding it to the combustor. This ratio of air consumption is about 68 times higher than that in the spark engine [37]. Combustor and turbine blades also have to be able to withstand the exceptional operation conditions. For this reason, the designing and manufacturing of gas turbine components mostly involve advanced processing techniques. Thus with a higher cost of materials used in the final product cost increases dramatically. Figure 1.11 [38] shows a gas turbine compressor. It is obvious how advanced manufacturing technology is important in its production.



Figure 1-11. compressor rotor [38].

1.4.1.2 Gas Turbine Combustors

Even though there are a great variety of gas turbine engines, they have the same parts in common. One of these parts is the combustor or combustion chamber. The main function of the combustor is to burn combustible mixtures of fuel and pressurised air to generate maximum heat release at constant pressure. The generated hot and accelerated stream of flow gases is delivered uniformly to the turbine blades. Hence the combustor (the most heat loaded component) is responsible for performing the most difficult part of the combustion process.
Therefore any combustor in any gas turbine has to achieve many requirements to be able to do its tasks properly, [39]

- High-combustion efficiency, which can be achieved by complete combustion of the fuel which leads to total conversion of chemical energy into heat.
- Reliable and smooth ignition in different operation environments (for example stationary gas turbines or those used in aviation tasks).
- Operating at wider equivalence ratios.
- Minimum pressure loss.
- An outlet temperature distribution that is tailored to maximize the lives of turbine blades and nozzle guide vanes.
- Low pollutant emissions.
- Easy maintenance and low-cost manufacturing, and
- Have good fuel flexibility (can work with different types of fuel).

1.4.1.3 Types of Combustors

Design and manufacturing techniques of gas turbine components have been changed significantly over the years to meet variable demands and regulations such as their duties, working efficiency and recently environmentally friendly regulations. Introducing advanced engineering programmes like computational fluid dynamics (CFD) and laser diagnostics contribute effectively in facilitating the design process. Common combustor types are [40]:

- Can type combustor. It is a self-contained combustor that has its fuel injector, igniter, liner and casing. Multiple cans can be arranged peripherally around the central axis, and the resultant exhaust of all of them is used to feed the turbine. The advantage of using can-type combustors is their easy maintenance. However, their use has some disadvantages, such as the pressure drop across one can combustor that can be higher than in other combustors. Also because they weigh more than other types they are not the preferred option for aircraft turbines.
- Can-annular. This combustor also has separated combustion zones contained in separate liners with their fuel injectors. It is different from annular combustors by having a common ring casing. Liner holes or tubes

enable sharing between combustion zones. Consequently, air is allowed to move circumferentially. This type has a more uniform temperature, which is ideal for turbine blades. It does not need separate igniters for each can, has less pressure drop than the can type and it is lighter. Nevertheless, it is more difficult to maintain than the can type.

• Fully Annular combustor. It is the most used type; it has a continuous liner and casing in a ring. Its advantages are more uniformly in combustion, lighter in weight, less surface area, uniform exit temperatures and less pressure drop than the other types. Figure 1.12 shows the three types [41].



Figure 1-12. different combustor configurations [41].

1.4.1.4 Combustor Features

Despite tremendous technological development in designing and manufacturing techniques of gas turbines, combustors have maintained some of its basic features and characteristics. Those features are as follows:

• Primary zone. Anchors the flame and provides adequate time to achieve complete combustion. This can be done by a generation of toroidal flow, which can produce perfect mixing and recirculation of hot combustion gases,

consequently having continuous ignition. Most gas turbine combustors use swirl burners to achieve this target.

• Secondary zone (Intermediate zone). The role of this zone is to accelerate the passage of incomplete combustion products (CO and H₂) to the dilution zone in order to complete combustion by injecting formidable amounts of air. Moreover, the injection will achieve the temperature drop that will help in preparing the gases before their entrance into the turbine [39].

In other words, the function of this zone is to oxidise CO into CO_2 , which can be done by accelerating the dull forward reaction through creating a generic lean equivalence ratio, keep the temperature at an elevated level, providing the residence time required for complete oxidisation [40].

 Dilution Zone. Reduces the temperature of the combustion products and improve mixing of the resultant exhaust gases to achieve acceptable temperature distribution that maintains long operating life for turbine blades. Figure 1.13 illustrates a combustor feature[40].



Figure 1-13. Combustor features (zones)[40].

1.4.2 Performance and Efficiency of Gas Turbines

Gas turbines have played the dominant role in power generation and propulsion for several decades, thanks to their ability to work using different fuels. They have been extensively employed in various applications starting from powering aircrafts, producing electricity, driving vehicles, trains and even for military applications. However, performance and overall working efficiency seem to be crucial factors that decide the suitability of each type of turbine. For example for aviation applications performance requirements should take into consideration the operation conditions at high altitudes such as a considerable drop in ambient temperature and even aerodynamic stresses where inlet distortions due to variation in the angle of attack can significantly affect the final performance [42].

Combustion efficiency has always been an important target for designers and manufacturers because it is the important feature that addresses the system's performance. Obtaining high operation efficiency mostly depends on the combustion processes. Nevertheless, achievement of complete combustion without unburnt products or low level of pollutants is still a challenging task [39].

Generally the factors that affect combustion efficiency hence the performance are [39],

- Air ambient inlet temperature and humidity.
- Ambient pressure.
- Inlet and outlet losses.
- Fuel type.
- Change in blade surface.
- Hot corrosion of the material.

1.5 Global Awareness and International Efforts

Even though energy which is mostly produced by burning fossil fuels is the lifeblood to sustain modern world life, many scientists and active climate protection organisations have warned against the potential dangers of excessive greenhouse emissions on the climate, consequently human's life on this planet. In 1827 the French scientist Jean-Baptiste Fourier was the pioneer in warning about greenhouse gas effects while the Swedish chemist Svante Arrhenius was the first who calculated the increase in concentrations of carbon dioxide and how it affects the global temperature [43].

Scientists' efforts have continued to address this issue. In 1957 Revelle and Suess [44] argued that the CO_2 concentration in the atmosphere have increased by 10% due to fuel combustion. In 1961 Charles Keeling identified a data curve which became

known later as the Keeling curve. This curve shows that carbon dioxide concentration in the atmosphere is rising steadily year by year [6].

Researches and reports have drawn the political attention on the problem, and it became globally famous by the 1990s. Many international efforts have been made to establish restrictions and control greenhouse emissions. The global efforts had a breakthrough by setting up the Kyoto protocol in 1997. This protocol extended the 1992 United Nations Framework Convention on Climate Change (UNFCCC) and admitted that global warming existed and was caused by human-made CO emissions. The main target of the protocol aims to control the emission of main greenhouse gases. The average target was cut about 5% relative to 1990 levels by 2012. However, this target has not been entirely achieved due to lack of commitment of many countries [45].

The importance of greenhouse gas emission threats has urged the European countries to establish the European Union Emission Trading Scheme (EUETS) which is the first and largest greenhouse gas emission trading scheme in the world. It was launched in order to combat climate change. The implementation of EUETS policies resulted in a reduction of greenhouse gas emission from big emitters about 8% based on 2005 levels [46]. Nowadays the issue is being addressed and became globally known, and it is predicted that the awareness about its potential dangers will increase considerably, especially when thinking about recent disasters that happened as a result of climate change. However, limited renewable energy resources, dependency on fossil fuels, and economic constraint are factors predicted to contribute to have the same emissions level at least for, the next few decades.

Thus in order to meet clean energy requirements as well as Kyoto and other organisations regulations it is highly important to develop combustion systems to achieve the maximum possible efficiency in addition to as much as low possible carbon and other greenhouse gas emissions reduction.

1.6 Summary

Gas turbines are used in a variety of applications, mostly in power plants and the aviation sector. Nevertheless, because of their huge consumption of fossil fuels, they

produce a significant amount of harmful emissions that cause the atmospheric temperature to rise due to greenhouse effects. Thus investigating their operation issues to find suitable solutions can enhance their performance and increase efficiency while lowering greenhouses gas emission.

1.7 Objectives

This thesis investigates swirling flame instabilities, especially flame flashback mechanisms that inherently occurs in swirl burners which are extensively used in gas turbine combustors. The potential effects of flashback on system hardware and consequently the operation stability and working efficiency are strong motivations to continue investigating this phenomenon and how to determine suitable treatments to tackle flashback when using alternative fuels. Although there have been many studies about swirling flame instabilities, the issues have not been totally diagnosed because of the complex nature of this phenomenon. The main objectives of this project can be summarised as follows;

- Determine the impact of different central fuel injector geometries on the flow field characteristics downstream the burner mouth, consequently the propagation of the combustion induced vortex breakdown flashback (CIVB).
- Investigate the effect of using axial air injection instead of central fuel injector on flashback limits and hence burner stability operation map. The effect of air injection position with respect to burner base plate will also determine. The characterization will be based on the measurements of change in velocity and turbulence downstream burner mouth.
- Investigate flashback resistance technique for both CIVB and BLF simultaneously. This includes using both air injection and micromesh surfaces for enhancement flashback resistance.

1.8 Thesis Structure

This thesis is divided into a number of chapters, which are as follow,

- Chapter 1. An introduction that reviews the problems of global warming and climate change additionally illustrates the basic principles of gas turbines and alternative fuels for power generation.
- Chapter 2. A comprehensive review of all previous work in this area, especially different flashback mechanism, as well as the combustion technologies, fuel blends and lean premixed combustion.
- Chapter 3. All experimental equipment and setups used during the test campaign.
- Chapter 4. Experiments and results of the impact of using different central fuel injector geometries with different swirl numbers on the combustion induced vortex breakdown are stated in this chapter.
- Chapter 5. This chapter describes the experiments of using central air streams instead of the central fuel injector, analysing and investigating the flow field characteristics and the nature of turbulence close to the burner mouth.
- Chapter 6. Characterization the effect of downstream velocity gradient on flame flashback are mentioned in this chapter.
- Chapter 7. This chapter outlines experimentally how to reduce wall boundary layer flashback arises due to central air injection.
- Chapter 8. Summary of discussions. This chapter highlights and discusses all the obtained results comprehensively.
- Chapter 9. Conclusions and future work. This chapter summarized the main conclusions of the study and the future work.



Swirl Combustor Characteristics and Combustion Instabilites of Gas Turbines

Chapter 2

The gift of mental power comes from God, and if we concentrate our minds on that truth, we become in turn with this great power. Nikola Tesla

2.1 Introduction

There has been an exceptional and enormous technological progress on new manufacturing techniques and performance characteristics of gas turbines since the 2nd half of the previous century. This evolution has encountered many issues, some of these related to availability and suitability of fuels, while others are linked to sophisticated designing and manufacturing of components. Interestingly, since the end of the previous century, another problem draw the attention of OEMs which is the concern of pollutant products from gas turbines and their potential effects on global warming. The former problem has enforced scientists and researchers to find out other manufacturing techniques and combustion technologies to meet these environmental issues.

Combustion processes inside gas turbine combustors are complicated due to the mutual interaction between physical features and chemical properties. Consequently, there has been extensive research in this field in order to cover the lack of knowledge in the complex area of fuel flexibility. Hence, important operational demands can be achieved with higher combustion efficiency, more uniform exit temperatures and low emissions. Lean premixed combustion has proved as a favourable technique that can fulfil some objectives such as low emissions, especially for NO_x reduction.

Furthermore, hydrogen or hydrogen blends and some syngases have been introduced successfully as promising alternative fuels in lean premixed combustion in order to meet high-performance requirements as well as less or no CO_2 emissions. However, upon these significant changes in combustion technologies, many new issues have appeared.

Combustion instabilities have been considered as an inherent drawback in gas turbine performance for swirling burners, their features and mechanisms have changed accordingly by introducing new combustion technologies and new fuel types. These undesirable problems represent special challenges for designers, producers, and operators. Gas turbine combustors have received the most effort amongst other components because it involves regions that can lead to flame flashback and blow off, combustion instabilities, significant variation of temperatures and fuel substitutability.

2.2 Combustion Systems and Combustion Technologies

Generally, combustion systems can be classified according to the technology used. Hence there are five types of combustion systems that are used in gas turbines;

2.2.1 Standard Systems

These systems are also known as stoichiometric combustion plus dilution. They are characterised by high NO_x levels but high performance. They are usually used in engines that need an effective performance like aircraft engines [47]. In these systems, part of the incoming air is used for complete combustion. The method of air injection is designed to achieve toroidal vortices, and the fuel is injected at the central axis of such vortices to obtain maximum contact, hence better mixing. Pilot flames are usually used to ignite the mixture; then the hot combustion products can fulfil continuous ignition of the nearby fuel-air mixture streams. The secondary or dilution air splits into two streams, one used to obtain uniform temperature profiles by producing sufficient mixing, while the other is used for liners cooling.

2.2.2 Catalytic Combustion Systems

This combustion technology uses catalysts in order to accelerate the desire oxidisation reaction of fuel. Consequently, a considerable reduction of pollutant emissions especially NO_x are produced by lowering the reaction temperature. The catalyst process depends mainly on inlet temperature and the type of mixture. The temperature should be equal to the adiabatic flame temperature for lean hydrocarbon mixtures such as CH₄, C₃H₈ since higher temperature can change the materials. In the case of fuels containing hydrogen, it is possible for catalytic temperatures to exceed the adiabatic temperature [48]. In general, the temperature range should be between 725 K and 1050 K whereas at higher

temperature the catalyst may yield decomposition [47]. The required inlet temperature is usually obtained through air preheating. One of the favourable features of this system is pressure loss around 2.5%, which is lower than that for standard combustors of 5% [49]. Moreover, the overall NO_x levels can be effectively reduced up to 5ppm [50]. Additionally employing hydrogen in this technology can increase combustion efficiency, and flame flashback can be controlled by lowering the combustion temperature [51].

2.2.3 Lean Premixed Combustion Systems

Diffusive flame combustors have been conventionally used in gas turbine engines for power generation based on their reliable performance and acceptable stability characteristics. However, the most disadvantageous feature of these types of combustors is their high and unacceptable NO_x thermal levels [52]. The global concerns about the unprecedented rise of pollutant emission levels have drastically led both designers and manufacturers to switch to more environmentally friendly combustion technologies.

Lean premixed combustion has become the most promising technology for emission reduction in gas turbine combustion systems [53-55], with NO_x levels that reach less than 10ppm [39]. In this technology, high levels of excess air are provided to the reaction which results in Φ <1 or λ >1 [56]. Local flame temperatures and residence time have a direct effect on the mechanism of NO_x formation in lean premixed mixtures. The majority of the fuel in lean premixed combustors is burned at lower temperatures than conventional combustors [57]. Lower flame temperatures mean lower nitrogen oxidation. However, chemical reaction rates will become slower at lower flame temperatures. In other words, lower flame temperatures require more time to complete the reaction to oxidise CO into CO₂, hence, it is possible to have unburned hydrocarbons (UHC) as well as high levels of CO emissions.

Therefore, lean premixed combustor should be designed o satisfy both CO and NO_x low emission requirements. This can be done through control of flame temperatures which is a function of the working equivalence ratio, as can be seen in figure 2.1 [58].



Figure 2-1 The effect of flame temperature on emissions [58].

Lean premixed combustion technology is widely used in swirl burners in different gas turbine systems to obtain stable operation and reduce harmful emissions [11, 52, 59]. However, swirling flows are complicated systems due to the appearance of coherent structures such as central recirculation zones (CRZ) and the precessing vortex core (PVC). These structures have direct effects on the stability and operation of combustion processes due to their nature [60] and lack of damping mechanisms compared to diffusive flames [61-63]. Common combustion instabilities that can arise when using lean premixed combustion with swirl combustors are blowoff, flashback, and thermoacoustic combustion instabilities.

Operation under lean conditions may result in complete flame blowoff [64]. These systems are also a prone to flashback an inherent drawback of lean premixed combustion [65-69]. This phenomenon is more likely to increase with increasing hydrogen concentration in the fuel blends [70-73]. Additionally, the low frequency produced from swirl flow structures can be coupled with the natural frequencies of the combustion system equipment producing resonance and damages [74-77].

In order to have low emissions and satisfy stable operation requirements represented by good flashback resistance and broad lean blowoff limits, new enhancement techniques

known as lean premixed prevaporised (LPP) have been developed successfully [78, 79]. However, the technique is restricted to liquid fuels. Hence it has been implemented to investigate new biofuels that have a promising reduction in CO emission [80]. Considerable work is still required with gaseous blends.

2.2.4 Fluidized-bed Combustion Systems (FBC)

It is the combustion technology used to burn solid fuels. A granular or lumpy solid fuel particles such as coal or woody biomass are suspended in hot bubbling fluid beds of ash or other materials such as sand and limestone are burned with oxygen when a jet of air is blown through the bed. The resultant mixture of gas and solids will lead to rapid heat transfer and chemical reactions within the bed. The air required for combustion is usually injected through porous holes at a high velocity to become airborne or fluidised. Using of FBC systems in power generation has increased because of their competitive cost, more environmentally friendliness, good fuel flexibility, high reliability and high efficiency [81].

One of the important features of this technology is that fuel burns at a temperature range between 750 – 900 0 C, which is much lower than the threshold for NO_x formation temperature (1370 0 C) [40]. Also, sulphur pollutants like sulphur oxide SO_x can be captured. This can be done when the flue gases are brought into contact with sulphur absorbing chemicals like limestone. The use of Oxyfuel combustion technologies with FBC systems can lower NO_x and SO_x levels around 50% and 90%, respectively [82, 83]. Although there are many advantages that outweigh FBC systems over the conventional ones. Erosion inside the boiler, no uniform residence time of solids in the reactor due to rapid mixing in the bed, difficulty in describing the flow of gas in the case of bubbling beds of fine particles and fuel flexibility limitations are some disadvantages that stain their performance [84, 85].

2.3 Combustible Mixtures Types

The mixing process of combustible mixtures in gas turbine systems is different from system to system, depending on many factors such as the required output power, efficiency, operation stability demands, nature of duty and even environment protection regulations. There are three mixing techniques used in gas turbines.

2.3.1 Premixed Mixtures

In premixed mixtures, fuel and oxidiser are intimately mixed at a molecular level before combustion is initiated. The mechanism of combustion is based on the propagation of the flame front into the unburnt mixture. The equilibrium between a chemical reaction and the amount of heat generated at the reaction zone determines the premixed flame structure and its characteristics, flame propagation speed or the so-called burning velocity, which is one of the characteristics that can address flame stability. In some cases, the burning velocity may become so fast that causes a dramatic transition to detonation, i.e. when the combustion wave moves at supersonic velocity with the possibility of severe damages by an explosion [86].

Premixed mixtures are relevant to gas phase combustion, and it is not applicable with liquid fuel droplets or solid particles. Until now it is commercially limited to natural gas and a small fraction of other active components. Other fuels are usually burned using diffusion systems or non-premixed combustion systems with special burner designs [35, 87].

2.3.2 Partially Premixed Mixtures

Partially premixed mixtures are used to achieve some operation requirements in certain systems. In this method, part of the fuel or air is injected axially with a premixed blend, or both fuel and air entered separately and mixed partially by turbulence [88]. Partially premixed is used when it is required to have both the operation features of premixed and non-premixed mixtures. It is utilized in some applications that need to reduce pollutant emissions as well as have wider operation stability. Partial diffusion injection can affect the process positively ensuring ignition at the injection plane in addition to a reduction of pressure fluctuations [89].

2.3.3 Non-Premixed Mixture

In non-premixed or diffusive combustion systems, fuel and oxidant enter separately to the reaction zone. This mode of mixing is used in some combustion systems like internal combustion engines, turbojet afterburners, oil refinery flares and pulverised coal furnaces. One of the important features of non-premixed combustion is the uniform temperature distribution inside the furnace that allows low flame temperatures, consequently low NO_x

emissions. It can be achieved through mutual dilution between recirculated burned gases and both fuel and oxidizer. Fuel and air jets are injected at high momentum into the combustion chamber, hence allowing a considerable increase in chemical residence time [90, 91]. In other words, the jet of gaseous fuel is mixed with sufficient quantities of air. Hence all fuel can burn at a suitable distance from the nozzle [88]. High stability and flame anchorage are also features of this technique.

However, despite some positive outcomes of using non-premixed combustion, soot formation seems to be one of the undesirable features of this mixture [92, 93]. Soot formation of a particular fuel is governed by physical and operating parameters. The amount of soot formation in such systems can be reduced by using oxy-fuel combustion and adding CO_2 to the fuel or oxidizer [83, 93, 94].

In order to achieve a deep understanding of non-premixed combustion systems, several experimental and numerical studies have been developed. Some mathematical models have proven to be the best tool for these investigations. The models are based on the assumption launched by Burke and Schumann [95] that the mixture fraction is a conserved scalar, hence the global properties of non-premixed combustion such as flame length is predicted and all scalars can be related to the mixture fraction.

2.3.4 Stoichiometry Ratio

Fuel to air ratio (f/a) is a significant parameter used to characterise combustion; it is expressed either in volume or mass basis. The ratio is referred to be stoichiometric when the quantity of air is precisely enough to consume all of the fuel theoretically. The stoichiometric reaction for any arbitrary fuel can be written as follow:

$$C_m H_n + \left(m + \frac{n}{4}\right) O_2 + \left(m + \frac{n}{4}\right) \left(\frac{79}{21}\right) N_2 \to m CO_2 + \frac{n}{2} H_2 O + \left(m + \frac{n}{4}\right) \left(\frac{79}{21}\right) N_2 \qquad (2-1)$$

Thus methane - air stoichiometric combustion reaction is

$$CH_4 + 2(O_2 + 3.76N_2) \rightarrow CO_2 + 2(H_2O) + 7.52N_2$$
 (2-2)

And stoichiometric fuel to air ratio is

$$FAR_{stoicio} = \frac{Weight of the Fuel}{Weight of the Air}$$
(2-3)

And the equivalence ratio is

$$\Phi = \frac{FAR_{act}}{FAR_{stoich}} \tag{2-4}$$

Equivalence ratio refers to the state of combustion whereas,

When $(\Phi = 1)$ ideal state will be satisfied

 $(\Phi < I)$ Lean mixture conditions

 $(\Phi > 1)$ Rich mixture conditions

However, some works use air access ratio $(\lambda = \frac{1}{\phi})$ to refer to the combustion status.

2.4 Swirl Burners

Swirl burners have been widely used in many industrial and technical applications for decades; this includes gas turbine combustors, internal combustion engines, pulverised-coal power station, refineries and process burners [74, 96]. Swirling flows have the ability to control the combustion process through mixing improvement rates between the reactants, reducing emissions with an increase of output power. Wider ranges of operational stability represented by flashback and blowoff limits can be achieved accompanied by high levels of turbulent flame speed and reduction of combustor size using these flows [74, 97-100].

Historically, there has been a significant development in designing and manufacturing swirl burners in order to meet operational requirements of combustion systems such as duty nature, working environment, fuel cost, running efficiency, maintenance cost, and recently low pollutant demands, therefore swirl burners used in the aviation sector are different to those used in stationary gas turbines or marine propulsion systems. Swirl burners can be used to burn different fuel types with different calorific values. Multi-inlet cyclone combustors and swirl burner furnace systems have the ability to burn even poor quality and low caloric value fuels (1.3-1.4 MJ/m³) without any supporting fuel [101].

2.4.1 Swirling Flow Characteristics

Swirling flows are characterised by their excellent flame holding with reasonably low emissions levels. They depend on the aerodynamically formed central recirculation zone (CRZ) which recirculate heat and active chemical species to the root of the flame and create low-velocity regions so that the flame speed can match the local velocity of the unburned mixture [65, 74, 96, 102].

Vortex breakdown phenomenon represents the crucial factor in the formation of the CRZ and determine its size and strength. The phenomenon initiates due to the formation of flow stagnation points downstream of the swirl exit when swirling flow yield a sudden expansion. Hence, there will be a considerable decay of intensity in the axial direction leading to negative pressure gradients. Consequently, the flow will stagnate at a certain point downstream resulting in the formation of a reverse flow region. This region is characterised by rapid flow changes in the axial direction accompanied by a reduction in the velocity of the approaching flow [74, 103]. Figure 2.2 illustrates the mechanism of CRZ formation [74].



Figure 2-2. Typical recirculation zone setup in the swirl burner [74].

The tangential velocity downstream burner exit follows a Rankine distribution, with a steady increase from the central axis to a peak in the forward shear layer (forced vortex) followed by a decay towards the walls (free vortex). This type of distribution continues downstream with steady decay of velocity levels, figure 2.3 [74].

Aerodynamically, swirling flows yield both a vortex and axial-tangential movement simultaneously. Furthermore, recirculation phenomena in any swirl combustor are complicated due to their three-dimensional and time-dependent nature. Thus the formation of the CRZ is governed by many factors such as swirl strength, burner configuration, equivalence ratio, flow field characteristics at the exit plane, etc. [96]. Swirling flows can also lead to the formation of three-dimensional unsteady vortex structures. These structures rotate around the swirling central axis affecting the mixing and subsequently the combustion process [104].



Figure 2-3 tangential velocity profiles [74].

Moreover, these structures have the ability to produce low-frequency modes matching the system's hardware natural frequencies, hence initiating oscillations that can damage the system. The precessing vortex core (PVC) represents one of these structures, arising from the random form of vortices [74, 96].

Other swirling flow structures appear depending on the burner configuration. Central body or central fuel injectors have been widely used in swirl burners in order to enhance flame stability. However, the use of these configurations can result in a considerable change of swirl structure characteristics downstream the exit plane. In these cases, outer recirculation zones (ORZS) located between the corners of the combustion chamber and annular jet or the high momentum flow region (HMFR) will appear. Shear layers (SL) located between internal recirculation zones (IRZ) and the high momentum flow region can be highly distorted [105-108]. Syred and Beér [102] found that cyclone combustors

could have three large recirculation zones that provide a longer residence time for the combustion of the reactants.

Swirl structures have been investigated extensively to achieve a deeper understanding of their interaction. The relation between PVC and other structures, especially CRZ_S has special importance due to their mutual coexistence. The presence of the (PVC) was firstly reported by Syred and Beér [102]. They detected the existence of this three-dimensional time-dependent structure by using pressure transducers. They found the PVC to be related to the formation of the CRZ, its frequency, and swirl number. Sarpkaya [109] found traces of the nature of the vortex breakdown, observing the possibility of the existence of two PVC_S precessing around the recirculation zone. However, the undesirable oscillations of the PVC have motivated many researchers to investigate stability damping mechanisms that can minimise the operation instabilities. Yazdabadi, et al. [110] found that by changing the downstream configuration of the cyclone combustor, it is possible to reduce the PVC frequency significantly. In addition the, PVC can be delayed or promoted by these configurations. Other studies have proved the distortion of the PVC when changing the nozzle [100, 111].

One of the most comprehensive studies that describe the correlation between swirling flow structures is done by Syred [74]. The study correlates the formation of the central recirculation zone (CRZ) and the precessing vortex core (PVC) with the burner geometry and swirl number. The higher the swirl number, the stronger the CRZ. Nevertheless, there is a lower limit of swirl to allow the generation of the vortex breakdown. Consequently, the formation of swirl coherent structures normally starts between a swirl of 0.6 and 0.7.

Moreover, the study also indicates that the behaviour of swirl structures, especially the PVC, becomes more complex under combustion conditions. In the case of combustion, the method of fuel entry, the confinement level, and the equivalence ratio play important roles in the occurrence and characteristics of coherent structures. For example, using axail fuel injection can play an important role in the suppression of the PVC amplitude, while the high level of confinement of the combustion chamber can alter the shape and size of the CRZ considerably. This effect of confinement on the swirl structures has been also proved by [100, 102, 105].

Modern visualisation techniques have been used effectively to investigate different types of swirling flows and coherent structures. Valera-Medina, et al. [60] recognised and identified new structures of swirling flows by using different measurement techniques; Hot wire anemometry, high-speed photography and particle image velocimetry studies were carried out on isothermal flows with tangential swirl burner. The co-existence of the PVC and CRZ, their shape and interaction depends on the burner geometry and flow regime. Interestingly, a significant finding of this study was the recognition of a canal formed by the PVC on the side of the CRZ which evidence the interaction between them. Furthermore, the generation of a second central recirculation zone (CRZ2) with confinement conditions was also demonstrated.

Many studies have proven the helical nature of the PVC, its centre precessing around the geometrical centre of the combustor at a fixed frequency and geometrical features [60, 74, 112] as can be seen from figure 2.4 [74]. Although the PVC promotes good flame stability through increasing the turbulence intensity, hence better mixing of fuel and oxidizer, its helical motion forced the flame stabilisation point to move at a certain frequency and in consequence asymmetric heat fluctuations hence initiate flame instabilities [113]. The main concerns about the effect of the PVC are manifested through the possibility of coupling between its frequency and thermoacoustic oscillations that could lead to severe damages to the system hardware.



Figure 2-4 Helical nature of the precessing vortex core [74].

2.4.2 Characterisation of Flow Field

2.4.2.1 Swirl Number

Swirling flows can be generated by the following three methods [63, 105, 114]

- 1. Tangential entry (axial-tangential swirl generator),
- 2. Guide vanes (axial swirling vanes with flat or twisted blades),
- 3. Direct rotation (rotating pipe).

It is quite common that swirling flames are arranged and controlled by Swirl number. Any change in its value affect the flame dynamics directly [63, 98]. Swirl number can be defined as the ratio of the axial flux of tangential momentum to the axial flux of axial momentum times the equivalent nozzle radius [105].

$$S = \frac{Axial \ flux \ of \ tangential \ momentum}{Axial \ flux \ of \ axial \ momentum \ \times Exit \ radius} = \frac{G_{\theta}}{G_{x}R_{0}}$$
(2-5)

Where :

$$G_{\theta} = \int_{\theta}^{R} (\rho u w + \rho u' w) r^2 dr \qquad (2-6)$$

$$G_{x} = \int_{0}^{R} (\rho u^{2} + \rho u^{2} + (p - p_{\infty} r^{2})) dr$$
(2-7)

- *u* : axial velocity [m/s]
- u': fluctuating axial velocity [m/s]
- w: tangential velocity [m/s]
- w': fluctuating tangential velocity [m/s]
- ρ : density [kg/m³]
- r : Radius [m]
- R_0 : nozzle exit radius

Nevertheless, determining the swirl number by using the above equation seems to be difficult because it requires an evaluation of all velocity and pressure profiles at different conditions which may give a different swirl number for each point in the flow field [100, 115].

Thus, according to Syred and Beér [102] and Fick [111] neglecting pressure variations across the flow (isothermal conditions), it is possible to evaluate a swirl number from the geometry of the swirl burner as follows:

$$S_{g} = \frac{R_{0} \pi R_{eff}}{A_{t}} \left[\frac{Tangential \, flow \, rate}{Total \, flow \, rate} \right]^{2}$$
(2-8)

Where:

 R_0 : the exit radius [m]

- R_{eff} : the radius at which the tangential inlets are attached with respect to the central Axis of the burner [m]
- A_t : the total area of the tangential inlets [m²]

For tangential swirl burner

Tangential flow rate = Total flow rate

$$S_g = \frac{R_0 \pi R_{eff}}{A_t} \tag{2-9}$$

However, determining the swirl number by the above equation has some drawbacks since it is based on idealised assumptions of perfect mixing between axial and tangential streams. Moreover, the swirl number may undergo significant decay downstream of the swirl generation plane, since combustion conditions cause thermal expansion which leads to significant increase in the axial velocity. Thus combustion swirl numbers can be calculated by the criterion proposed by Syred and Beér [102],

$$S_{g, comb} = S_{g, iso} \frac{T_{inlet}}{T_{outlet}}$$
(2-10)

Where:

 $S_{g, comb}$ = the combustion swirl number [-] $S_{g, iso}$ = swirl number at isothermal conditions [-] T_{inlet} = the inlet temperature of the fluid [K] T_{outlet} = the outlet temperature of the fluid [K] Although there are difficulties when determining the swirl number during combustion, some studies [98, 116] have used the Laser Doppler Velocimetry (LDV) to measure the velocity profiles to evaluate the actual swirl number. Interestingly the calculated swirl number was lower than that estimated from standard calculations based on geometrical parameters. Experimental evaluation of swirl numbers is becoming more important in determining the fluctuation of heat release and as a consequence more accurate diagnosis of combustion dynamics and instabilities.

Swirl numbers can also be theoretically calculated by using equation (2-5) by calculating the axial flux of tangential momentum and axial flux of axial momentum using Reynold Navier-Stock equations under boundary layer assumptions. However, this method is limited in applications with weak swirling flows [113]. Swirl generation efficiency represents an important feature that characterises swirl combustor performance, and it is a function of swirl number. It is defined as the ratio of kinetic energy flux of swirling flow through the burner throat to the drop of static pressure energy between air inlets and burner throat. Thus it varies depending on the swirl configuration or geometry. Efficiency of axial –tangential swirl generators decreases to about 40% at S=1 while its value are about 70-80% for radial-flow guide vane systems [114].

Swirl number is the crucial factor for the onset of the vortex breakdown and hence the size and shape of the CRZ. Syred and Beér [102], Beer and Chigier [114] stated that the required swirl number for the onset of the (CRZ) is S<0.6. Other studies [19, 117] found that it is possible to obtain (CRZ) even at low swirl numbers. However, the swirl degrees can be classified as follow[105]:

- ➤ Very weak swirl (S≤0.2) in this case swirl strength and related velocities yield rapid decay consequently the amount of generated adverse pressure gradient is not sufficient to generate axial recirculation.
- ➤ Weak swirl (S≤0.4) up to this swirl number the axial velocity profiles are in Gaussian form, and the maximum velocity values are on the jet axis. Nevertheless, these maximum values are displaced from the jet axis with swirl number greater than 0.5 with formation of the recirculation zone [114].
- Strong swirl (S≥0.6) at this swirl number the radial and axial pressure gradient reach higher values, and as a consequence the flow cannot be overcome by fluid kinetic energy, enabling the formation of the recirculation zone.

However, determining the minimum swirl number or the so called (S_{cr}) which is required for the formation of the vortex breakdown and hence the swirl structures and set it as a universal value for all swirling flows, seems to be a difficult task. Because it depends on the particular flow features, which are different for all burner geometries [118].

2.4.2.2 Strouhal Number

It is a dimensionless number describing an oscillating flow mechanism. Thus it can be used to characterise and analyse swirling flows, and it can be defined by the following formula:

$$Sr = \frac{fD}{U_b} \tag{2-11}$$

Where:

f: Oscillation frequency (vortex shedding) [1/s]

D : Characteristic length (exhaust diameter of swirl burner) [m]

 U_b : Average bulk burner exit axial velocity [m/s]

Strouhal number can be used successfully to describe flow unsteadiness and periodic vortex shedding; it can be correlated with other dimensionless numbers such as Swirl and Reynolds numbers to characterise the PVC [74, 97, 100].

2.5 Operability Issues of Swirl Combustors

All swirl combustors are designed to work efficiently, safely and reliably with reasonable operability limits. This means the burner should have the ability to hold a flame at a variety of equivalence ratios, different fuel types and different environmental conditions (stationary gas turbines or aviation applications). However, most industrial combustors have undergone some operational issue that can cause considerable degradation of their performance and as consequence overall system efficiency. The common operability problems of swirl combustors are as follows;

2.5.1 Blowoff

Blowoff means the flame is being detached from its stable position and physically blown out. The most common reason that can cause flame blowoff is very lean mixtures or robust combustion instabilities. Recently swirl combustors have been extensively employed with lean premixed combustion to reduce NO_x emissions, which are required to operate close to their blowoff limits, hence the strong possibility of being subjected to this issue. This problem becomes significantly important, especially for aircraft engines. Blowoff is a complicated phenomenon that involves interaction between chemical reactions, flame propagation and flow dynamics [35, 119]. Blowoff phenomena have been extensively investigated, and a variety of experimental techniques and numerical simulations have been employed to achieve better understanding to prevent its occurrence [19, 59, 119-127].

It is thought that blowoff is governed by the equilibrium between chemical kinetics and fluid mechanics. Thus this relation can be described by using a Damköhler number based on a well-stirred reactor [35, 121] which represents the relation between residence time and chemical kinetic time.

$$Da = \frac{\tau_u}{\tau_c} \tag{2-12}$$

Where

- τ_u : The period that reactants remain in the reaction zone or flow time scale [s]
- τ_c : The required time for chemical reaction [s]

Thus blowoff will occur if the residence time is shorter than the chemical time [40, 128]. The residence time is denoted by the ratio between the characteristic length (recirculation zone), and the characteristic velocity scale, while chemical time scale represents the ratio between thermal diffusivity, and laminar flame speed S_L [129], thus blowoff can be correlated with the Damköhler number as follows:

$$Da = \frac{d_{U_b}}{\alpha_{S_L^2}} = \frac{S_L^2 d}{\alpha U_{ref}}$$
(2-13)

Where

Da : Damköhler number [-]

- *d* : characteristic length (recirculation zone length) [m]
- U_b : average bulk burner exit axial velocity [m/s]
- α : thermal diffusivity [m²/s]

S_L : Laminar flame speed [m/s]

However, determining the characteristic velocity, the characteristic length and hence the residence time is not simple. The characteristic velocity is a function of approaching flow velocity U_b which may change with Reynolds number due to changes in burning gas temperature while the characteristic length is a function of the recirculation zone shape. Additionally, calculations of thermal diffusivity for a fuel of different blends is sophisticated. Hence the evaluation of chemical time scale becomes more complex. This investigation of the chemical kinetic is not sufficient to have a comprehensive understanding about blowoff; fluid mechanics must be accounted too [40, 129, 130]. Nonetheless, blowoff trends of a wide range of fuel compositions can be determined by using the Damköhler number [33].

The relation between flame strain and blowoff can be used to predict the mechanism of blowoff. Whereas flame strain occurs under certain levels of flow velocity gradients and it can withstand limited levels of strain before extinction, this level is denoted by an extinction strain rate K_{ext} which is a function of the fuel composition, chemical time scale and stoichiometry [121, 131], thus Damköhler number can be written as follows;

$$Da = \frac{K_{ext}}{K} \tag{2-14}$$

 K_{ext} : extinction strain rate [1/s]

K : flame stretch

Many mechanisms can be used to enhance flame resistance to blowoff. Swirl burners and bluff-body techniques have been employed in many applications to achieve flame stability. They can work as premixed flame holders, enabling mixing between combustible mixture with hot combustion gases and promote the continuous ignition of reactants [122, 123, 125].

The flame holding represents a crucial factor in flame stability. Consequently, many studies have demonstrated different characteristics of flame holding and its effect on blowoff limits by using bluff-bodies. Lieuwen, et al. [132] studied the effect of bluff-body geometries. The team pointed out that vortex shedding results from the bluff-body playing an important role in the ratio of gas expansion across the flame and hence the

level of flame extinction, and as a consequence the blowoff mechanism. Roy Chowdhury and Cetegen [133]correlated the turbulence level with strain rate; they found that strain rate increased with increasing turbulence, consequently leading to localised extinction zones along the flame surface. Chaparro and Cetegen [122] also indicated that blowoff equivalence ratios are affected directly by upstream flow extinction.

Due to recent trends from using lean premixed combustion and alternative fuels, introducing hydrogen or high hydrogen blends as alternative fuels have increased significantly. Blowoff and flame extinction limits are directly proportional to the percentage of the hydrogen content of a certain fuel. Many studies have proven that the equivalence ratio at which blowoff occurs can be decreased by increasing hydrogen percentage. Working at this conditions enables burners to operate at lower flame temperatures consequently low NO_x emissions [19, 134-136].

2.5.2 Autoignition

Under certain conditions inside the premixer such as high pressure or high temperature, it is possible that the mixture ignites without an external energy source this occurs especially in premixed flows an increase in temperature may be caused by early incidence of chemical reactions of the mixture inside the premixer at high pressure and/or high temperature. This results in considerable amount of heat generation leading to spontaneous ignition. In general, autoignition is defined as the homogeneous ignition of the combustible mixture upstream of the combustion chamber [35].

Autoignition is a serious issue that can stain the performance of gas turbines and can cause considerable damage to the system. Thus it is crucial to predicting the ignition time of each mixture to achieve more secure operation conditions. Autoignition time delay is defined as the time interval between the injection of the mixture into the combustion chamber and the onset of the flame. This time can be represented using Wolfer equation [39].

$$t_i = 0.43P^{-1.19} \exp(4650/T_m) \tag{2-15}$$

Where

 t_i : the ignition time delay [ms]

P : the pressure [bar]

T_m : the initial mixture temperature [K]

Nevertheless, in addition to the effect of pressure and temperature, fuel composition plays an important role in determining the autoignition time delay [33]. Natural gas or methane is the favourable fuel for stationary gas turbines and has reasonably weak ignition behaviour at low temperatures. However, this behaviour can be affected drastically when adding higher hydrocarbon blends [137]. Recently, hydrogen has proved to be a clean energy carrier, and it can be added to natural gas or methane to enhance some operational requirement, as a consequence of its delay time, wide flammability limits, high energy density and high laminar flame speed [138]. However, the percentage of hydrogen addition significantly impact the autoignition mechanisms. The higher the hydrogen concentration, the more the hydrogen chemistry dominates the mixture ignition, whereas the non-linear relationship between logarithmic of ignition delay and equivalence ratio becomes increasingly evident with increasing hydrogen fraction [139].

In swirl combustors, strong aerodynamic effects are directly linked to predicting autoignition times. However, weak aerodynamic system designs make them more likely to suffer flashback. Thus, ignition delay times of the fuel used and the local conditions should be taken into consideration when designing the premixer to avoid autoignition [33].

2.5.3 Flashback

Flashback is the upstream flame propagation from the combustion zone to the premixing zone inside the combustor. It occurs when the flame speed exceeds the approaching flow velocity [39]. Flashback is a main operability issue and an inherent instability problem of lean premixed combustion systems which can cause the burner to overheat with considerable damages to the system hardware and increase in pollutant emissions [68, 140, 141] as can be seen from figure 2.5.

Fuel composition represents an important factor affecting flashback mechanisms. This effect is characteristic of turbulent flame speed S_T variations [33]. Thus introducing hydrogen or hydrogen fuels to premixed combustion systems that use swirl combustors have brought many concerns regarding flashback. More reactive gases than a natural gas

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such as hydrogen can present flame speeds about four times greater than those of natural gas [118, 142, 143].

Swirl lean premixed flames are more likely to be subjected to flashback [140, 144, 145]. Swirling flows are complex and can interact with the flame in these systems causing flame flashback [146].



Figure 2-5. System hardware damages due to flashback.

According to previous literature [65, 66], there are four types of flashback mechanisms,

2.5.3.1 Boundary Layer Flashback

This type of flashback was firstly described by [147]. They investigated laminar nonswirling flow flashback, pointing out that the reason of flame flashback in the wall boundary layer is due to the reduction of the approaching flow velocity near the wall in such way that it became lower than the flame speed. Thus, they proposed the so-called "Critical gradient" model. This model defines the flashback limits according to the equilibrium between the wall-parallel boundary layer velocity u(y), where (y) represents the wall - normal coordinate and the flame speed S_f (y) at a distance dq from the wall which represent the quenching distance, as can be seen in figure 2.6 [148].



Figure 2-6. mechanism of wall boundary layer flashback [148]

Thus the critical velocity gradient g_c at the wall can be used as a measure of laminar wall flashback limits and can be represented by the following equation;

$$g_c = \frac{\partial u}{\partial y} = \frac{S_f(y)}{d_q} \tag{2-16}$$

 g_c : critical velocity gradient [1/s]

 $S_f(y)$: flame speed [m/s]

 d_q : quenching distance [m]

The critical velocity gradient g_c , in general, depends on equivalence ratio, fuel-oxidiser kinetics, static pressure, static temperature, wall temperature, the state of the boundary layer and the geometry. According to equation (2-16), flashback can take place when the velocity gradients at a distance d_q from the wall, becomes lower than the flame speed. Lewis and von Elbe [147] proposed the criterion of critical velocity gradient for laminar poiseuille flow in tubes according to the following equation;

$$g_t = \frac{\partial u}{\partial r}\Big|_{r=R_t} = \frac{4U_b}{\pi r^3}$$
(2-17)

r: the radial coordinate [m]

 R_t : radius of the tube [m]

 U_b : bulk flow velocity [m/s]

Although the above mentioned critical velocity gradient model in equation (2-16) is considered as state of knowledge been intensively used to describe the physical situation of BLF, it does not take into consideration the flame flow interaction. Moreover the flame speed could be considerably affected by heat transfer between flame and burner wall in addition to flame stretch effects. Therefore, critical velocity gradient must be determined experimentally for different fuels [53].

Turbulent flows lead to increase in flashback propensity, whereas the critical velocity gradient increases during the transition from laminar to turbulent flow [149]. Many studies have investigated different mechanisms to reduce or avoid wall boundary layer flashback. Eichler and Sattelmayer [148] investigated flame flashback of hydrogen rich mixtures on a flat plate by using Reynolds-average Navier Stock computations. Interestingly, they found that there is considerable deviation of flashback results than those in the case of tube burners and the effect of adverse pressure gradients on the boundary layer structure, resulting in a totally different interaction between flame and flow. This fact has also been proved by Eichler and Sattelmayer [67], who found that flashback limits of wall boundary layer flashback are governed by flame tip quenching in the backflow region. This region is continuously dislocated down the upstream direction by laminar flames while it develops low-velocity structures close to the wall in the case of turbulent flames.

Baumgartner, et al. [53] investigated the transient phenomena during the onset of BLF; they found that for unconfined flames the flashback propensity is governed by flame speed and heat exchange between the flame base, leakage flow and burner rim. The most interesting finding of this study is that the adverse pressure generated from flame backpressure, leads to retardation and reflection of the approach flow upstream of the leading flame tip, at flashback conditions this flow reflection drive the subsequent flame propagation. The study also revealed that for unconfined flames, the leading flame tip propagates with some vertical offset from burner wall, this offset, in turn, lead to increase the leakage flow of fresh mixture between flame and wall.

Recently Lin, et al. [149] proposed a methodology for correlating the flashback propensity with turbulent flame speeds (S_T) for H₂ rich gases based on measurement of critical velocity of the flame and flow as follows;

$$g_c = \frac{S_T}{(Le \times \delta_{lo})} \tag{2-18}$$

- g_c : the critical velocity gradient of the flame [1/s]
- *L_e* : Lewis number which represents the ratio between thermal diffusivity to mass diffusivity [-]
- δ_{lo} : (un-stretched) laminar flame thickness [m]

The critical velocity gradient of the flow is represented by the following equation;

$$g_f = 0.03955 \times U_b^{7/4} \times v^{-3/4} \times d^{1/4}$$
(2-19)

- g_f : the flow velocity gradient [1/s]
- U_b : the bulk velocity at the combustor inlet [m/s]
- v : the kinematic velocity of the combustible mixture [m²/s]
- *d* : the diameter of combustor inlet [m]

Thus, according to these relations, the flame stability depends on the balancing between (g_f) and (g_c) , whereas flashback takes place when (g_c) exceeds (g_f) .

2.5.3.2 Turbulent Core Flashback

Turbulent core flashback takes place when the turbulent flame velocity (S_T) exceeds the local flow velocity in the core flow. Then the flame can propagate upstream in all the zones of the burner that have lower velocities than the turbulent flame speed. Thus one of the fundamental design parameters of new burners is how to ensure that flow fields do not have strong local velocity deficits with axial flow velocities that must be higher than the turbulent flame speed (S_T) [35]. Turbulent flame velocities depend on chemical kinetics, turbulent structures, length scales and velocity fluctuations.

Determination of turbulent flame speed has encountered many difficulties represented by complex interactions between turbulence and chemistry [65]. In swirling flows, the relation between turbulent flame velocity and level of turbulence fluctuation is crucial in determining the flashback. The flame can propagate if turbulent velocity is higher than characteristic turbulent velocity fluctuations, whereas the degree of turbulence fluctuation in the mixing zone is governed by the swirl strength [35]. Furthermore, swirl flame structures are highly wrinkled and corrugated. Thus flame surface increases considerably. Consequently, this effect results in turbulent flame velocities being higher than the

laminar value. According to this correlation between turbulent and laminar flame speed, it can be concluded that the burners of low turbulence levels are more suitable for burning fuels of high laminar burning velocities than high swirl designs to achieve good flashback resistance [65].

Many studies about the determination of turbulent flame velocity have been published Lipatnikov and Chomiak [150] found that the turbulent flame velocity increases with decreasing Lewis number which depends on fuel composition. Kobayashi, et al. [151] measured the turbulent flame velocity of gasification syngas flames burned with air $(CO/H_2/CO_2/air flame)$ by using OH-PLIF and flame radiation. The results were compared with those of $(CH_4/air flames)$. They found that the ratio of turbulent flame velocity to the laminar value (S_T/S_L) of gasification syngas flames $CO/H_2/CO_2/air$ flame are higher than that of CH_4/air flames. Beerer, et al. [152] used Laser Doppler anemometer (LDA) to detect the local flame displacement and hence determine the turbulent flame velocity. They found the turbulent flame velocity for hydrogen/methane flames in low swirl is a function of fuel composition and turbulence intensity.

Although, there is vast literature on different techniques and simulations to predict the turbulent flame velocity of various combustion systems, there is still a crucial need to have a more practical and robust methodology for investigating turbulent flame velocity especially for H₂-fuels [149]. Consequently, designing combustion systems to have a higher axial flow velocity, hence avoid turbulent core flashback, is still based on practical experience.

2.5.3.3 Flashback due to Combustion Instabilities and Coherent Structures

Combustion or thermo-acoustic instabilities inside the combustion chamber can cause upstream flame propagation both in the turbulent core and boundary layer, consequently flame flashback. These types of instabilities arise when unsteady or periodic heat release couples with one or more acoustic modes of the combustor [62, 153]. It is possible for these oscillations to interfere with the combustion system causing a considerable amount of vibrations which lead to failure or collapse of the combustion system [52, 154].

Fluctuations in the combustor are attributed to different mechanisms, one of these mechanisms is high pulsation levels, which cause a modulation of the flow field in the

burner. This modulation flow velocity undergoes a periodic drop below the time average, and large scales of vortices will be generated. If the frequency is low enough, there will be upstream flame propagation [35]. Another cause of instabilities is equivalence ratio perturbations which have a direct effect on a chemical reaction, flame speed and indirectly on the flame area which in turns control the amount of heat release response [155, 156]. Inlet velocity fluctuation is also considered as one of heat release fluctuations, Balachandran, et al. [157] found that heat release becomes nonlinear after a certain level of inlet velocity amplitude.

Swirling flows and their coherent structures, especially the precessing vortex core (PVC), also represent another source of unsteadiness or fluctuation. The (PVC) can trigger the formation of radial eddies and alter patterns of lean and rich combustion which in turn consolidate combustions oscillations [74]. The process of development of these fluctuations inside the combustor is governed by a feedback cycle where heat release rate fluctuations add energy to the acoustic field. Consequently there will be acoustic pressure and velocity fluctuations propagating throughout the combustion chamber. These fluctuations in turn trigger coherent structures and equivalence ratio oscillations leading to further heat fluctuations closing to the feedback loop [62].

Protection of combustion systems from being subjected to combustion instabilities during stable operation represents an urgent need from the point of view of high-performance efficiency and elongating system life. Thus, it is required to have a suitable damping mechanism that can mitigate or totally eliminate the effect of excitation mechanisms. Combustion instabilities can be suppressed by either active or passive techniques. Some common damping mechanisms used in industry are [158],

- Radiation when some acoustic energy is propagated to the far field.
- Convection of acoustic energy.
- Thermal and viscous dissipation which is done by either imperfect reflection of acoustic waves due to boundary layer near rigid walls or by flow separation and vortex shedding at the sharp edges.
- Energy transfer between acoustic modes. Combustion instability occurs at different frequencies, and since nonlinear mechanisms allow for transferring energy from unstable modes to other modes, thus this extraction of energy from unstable mode has stabilising effects on the system.

However, the degree of effectiveness of combustion instabilities controllers is different from system to system. Some techniques perform differently with different combustion conditions. Although a certain level of combustion instabilities suppression is possible, more research is required for new systems to achieve a deep understanding of their combustion dynamics, hence effectively control combustion instabilities in a more generic way [159].

2.5.3.4 Flashback due to Combustion Induced Vortex Breakdown (CIVB)

Combustion induced vortex breakdown has been identified as one source of sudden flame transition and flashback in swirl combustors. This sudden transition can cause severe system failure due to the risk of overheating of upstream burner positioned components. This type of flashback is substantially different from the previous three flashback types. It occurs primarily in swirl stabilised combustors without a centre body at stable operation due to small changes in the flow field and the interaction of the turbulence and chemistry at the tip of the central recirculation zone [65, 160, 161]. It can occur even if the flow velocity is higher than the turbulent flame speed [162]. The earliest study that denotes this phenomenon was done by Kröner, et al. [66]. The group investigated the flashback limits for swirl burners with cylindrical premixing tubes without using centrebodies.

They found that the quenching effect of chemical reactants plays the dominatn role in flashback limits. Another study [65] found that flame propagation is governed by mutual feedback between chemical reaction and flow. In other words, flashback has high sensitivity to the momentum distribution in the vortex core, whereas the chemical reaction leads to variation in pressure boundary conditions at the inlet of the combustion chamber, consequently, promoting the propagation of the vortex breakdown. The study has proved the possibility of significant flashback resistance mechanisms when introducing axial jets, however, altering swirl levels. The effect of the axial jet has also been investigated by [163] who found that the increase in axial jet diameter leads to thicker vortex cores. Hence the vortex breakdown will be shifted downstream and consequently stabilise the burner operation window.

Vorticity transport equations have been employed by [141, 162, 163] to analyse the source terms of azimuthal components of the flow field. They argue that the baroclinic torque

produces a considerable amount of negative axial velocity in the vortex core which in turn results in the CIVB flashback. They also correlated the phenomenon of stagnation of the recirculation zone with the effect of changing the equivalence ratio of the axial position. Based on their investigations the upstream flame envelope and the stagnation point play the dominatn role in the onset of the CIVB. If the recirculation zone can pass the stagnation point in the upstream direction, the volume expansion upstream of the recirculation zone generates positive azimuthal vorticity consequently positive axial velocity. Thus stable flame prevents the occurrence of the CIVB. However, upon further upstream flame propagation (a further increase of equivalence ratio) baroclinic torque will increase, thus the high possibility of CIVB occurrence increases as well.

Since swirling flows are characterised by high complexity turbulent flows, the CIVB flashback cannot be only correlated to the balance between flow velocity and flame speed. Thermal and fluid dynamic properties are also crucial to the stable reaction during upstream flame propagation. The high stretch rates and heat loss in the reaction zone can also cause extinction or flame quenching, whereas with increasing the stretch rate the heat release becomes limited by reaction kinetics and as a consequence, disturbing the balance between the generation of products and removal from the reaction zone. Kröner, et al. [66] proposed a new correlation to define flashback limits. This correlation stated that flame quenching occurs at a critical value of the quench parameter (C_{quench}) which represents the ratio between the characteristic chemical time scale (τ_c) and flow or mixing time (τ_u)

$$\frac{\tau_c}{\tau_u} \ge C_{quench} \tag{2-20}$$

Where

 $C_{quench} : \text{ quench parameter [-]}$ $\tau_c : \text{ chemical time scale} = \frac{\alpha}{S_L^2} \text{ [s]}$ $\alpha : \text{ thermal diffusivity [m^2/s]}$ $S_L : \text{ laminar burning velocity [m/s]}$ And $\tau_u : \text{ flow time scale} = \frac{d}{\overline{u}} \text{ [s]}$ Where
- *d* : mixing tube diameter [m]
- \overline{u} : Characteristic or the mean axial velocity [m/s]

Thus, the quenching constant can also be defined as follows;

$$\frac{\alpha \bar{u}}{S_L^2.d} \ge C_{quench} \tag{2-21}$$

The above equation has been modified according to the Peclet number approach,

$$\frac{\bar{u}.d}{\alpha} = C_{quench} \cdot \frac{S_{L}^2.d^2}{\alpha^2}$$
(2-22)

$$Pe_u = C_{quench} \cdot Pe_{SL}^2 \tag{2-23}$$

Where (Pe_u) is the flow Peclet number and (Pe_{SL}) is the flame Peclet number.

Thus, according to the above correlation flashback behaviour of a certain burner can be characterised by using the quenching parameter, or in other words, by knowing the quenching parameter or quench distance for a mixing tube, flashback can be avoided. The above correlations have been further modified by Kroner, et al. [164], chemical time scale has been defined according to the relation between the thickness of the reaction zone and the transport due to thermal diffusivity. In this relation, the time scale is the measure of residence of the hot reacting gases inside the reaction zone of thickness δ_R where,

$$\delta_R = \frac{\alpha}{S_L} \tag{2-24}$$

Thus, chemical time scales can be defined according to the following equation,

$$\tau_c = \frac{\delta_R^2}{\alpha} \tag{2-26}$$

Almost all the above works showed that the stability regime close to the burner mouth is a function of the interaction between flow field and heat release from the combustion process. This interaction is governed by fuel type, composition, unburned mixture conditions and local stoichiometry, which mean it is crucial to understand the effect of fuel type and swirl burner geometry regarding their impact on the turbulence characteristics close to the burner mouth. The fuel type or fuel blends play an important role in the CIVB flashback mechanism, whereas using high hydrogen increase the onset of the CIVB. Several studies [33, 71, 143] stated that the behaviour of fuel mixtures could be considerably different from that of individual constituents, while [73, 165, 166] showed significant differences in flashback limits upon increasing the hydrogen content of the fuel.

Swirl flow characteristics and local equivalence ratios directly affect the interaction at the burner mouth and hence the governing conditions of the CIVB flashback. There are some studies have investigated the effect of swirl strength and mixing degree, Baumgartner and Sattelmayer [167] found that the flashback mechanism is altered from boundary layer flashback BLF to CIVB flashback with increasing equivalence ratios at low swirl intensities. Sayad, et al. [72] found that CIVB flashback occurs at higher swirl numbers than that of boundary layer flashback BLF. The effect of geometrical issues on CIVB flashback has also been investigated, and many studies about the effect of changing geometries on the enhancement of CIVB flashback resistance have also been implemented. Central bluff bodies or lances are widely used in swirl burners to avoid upstream flame propagation or CIVB flashback.

However, some works used axial injection to increase flashback resistance albeit the possibility of working with partial premixing and swirl number reduction. Reichel, et al. [70] found that high amounts of axial air injection reduce the defect in axial velocity, which influences the vortex breakdown (VB) Position and consequently enhance the CIVB resistance.

Other works, Konle and Sattelmayer [163] Have found that the increase in the diameter of the central injector will increase the thickness of the vortex cores and strengthens the axial flow velocity. Thus the vortex breakdown is shifted downstream producing better flame stability. The findings have also been stated by [141, 145] who found that axial injection results in wider vortex cores lead to lower pressure gradients hence baroclinic torque reduction, and improving system resistance to CIVB flashback. Mayer, et al. [168] found an improvement of the flashback resistance by using axial fuel injection when compared to trailing injection where the trailing injection was done by adding rows of injector holes placed along four trailing edges of a vanned type swirler. Lewis, et al. [11] studied the effect of axial injection of methane, air and carbon dioxide, finding that the coherent structures changed significantly according to the type of gas injected. Furthermore, they proposed a correlation between the reciprocal effect of the high momentum flow region (HMFR) and the central recirculation zone (CRZ). According to this relation, the CIVB flashback occurs when the CRZ is squeezed by the high momentum flow region HMFR that evolves through the flow. Another study in this context was done by Sattelmayer, et al. [87]. They found that using diffuser between the burner mixing tube and the combustor can alter the flashback mechanism from CIVB flashback to wall boundary layer BLF. Moreover, they stated that similar flashback limits for BLF and that initiated by CIVB for certain swirl combustor could be achieved by controlling its aerodynamics, and this represent the optimum leads to deterioration of the flashback resistance of one of the two types consequently worsen flashback limits.

However, although many works have been implemented to have a deep understanding of the CIVB flashback mechanism, there are still many unknowns. The effect of different geometries and fuels on flashback phenomena are still generally predicted in a qualitative sense.

2.6 Fuel Blends Used for Gas Turbines

One of the significant features that characterise gas turbines is their high fuel flexibility. They can consume a vast spectrum of fuels ranging from high to low-level quality. Nonetheless, the concerns about the possible pollutant products that may generate from burning these fuels and their impact on climate have represented a serious challenge for fuel flexibility. Of particular interest to flashback, there are some specific fuel compositions, generally the most common fuels and fuel blends used in gas turbines that concern reserves underlying on flashback are,

2.6.1 Natural Gas or Pure Methane CH₄

Natural gas or pure methane represents the main fossil fuel used to operate gas turbines. It is the cleanest of all fossil-based fuels, and it has lowest carbon intensity, emitting less CO_2 per unit energy than other petroleum derivate fuels. Furthermore, it is characterised by high flexibility and abundance [169-171]. The lower heating value (LHV) and higher

heating values (HHV) of pure methane are 48 and 55 MJ/kg respectively making it suitable for firing many gas turbines.

Nowadays natural gas is the fastest-growing primary energy source and a vital component of energy supply in the world. It has a considerable share amongst other energy resources in the energy landscape, and this ratio is predicted to increase in the future [7], whereas the global consumption of natural gas is predicted to increase up to (50%) by 2035 [172]. Natural gas mainly consists of 85%-95% methane and can include small quantities of ethane, propane, butane, nitrogen, carbon dioxide and sulphur compounds. Natural gas composition and purity can vary depending on many factors such as the production region, season and climate [173]. The chemical composition of natural gas is given in Table 2-1 [174].

Component	Typical analysis (vol%)	Range (vol%)	
Methane	94.9	87.0-96.0	
Ethane	2.5	1.8-5.1	
Propane	0.2	0.1-1.5	
Isobutane	0.03	0.01-0.3	
<i>n</i> -Butane	0.03	0.01-0.3	
Isopentane	0.01	Trace to 0.14	
<i>n</i> -Pentane	0.01	Trace to 0.14	
Hexane	0.01	Trace to 0.06	
Nitrogen	1.6	1.3-5.6	
Carbon dioxide	0.7	0.1 - 1.0	
Oxygen	0.02	0.01-0.1	
Hydrogen	Trace	Trace to 0.02	

Table 2-1 Chemical composition of natural gas.

2.6.2 Methane - Carbon Dioxide Blends

Carbon dioxide (CO₂) usually added to methane (CH₄) in order to reduce combustion temperature and hence reduce the nitrogen oxides (NO_x) emissions in addition to achieve better flame stability especially flame flashback [11, 170, 175]. The specific content of the fuel can be varied according to the operation conditions, in general, the percentage ratio of CH₄ is (50-75%) while for CO₂ is (25-50 %) [176].

2.6.3 Pure Hydrogen

Hydrogen has been considered as a promising clean energy carrier for the future and plays an important role regarding long-term sustainable energy supply. It characterises by high heating value around 141 MJ/kg, high combustibility, easy ignition and significantly zero or low level of harmful emissions to the environment compared with fossil fuels [87, 138, 177, 178].

Although its plentiful in nature, hydrogen is not found separately. Therefore its domestic production still faces some serious obstacles regarding cost, safety storage and operation. Moreover, its application is still restricted by geographic location, seasons and surroundings [179]. In general, hydrogen can be produced by using different technologies. Such as covering coal gasification with and without carbon capture and storage (CCS), natural gas reforming with and without CCS, water electrolysis, thermochemical water splitting with copper-chlorine (CU-CI) cycle, advanced nuclear technologies, biomass gasification and solar hydrogen production technologies [179-181].

2.6.4 Hydrogen – Hydrocarbon Blends

The unique characteristics of hydrogen represented by wide flammability range and high flame speed lead to many difficulties in terms of security of storage, transportation and combustion stability. Thus blending hydrogen with other hydrocarbon fuels represent a suitable option that can attain low emission requirements, high output power and reasonable stability margins [64, 182, 183].

Hydrogen is blended with different hydrocarbons in various proportions to produce the so-called syngas fuel. Most common and widespread fuels are hydrogen-methane blends. Those blends receive particular importance as alternative fuels used for gas turbine operation. The proportion of hydrogen in the mixture ranges from 5% to 90% [143]. However, the most obvious challenge to use blends is their increasing the propensity to flashback because of the considerable increase of flame speed [35, 71-73, 148, 149, 184].

2.7 Summary

A literature review of some topics related to swirl combustor characteristics and associated combustion instabilities such as blowoff and flashback, especially in the context of various alternative fuels and different geometries, has carried out in this chapter. It can be concluded that,

- Combustion technologies have been developed in order to meet minimum pollutant level requirements while keeping high operation efficiency of gas turbines. Lean premixed combustion using alternative fuels has proved as the most promising technology in this context.
- Swirl combustors are widely used in gas turbines owing to their great stability range of operability because of the unique nature of the generated swirling flows.

The important feature of these flows is the central recirculation zone (CRZ) which recirculates heat and active chemical species to the flame root, hence promoting the combustion process by excellent flame holding with reasonably low emissions.

- Despite their significant features, Swirl combustors are frequently subjected to different types of combustion instabilities, especially when they employ lean premixed combustion systems. These instabilities have arisen when using some alternative fuels especially high hydrogen blends. Common operation instability problems are blowoff, autoignition, combustion instabilities and flashback.
- Flame flashback has been considered as an inherent operability issue in swirl combustors that use premixed combustion. It can lead to severe damages in the system hardware; flame flashback occurs through four mechanisms; wall boundary layer flashback (BLF), turbulent core flashback, combustion instabilities and combustion induced vortex breakdown (CIVB) flashback.
- (CIVB) Represent state of the art amongst other types of flashback, because it can be initiated even at stable operation conditions and this mechanism is not

yet completely understood. Thus more investigation is required to have a deep understanding of this phenomenon and consequently find suitable techniques that can reduce or totally annihilated this type of flashback.



Experimental Setup and Methodology

Chapter 3

Any sufficiently advanced technology is indistinguishable from magic

Arthur C. Clarke

3.1 Introduction

Swirl combustors have particular importance in terms of performance and efficiency for gas turbines. Thus the accurate determination of their operating characteristics is considered a crucial factor in addressing many issues regarding their duties. However, swirling flows are so complicated that despite enormous development in computer softwares, precise approach still require a combination of numerical analyses with experimental trials. Laboratory facilities should have almost similar features to those of the real combustion systems in order to simulate their operation and consequently the possible operation instabilities. Thus, in this work, a 150 kW tangential swirl burner with other characterization systems has been used to study different flashback mechanisms to enhance flame flashback resistance.

3.2 Experimental Burner Rig Setup

A tangential generic swirl burner of a thermal capacity up to 150 kW, figure 3-1, has been used to investigate different flame flashback mechanisms [100, 185]. The system has two circular tangential inlets of 67 mm diameter each. The inlet area can be varied by using different inserts (25-25), (50-50) and (70-70) respectively in order to achieve variable swirl numbers. The system can also be used with three different nozzle configurations varying the exhaust diameter from 62 to 76 mm. Air was provided by a centrifugal fan via flexible hoses and two banks of rotameters to control the air flow rate. Another bank was used for natural gas injection.

The burner baseplate is designed mainly to allow central fuel injection in order to start burner operation. During stable operation, both air and gas are injected tangentially. Premixed gas injectors extend across the inlet ducts to allow adequate mixing. central fuel injectors are extended from the baseplate through the centre of the burner plenum opening just before the nozzle.



Figure 3-1. Experimental setup.

Swirl strength is an important factor to describe the performance of any swirl combustor. This strength is represented by the swirl number which was described in chapter 2. Determination swirl number by using equation (2-5) is difficult due to the complexity of the flows. Thus a geometrical swirl correlation suggested by [114], equation (2-9), has been used as an approach to defining the swirl number. The approach considers the following assumptions:

- 1. Isothermal conditions
- 2. Constant static pressure across the exit
- 3. Constant density throughout the burner
- 4. Uniform axial velocity in the exhaust

Figure 3-2 shows the dimensions used for geometric swirl number calculations.



Figure 3-2 Tangential swirl burner configuration.

3.3 Laser Doppler Anemometry (LDA)

Laser Doppler anemometry (LDA) is considered an effective measurement technique for fluid dynamic investigations. The non-intrusive principle, directional sensitivity, absolute measurements, high accuracy, high spatial and temporal resolution make it suitable for reversing flow and high-temperature chemically reacting media applications such as laminar and turbulent flows, aerodynamic measurements, supersonic flows, liquid flows, combustion characteristics of gas turbines[186]. Thus it has a good ability to determine turbulence mode parameters and consequently can characterise different features of swirling flows effectively.

Like other laser applications, LDA needs seeding particles in order to trace the flow, the type and size of those particles depend on the kind of measurement application.

3.3.1 LDA Principle of Operation

The principle of operation of the LDA is based on the well-known Doppler shift of frequency between transmitting and reflecting signals. When naturally buoyant particles are seeded within the flow, they scatter the light. Those particles are illuminated by a known frequency monochromatic laser light. The scattered light is

detected by photomultiplier tubes (PMT) which generate a current proportional to the absorbed photon energy, figure 3-3. The difference between the incident and scattered light frequencies is the Doppler shift [186].



Figure 3-3 Doppler phase shift of light scattered by a moving particle [186].

The Doppler frequency is a function of the particle velocity *V*, the direction of particle motion which is defined by an angle β , the light wavelength λ , the orientation of the observer which is defined by the angle α (the angle between the photodetector and incident light wave). Thus, the Doppler frequency can be described according to the following equation:

$$f_d = \frac{2V}{\lambda} * \cos\beta * \sin\frac{\alpha}{2} \tag{3-1}$$

It can also be measured directly by analysing the incident frequency f_i , and the observed frequency f_0 , where the difference between them is the Doppler frequency. However, the Doppler shift is very small compared to the incident light frequency (f_i) . Therefore this process has a high degree of uncertainty, and it is hard to have a direct derivation of velocity. Thus many LDA systems use dual-beam configurations; this configuration involves two coplanar laser beams with different directions intersecting at one point in space [186, 187]. The point of intersection is the measurement volume, and hence the particles that pass through this volume scatter the light from both beams as can be seen from figure 3-4.



Figure 3-4 intersections of two incident beams forming a measurement volume [186].

The crossing of two beams results in the generation of the so-called interference fringe patterns. Those patterns consist of alternating zones of brightness and darkness and the distance between two sequential bright or dark fringes is known as the fringe spacing (d_f) which can be determined by the following equation:

$$d_f = \frac{\lambda}{2\sin(\theta/2)} \tag{3-2}$$

Where (θ) : the angle between two incident beams.

In this case, the intensity of scattered light varies with the intensity of the fringes when a flow particles crosses the fringe patterns. Thus, as can be seen from figure 3-5, the signal burst varies with the time scale. This fluctuation time scale is $\left(\frac{d_f}{V}\right)$, where V is the velocity component perpendicular to the fringe patterns, thus equation (3-2) can be re-written as follows:

$$\frac{V}{d_f} = \frac{2V}{\lambda} \sin(\frac{\theta}{2}) \tag{3-3}$$

This value $\left(\frac{V}{d_f}\right)$ represents the Doppler frequency



Figure 3-5 signal burst record as one particle passes through fringes [186].

Noticeably, determining Doppler frequencies by two beams systems is not dependent on the position of the photodetector, compared to that determined by using equation (3-1). However, the Doppler frequency is only a function of the magnitude of Vregardless the direction, in other words, positive and negative values of velocity lead to the same Doppler frequency. This directional ambiguity can be corrected by making the frequency of one of the beams shifted by a known value, f_s . This value makes fringe patterns move at speed V_s towards the incoming unshifted beam as can be seen from figure 3-6, where

$$V_{S} = f_{S} d_{f} \tag{3-4}$$

Thus the frequency recorded by the photodetector is:

$$f_d = \left| f_s + \frac{2V}{\lambda} \sin(\theta/2) \right| \tag{3-5}$$

A single pair of incident beams can measure single velocity components perpendicular to the bisector of the two beams. Hence to measure another velocity component, it is possible either by rotate the laser probe by 90^{0} or add a pair of beams intersecting at the same point. However, each pair must have its own wavelength so that the burst signals can be recognised by the filters. The most common lasers used for multi-component measurement are blue, green and violet lasers. The measurement volume

size and shape is determined by the intersection of all beams, and its shape is almost elliptical.



Figure 3-6 removing directional ambiguity by frequency shifting [186].

3.3.2 LDA Configurations

In this study, a BSA flow light LDA has been used. It consists of the following components:

- A continuous wave laser generator Producing a monochrome, coherent, linearly polarised, low divergence (collimator), Gaussian intensity distribution laser. Mainly it is ND: Yag Laser
- > Transmitting optics, including beam splitter, Brag cell and focusing lenses
- Receiving system consist of receiving optics, spatial filter, interference filter and photodetector
- > A signal conditioner that contains an amplifier filter
- Signal processor that includes spectrum analyser, correlator, counter and tracker
- A Personal computer using BSA Flow software v5:20 DANTEC DYNAMICS
- Traverse system

LDA can be set up in forwarding scattering mode where (difficult to align and sensitive to vibration), or backwards scattering (which is easy to align). In this study backwards scattering has been used. Fig 3-7 illustrates the backwards scatter arrangement.



Figure 3-7 backscatter arrangement [186].

3.3.3 Laser Doppler Anemometry (LDA) Settings

Instantaneous velocity components downstream the burner mouth have been measured by Laser Doppler Anemometry (LDA). The LDA system was a one component Flowlight LDA system (Dantec) operated in backscatter mode. The tangential velocity component can be measured by turning the laser probe 90° .

The complex nature of swirl flows and the change in flow structure between isothermal and combustion cases required using two different settings. For isothermal flows, it was possible to acquire much more samples than that in the combustion case using a larger control volume (probe volume). However, during combustion the swirl structures impose different setting and strict control on the method to introduce seeding to the burner plenum. However, introducing less seeding can reduce the validity of results, while more seeding can lead to change in swirl flame structure or even initiate earlier flashback through changes in the flame local temperature hence affecting the residence time of the reaction.

The different settings of isothermal and combustion cases are illustrated in table 3-1. The seeding particle used was aluminium oxide AL_3O_2 of size ~ 10µm. Velocity measurements have been done at different positions Y = 5, 15, and 25 mm downstream the burner dump plane.

Parameters	Isothermal	Combustion			
Wavelength(nm)	532	532			
Focal length (nm)	500	500			
Beam separation (mm)	38	38			
Expander ratio	1	1.5			
Number of fringes	48	48			
Fringes spacing (µm)	7.005	4.674			
Beam half angle (deg)	2.176	3.262			
Probe volume dx (mm)	0.3389	0.2262			
Probe volume dy (mm)	0.3387	0.2258			
Probe volume dz (mm)	8.919	3.968			

Table 3-1. Laser Doppler Anemometry (LDA) settings

3.4 Flame Photography

Almost all combustion processes have undergone dramatic and fast temporal and spatial variations in their characteristics during very short periods of time, especially flame flashback. Thus high-speed photography is considered as a proper technique for investigation such combustion phenomenon. However, other imaging techniques even that of low f/s can give some beneficial diagnostics about flame propagation. In this work, Nikon D7200 digital camera has been used for flame imaging to detect some flame flashback diagnostics that can support the investigations. This camera characterised by the following characteristics;

- Resolution (24.2 megapixel)
- Sensor size (23.5 X 15.6 mm)
- ➢ Kit lens (27- 210 mm) zoom
- Viewfinder (optical/LCD)
- ➢ Native ISO (100-25600)
- ➢ 6 fps (continuous shooting)
- Extended ISO (100-25600)
- ➢ Shutter (1/8000-30 s)

3.5 Summary

The above measurement instrumentation can be considered as vital in terms of characterisation and verification of different combustion instability problems and in particular, for the understanding of flame flashback propagation.

The 150 kW tangential swirl burner has been used for investigating different combustion instability issues, especially blowoff and flashback. The configuration of this burner utilises working at wide swirl numbers with different arrangements. Moreover, the acquired data can be scaled up to be used with burners of highest output power.

Laser Doppler Anemometry (LDA) has the ability for investigating small control volumes and consequently provide an effective mean to have a deep insight and clear vision about the instantaneous velocity components and local turbulences that play an important role in flame flashback mechanisms.

CHAPTER 4

Characterisation of the Effect of Injector Geometry on Flashback Mechanisms

Chapter 4

Science walks forward on two feet, namely theory and experiment...But continuous progress is only made by the use of both.

Robert A. Millikan

4.1 Introduction

The global trend towards using alternative fuels that produce low harmful emissions has increased significantly. However, use of these new sources of energy faces many issues especially in terms of operational stability, with instability issues causing severe damages to the system hardware. Consequently, the motivation to develop new flexible combustion systems that meet the demands of switching to different fuels accompanied with high-reliability operation has augmented. However, development of these systems still faces many difficulties.

Swirl-stabilized combustion is the most widely spread deployed technology used to stabilise and control combustion in gas turbines and numerous other systems. However, the interaction of the swirling flows with the burner geometries is very complex, and it has been proved that any change in the burner geometry can affect the flow field inside the combustion chamber, close to the burner mouth and downstream the combustion zone. Thus, most burners are provided with a central fuel injector that centrally delivers well-known fuels allowing the stabilisation of the system previous to operation under entirely premixed conditions. Moreover, the injector anchors the central recirculation zone formed downstream of the nozzle [123].

However, swirl combustors can be subject to four different flashback mechanisms [65]. Recently flame flashback due to combustion induced vortex breakdown (CIVB) acquired particular importance because the possibility of its onset even during stable

operation of the combustor or in other words when the flow velocity is higher than the turbulent flame speed [35, 162]. Thus, the use of injectors can also affect the stability limits of the system, especially propagation of flashback through changes of shape of the shear layer. Generally, adapting the burner geometry to meet flashback safety requirements is of particular interest in this context [188, 189]. However, the characterisation of the flow and its impacts on the flame flashback propagation and other swirling structures using different injectors has been briefly documented. Therefore, this chapter focuses on obtaining data to support the study and demonstrate the deterioration of the CIVB with these injectors.

4.2 Experiments and Diagnostics

The tangential swirl burner described in chapter 3 has been used during the experiments. Seven central fuel injectors of different outside diameters (central fuel injectors; 7, 12.5, 16, 18, 19, 21.3, 23.5 mm) and the same length (175 mm) were used. Figure 4.1 shows the position of the injector inside the swirl burner. Experiments were done using all these injectors, nozzles and inserts. Figure 4.2 shows different nozzle configurations of 0.8D, 0.9D and a Quarl used to change burner exit diameter. By changing the exit nozzles configuration and inserts, variable swirl numbers from 0.91 to 3.65 have been achieved. Table 4.1 illustrates all experimental conditions.



Figure 4-1 Central fuel injector inside swirl burner.



Figure 4-2 Different nozzle configurations.

Type of insert	0.8D Nozzle	0.9D Nozzle	Quarl	Swirl number
25%	yes	-	-	0.91
25%	-	yes		1.3
25%	-	-	yes	1.12
50%	yes	-	-	1.59
50%	-	yes		1.80
50%	-	-	yes	1.98
70%	yes	-	-	2.96
70%	_	yes	_	3.34
70%	_	_	yes	3.65

Table 4-1 swirl numbers for the different nozzles and inserts used.

Two modes of fuel injection were utilised; a central injection mode where the fuel is injected via the central injector, and premixed injection where the fuel is injected through the tangential inlets just before the inserts, fuel always being natural gas.

To start the burner, firstly, a low fuel flow rate is injected axialy via the central fuel injector, and once the diffusive flame is established, low air flow rates are injected tangentially. Then a low amount of fuel is injected tangentially while central fuel

injection is reduced gradually until it is totally shut. At this point, a stable swirl flame is achieved under completely premixed conditions. From this point, flame flashback is achieved by increasing tangential fuel flow rate at constant tangential air flow rate, upon increasing tangential fuel flow rate the flame base propagates slightly to be in contact with the fuel injector for small size injectors, and in touch with the nozzle for big size injectors, the flame still in this position for a period at constant Φ , after that it propagates entirely inside burner plenum. Once the flame propagates against the incombustible mixture towards the burner plenum, tangential air and fuel flow rates are recorded.

The flashback is expected to be CIVB, whereas the minimum bulk velocity of the unburned mixture at low flow rates is (2 m/s) which is much higher than the laminar flame speed of methane (0.4 m/s). And it is not predicted the turbulence fluctuation is too high to make turbulent flame speed exceeds the flow velocity. The blowoff limits were measured by gradually reducing the tangential fuel flow rate, but maintaining the air constantly flows through the tangential inlets. The blowoff point was determined when the flame zone was visibly lifted from the burner mouth and blown off. The stable operation region is to the right of blowoff points, and to the left of flashback points. In order to check the accuracy of results, experiments were repeated five times, showing deviations not greater than 5%. This method is repeated for another flow rate values (higher than the previous one) until obtaining all flashback and blowoff results for each fuel injector at a certain swirl number. All experiments for other swirl number is repeated by the same procedure.

4.3 Results

Flashback trends are correlated according to the inlet tangential velocity where using this correlation allows a comparable analysis between cases and reveals the effects of outside central fuel injector diameters with different types of swirl burners. Moreover, this approach allows a fair comparison between cases without considering the change of Re at the burner mouth caused by the different injectors [185]. To make the data more generalised, a dimensionless number (χ) which represents the ratio of the outside injector diameter to the inside nozzle diameter has been used for the analyses. Those values are illustrated in Table 4.2.

	Ĩ		
Injector type	0.8D Nozzle	0.9D Nozzle	Quarl
OD in mm	D=60.8 mm	D= 68.4mm	D=75.0 mm
7.0	0.115	0.102	0.090
12.5	0.205	0.182	0.160
16.0	0.263	0.233	0.213
18.0	0.290	0.260	0.240
19.0	0.312	0.270	0.253
21.0	0.345	0.307	0.28
23.0	0.378	0.336	0.306

Table 4-2 Values of fraction χ (Injector OD/Nozzle inside diameter)

Figure 4.3 shows the flashback curves for different χ using a swirl number of 1.12 (Quarl Nozzle). It can be seen that with the reduction of χ the flashback curves move to leaner regions in the inlet tangential velocity range (W_{in}) of 2.0-3.5 m/s over an equivalence ratio (Φ) ranging from 0.50 to 0.75.



Figure 4-3. Effects of central fuel injector outside diameter on flashback limits, Quarl, swirl number 1.12.

Noteworthy, considering the flashback results of χ =0.09 and χ =0.016, as the trends follow one curve and by visual observation, it appears that central recirculation zones extend over the central fuel injector towards the baseplate for all tangential inlet velocities. The flashback mechanism can be seen occurring radially from the outer boundary of the CRZ to the radial aligned tangential inlets. This type of flashback has been previously denoted by partial flashback [190]. Flashback resistance under these conditions is weak because it depends on the relatively low radial velocity in the flat swirl chamber. The effect of the injector diameter has an almost little effect regarding CRZ formation.

However, as the diameter of the injector is increased, flashback limits generally improve (providing wider burner stable operation) until χ =0.253 where flashback equivalence ratio has improved to ~ 0.6 over a wide range of tangential velocities. The CRZ/burning vortex region still extend around the fuel injector but not always to the burner baseplate and the flashback mechanism still appears to be a radial flashback in the flat swirl chamber. Upon further increase of the injector diameter to χ =0.280 and χ =0.306 flashback resistance increases, while the location of the flame front and flame stabilisation mechanism change significantly.

The CRZ is almost unstable at low tangential velocities, close to vortex breakdown, leading to flame fluctuation and an early flashback. However, higher tangential velocities stabilise the flow field, and the flashback mechanism has changed. This occurs via the outer boundary layer between the shear layer and the nozzle wall. Now two flames are formed, one stabilised on the central fuel Injector while the other on the outer lip of the burner. Figure 4.4 shows the transition of flashback mechanisms when injector diameter is increased.



Figure 4-4. Flames close to a flashback in Premixed Swirl Burners with central fuel Injectors; A) Flame burns in the boundary of the CRZ which extends down over the fuel injector to the baseplate; B) Two flames formed, one stabilised by the central fuel injector, the other located on the burner lip.

It is worth mentioning here that although working at leaner conditions is important for lean premixed combustion systems, movement of flashback curves to the leaner region for the case of small injectors leads to a considerable reduction of stability map (the blowoff and flashback curves became close to each other) hence more likely to subject to blowoff and flashback, as can be seen from figure 4-5. This narrow operation area is not favourable regarding the swirl combustors safe operation, any small change in equivalence ratio or mass flow rate could lead to one of the instability cases.



Figure 4-5. Small injector diameter leads to considerable reduction the of the stable operation region; big diameter injectors keep wider stability operation region. Quarl, swirl number 1.12.

The change of annular passage area of the flow between the central injector wall and inside wall of the nozzle plays an important role in this context. At small values of χ (small central injector diameter) a large annular passage area is available. Thus adverse pressure gradient across the combustion zone due to sudden density rise across the flame front and hence the static pressure generated at the flame tip leads to a boundary layer separation from the wall of the injector, consequently leading to upstream flame flashback propagation. However, the interaction with the PVC may keep the flame close to the injector wall and extend to the baseplate [146]. The significant pressure difference generated due to the higher jump in density may also lead to increase the baroclinic torque which in turns initiate negative velocity regions at the flame tip leading to upstream flame flame flame propagation [144, 191].

However, upon increasing the central injector diameter, the passage area is reduced, and hence the unburned mixture will have a slightly higher momentum which provides a better impedance to stagnation pressure and hence reduced boundary layer separation and consequently upstream flame flashback propagation. Moreover, the increase or reduction of the injector diameter at certain equivalence ratio can affect the flow field velocity at the burner exit, hence the balance between the flow velocity and flame speed. Consequently the flame flashback [192]. Figure 4.6 shows the difference in flashback conditions based on the burner exit velocity (directly downstream the nozzle) between a medium size injector diameter (χ =0.213) and a large size injector diameter (χ =0.306). It appears that for big diameter injectors flashback occurs at slightly higher exit velocities than for the small diameter injector at approximately the same mass flowrates. This suggests that small annular passage areas provide better flashback resistance by providing the required matching between the axial flow velocity and flame speed, as expected.



Figure 4-6. Effect of annular passage area on flashback limits, Quarl, swirl number 1.12.

Moreover, the propagation and extension of the CRZ around the central injector at small χ values can lead to a considerable amount of heat dissipation. Whereas this can, in turn, degrade the ability of the CRZ to recirculate the heat and active chemical species to the root of the flame which is one of the important features of swirling flames stabilisation. Moreover, long-term exposure of the central fuel injector for this

heat, especially in a practical operation like that in stationary gas turbines, will lead to increase maintenance cost requirements because it has a negative impact in terms of the central injector life spam [108] in addition to the possibility of increased pollutant levels.

For lower swirl numbers 0.91, the flashback trends are still almost the same, Figure 4-7. It appears that flashback curves located in tangential velocity ranges (Win) of 2.30-4.30 m/s over an equivalence ratio (Φ) range from 0.50 to 0.70. The flashback limits of the low injector diameter ranges (χ = 0.115, 0.205 and 0.263) are located in the same equivalence ratio while increasing injector diameter to χ = 0.345 and 0.378 gives better flashback trends in terms of stability operation. These results were evident in all the analysis done for all swirl numbers and geometries which meaning that the mechanism of flashback is similar for similar injector sizes. Interestingly, a critical of χ is around 0.280 for S=1.12 and 0.320 for S=0.91. This suggests that flashback resistance is improved with lower swirl numbers.



Figure 4-7. Effect of central fuel injector diameter on flashback limits, 0.8 D nozzle, swirl number 0.91.

However, in general, the flashback limits for this configuration occur at a slightly higher tangential velocity than the previous one, Figure 4-8.



Figure 4-8. Comparison of flashback limits for the two swirl numbers.

Fig 4-8 shows the flashback limits for two different central injector diameters at two different swirl numbers. The sudden flow expansion for 0.8D Nozzle,(S=0.913) can provide better pressure gradient conditions at the burner lip. For this condition the CRZ can be accompined by an outer recirculation zone (ORZ) in addition to two annular shear flow regions, the inner shear layer (ISL) located between the high momentum annular flow region and CRZ, while the other is the outer shear layer (OSL) between the high momentum annular flow region and ORZ. In this configuration the flame stabilises on both the ISL and OSL, this aerodynamic stabilisation can increase flame resistance to flame flashback propagation. Nonetheless, other factors such as local straining of the flame by a high shear layer may be consider as a crucial factor in flame stabilisation in this context [108].

However, when using the Quarl nozzle, (S=1.12) flashback occurs slightly earlier in terms of swirl inlet tangential velocity, as the flow expansion becomes rather gradually compared with a sudden expansion case. Thus, the flame stabilisation by the outer shear layer becomes less pronounced , hence less aerodynamicly stable, and consequently, more prone to flashback at higher inlet tangential velocities.

4.4 Summary

Development and modification of combustion systems geometries are of special importance to achieve many requirements regarding enhancement of gas turbine performance. One of this demands is enabling swirl combustors to work at wider stability margins especially upon introducing alternative fuels.

Thus, this chapter presents an experimental approach for the effect of changing the outside diameter of the central fuel injector in swirl burner on flame flashback propensity, especially combustion induced vortex breakdown flashback (CIVB).

A dimensionless number (χ) which represent the ratio between the nozzle inside diameter and the central injector outside diameter has been used to correlate results. This correlation makes the result more global and can be applied regardless the burner size. Results showed that this geometrical variation can affect the flashback mechanism significantly. Reduction of the central injector outside diameter moves the flame flashback limits to the leaner region. However, there is a limit for this reduction. Flashback results for small size injectors give almost the same result and in general very small injector diameter gives result to those without injector.

A considerable alteration in flashback mechanism has been observed in terms of CRZ stabilisation, when using small diameter injectors generally the CRZ extends around the central fuel injector towards the burner baseplate, this upstream, CRZ propagation is considered as an undesirable flame propagation because it causes a considerable rise in temperature around the injector which in turns lead to increase pollutant levels accompanied with reduction of injector life and increase maintenance cost.

However, wider injector diameters provide better flashback resistance through providing better CRZ stability. The flow stagnation point is shifted downstream because of the slightly increase in the axial flow velocity due to the reduction in the annular passage area. CRZ generally stabilizes on the inner and outer shear layers downstream the burner nozzle and injector.

The results for the two swirl numbers are almost identical despite some differences in terms of working at a slightly higher inlet tangential velocities for S=0.913. Moreover, a critical value of χ for a low swirl number S=0.91 is χ = 0.320 which is higher than that for other swirl number S=1.12, i.e. χ = 0.280. The reason for this effect is low

swirl numbers result in weak CRZ, hence less likely for CIVB propagation. Consequently, the transition from CIVB to BLF delayed to higher equivalence ratios and enables working at higher flowrates, contrary to higher swirl numbers, i.e. S=1.12 as the CRZ is bigger, hence more likely for CIVB propagation. As a result, early transition to CIVB in terms of equivalence ratios is observed. It can be concluded that flashback resistance is enhanced at low swirl numbers. Thus, determining this value is of high importance in terms of burner design and hence operation stability requirements.

Results evidence that using a central fuel injector as bluff body can significantly annihilated the CIVB flashback to some extent. Nonetheless, when choosing central fuel injector it is important to take into consideration that it is big enough to provide central fuel for stability purposes and at the same time small enough to destroyed CIVB.



Effect of Axial Air Injection on Flame Characteristics and Flashback Resistance

Chapter 5

The science of today is the technology of tomorrow. Edward Teller

5.1 Introduction

Swirl combustors have been proven as effective flame stabilisers over a wide range of operation conditions thanks to the formation of well-known swirl coherent structures which provide low-velocity regions that enable flame anchoring. However, the interaction between swirl structures incoming flow conditions and swirl burner geometries is so complicated that any change in the later can considerably alter the stability regime downstream the burner exit plane.

Using central injectors either as a central body or to inject fuel axially has been used successfully to achieve good stability limits through transient conditions and hence preventing upstream flame propagation. However, injecting either fuel or mixtures of air and fuel axially can degrade the degree of mixing, consequently increasing pollutant levels significantly [87]. Furthermore, employing bluff- bodies in gas turbines does not achieve whole safe operation with risks subjected to different flashback mechanisms including combustion induced vortex breakdown (CIVB) [35, 193]. In addition, centre bodies, bluff-bodies and central fuel injectors can undergo material degradation due to harsh environments produced by high flame temperatures caused by the surrounding of the central recirculation zone (CRZ) especially when high hydrogen content blends are used [73, 194]. Figure 5.1 shows the effect of high temperature on central fuel injectors.

Air injection can provide many solutions for different operational and maintenance problems in the gas turbine sector. For instance, this technique can replace the use of central or bluff-bodies which frequently are subjected to overheating conditions. Moreover, air injection allows wide operation stability margins compared to stationary fuel injectors since the diameter of air jets can be varied during combustor operation to meet stability requirements. Thus, this chapter describes the effect of using axial air injection on flow field characteristics and how it can affect the lower instability limits through altering the flashback mechanism produced by the induced vortex breakdown flashback (CIVB).



Figure 5-1. Effects of high temperature on Central Fuel injectors under flame flashback Conditions

5.2 Apparatus and Experimental Procedure

5.2.1 Air Injection Facility

The description of experimental setup has been mentioned in chapters 3 and 4. However, a new configuration has been utilised to enable central air injection from the burner plenum bottom instead of the central fuel injector. Thus the original baseplate which contains the central fuel injector (Central injector) has been replaced by a new one that allows axial air injection in addition to the fuel, Figure 5.2. The inside diameter of the central air injector has been chosen to be 19 mm based on previous results from chapter 4 where this diameter produced the region of transition of the behaviour of flame flashback from wall boundary layer flashback (BLF) to combustion induced vortex breakdown flashback (CIVB). Figure 5.3 illustrates this considerable change of flame flashback mechanism due to the change of central fuel injector diameter.



Figure 5-2 Experimental setup.



Figure 5-3. variations of flame flashback mechanism with injector outside diameter.

The air injector is fitted by external screws inside a cylindrical pipe connected to the burner baseplate. This allows its vertical movement inside the burner plenum to have different positions (X) with respect to the baseplate. Figure 5.4 shows the position of the air injector inside the burner plenum.



Figure 5-4. Different air injector positions inside Burner plenum.

Six different positions of the air injector inside the burner plenum have been investigated. Table 5.1 illustrates those positions.

No	X (mm)	From outlet (mm)	notes
1	0	205	Air injection directly from burner baseplate
2	29	176	Downstream baseplate
3	48	157	=
4	75	130	=
5	110	95	=
6	150	55	=

Table 5-1 Ai	r iniector	positions	with	respect	burner	baseplate.
Table 3-1 Al	I mjector	positions	WILLI	respect	Durner	vaseplate.
5.2.2 Obtaining Burner Stable Operation

The process of obtaining stable swirl flames without central fuel injector in relatively high power laboratory tangential swirl burners is difficult, especially when the air is injected at X=0 (air injection position parallel to the burner baseplate). The absence of a bluff body complicates the mechanism of flame anchoring and hence central recirculation zone (CRZ) generation. Moreover, the difficulty includes starting the swirl burner without central fuel injection. Thus, many experiments have been done to determine the proper procedure to start the burner and have stable swirl flames. The first set of experiments was implemented at low flow rates to avoid the risk of subjected to sever flashback that could lead to damaging of the system. The bestobtained procedure for burner operation under these conditions was to inject fuel firstly. This was done by shutting the air valve and enabling fuel injection through the air injector pipe. Then the tangential premixed injection was started at low flow rate with simultaneous decreasing of the central fuel injection. Once a stable flame is achieved the central fuel was shut down. At this point the axial air injection can be started with simultaneous increasing of the tangential premixed mixture until having a stable swirl flame was observed.

However, the appropriate amount of air injected axially is also crucial in terms of the ability of this axial jet to physically simulate not only the shape of the central fuel injector but also act in the same way as a solid body. Thus, different amounts of axial air injection at different tangential premixed mixtures have been investigated. The most suitable axial air flowrate was 50 LPM, this value represents almost (3 - 10 %) from the total volume flow rate of the burner.

In order to obtain stability limits, firstly blow off limits were determined. At stable operation condition for a certain flowrate the tangential gas was reduced at constant tangential and axial air injection until the flame was totally blown off. This procedure was repeated for all flow rates. Flashback limits were evaluated using the same procedure except the tangential gas was increased at constant tangential and axial air injection until the flame at constant tangential and axial air injection until the flame was increased at constant tangential and axial air injection until the flame propagated upstream and into the plenum.

5.3 Results and Discussion

5.3.1 Burner Stability Operation Map Using the Central Air Injection

Compared to actual injectors, central air injection affects the flashback trends significantly. The use of the central air injection simulates the physical shape of a central fuel injector allowing wider operation ($0.55 > \Phi < 0.7$) over a tangential velocity range from 2.5 m/s to 7.5 m/s. While when using the central fuel injector as central body the burner stability operation is narrower at lower equivalence ratios ($0.48 > \Phi < 0.57$) over a tangential velocity range 2.5 - 4 m/s after which no stable flame can be achieved. Figure 5-5 shows the different stability limits of the burner when using air stream injected from the bottom of the burner baseplate instead of a central fuel injector. The stable operation region is to the right of the blowoff points, and to the left of the flashback points. In order to check the accuracy of results, experiments were repeated five times, showing deviations not greater than 5%.



Figure 5-5. Difference in the stability operation when the air stream is injected from the baseplate (X=0) and the use of a central injector, 0.8D nozzle, S = 0.9.

This impact of air injection has very important potentials regarding increasing the output power and hence the total efficiency of the combustor, as seen from figure 5.5. Using air stream enables the burner to operate at higher tangential velocity values (almost three times greater than the case when using a central fuel injector to promote

the stability operation) and higher equivalence ratios which in turn increase the overall stability operation map.

Moreover, it appears that both blowoff and flashback limits of the air injection case are located at relatively more constant equivalence ratios at variable tangential velocities. This finding is important because flashback due to CIVB can initiate due to an increase in equivalence ratio or mass flow rate fluctuation [160].

Additionally, this can reduce combustion instabilities that can arise from equivalence ratio fluctuations, while this fluctuation is evident when the central fuel injector is used. Moreover, it is possible to achieve the reasonable amount of tangential velocity at relatively constant equivalence ratios while switching to a higher power and maintaining constant equivalence ratios.

This effect of axial air injection on the burner stability map is still achievable at higher swirl numbers (S = 1.2, S = 1.59 and S = 3.65) as can be seen from Figures 5.6, 5.7 and 5.8, which suggests the validity of using this flame stabilisation technique at even higher swirl strength burners.



Figure 5-6. Difference in the stability operation when the air stream is injected from the baseplate (X= 0) and the use of a central injector, Quarl, S = 1.12.

The central air injection promotes flame stability by affecting the aerodynamic characteristics of the flow field downstream the burner mouth. It reduces the defect in the axial velocity at the tip of the recirculation zone which is one of the main reasons leading to CIVB flashback. Figure 5.9 illustrates the effect of central air injection on

the axial velocity at the central axis if the burner is used for variable flow rates under isothermal conditions. All flow rate values of isothermal tests are illustrated in appendix A



Figure 5-7. Difference in the stability operation when the air stream is injected downstream baseplate (X= 29 mm) and the use of a central injector, 0.8 D nozzle, S = 1.59.



Figure 5-8. Difference in the stability operation when the air stream is injected downstream the baseplate (X= 29 mm) and the use of a central injector, Quarl, S = 3.65.



Figure 5-9. LDA results, effect of axial air injection on the defect of axial velocity downstream burner mouth X= 0 mm, Y = 5mm (Isothermal condition).

This effect is almost the same at downstream levels Y = 15 mm and Y = 25 mm as can be seen from figures 5.10 and 5.11



Figure 5-10. LDA results, effect of axial air injection on the defect of axial velocity downstream burner mouth. X= 0 mm, Y = 15mm (Isothermal condition).

The effect of axial air injection on the centreline velocity defect became more evident at low flow rates as the air injector was moved downstream the burner baseplate. As can be seen from figure 5.12, with air injector position being at X = 29 mm and figure 5.13 with air injector is located at X = 150 mm. It is worth mentioning here that the closer the air injection position to the nozzle exit, the more symmetry of the velocity profile. This change of symmetry of swirl flow is because the axial air jet is considerably affected by inlet tangential jets at X=0 mm, whereas this effect is almost disappeared at X=150 mm, as the air axial air jet became entirely protected by injector body and the sleeve.



Figure 5-11. LDA results, the effect of axial air injection on the defect of axial velocity downstream burner mouth X = 0 mm, Y = 25mm (Isothermal condition).

However the change in axial velocity profile is almost the same at high flow rates as can be seen from figure 5.14, this being a consequence of the weak CRZ at low flow rates.



Figure 5-12. LDA results, the effect of axial air injection on the defect of axial velocity downstream burner mouth, X = 29 mm, Y = 5 mm (Isothermal condition).



Figure 5-13. LDA results, the effect of axial air injection on the defect of axial velocity downstream burner mouth X = 150 mm, Y = 5 mm (Isothermal condition).



Figure 5-14 LDA results, effect of axial air injection on the defect of axial velocity downstream burner mouth at high flow rates X = 150 mm, Y = 5 mm (Isothermal condition)

However, unlike the isothermal case, swirling flows under combustion conditions yield dramatic changes. The recirculation bubble formed at the tip of the recirculation zone can move further upstream under the effect of the heat generated that plays an important role in the interaction between volume expansion and baroclinic torque which in turn represent the crucial factors that promote flame stability in this region.

Moreover, the pressure variations inside the burner plenum arise due to chemical reactions downstream the nozzle. This can be another reason for this upstream shift in the position of low or negative velocity regions [65, 162].

It is worth mentioning here that these dramatic changes in the flow field are not taking place under isothermal conditions. Thus the axial velocity profile can change significantly, although some differences in total mass flow rate between two cases, as in the combustion case fuel mass flow rate is added. Figures 5.15 and 5.16 show the difference in axial velocity profile between the two cases. This difference has also been demonstrated previously by [163].

Under combustion conditions the effect of axial air injection on negative velocity defect still the same as that of isothermal conditions for low flow rates (400 and 600 LPM). However, at a higher flow rate 800 LPM the effect is altered, this evinces of propagation of the CRZ. The axial air injection pushes the CRZ down, and its tip became parallel to the measurement level, while without air injection CRZ propagates inside the sleeve. Hence measurement level became bounded by CRZ in less negative velocity regions than that at the tip of the CRZ. Figure 5.17 illustrates this difference. All flow rate values of combustion tests are illustrated in appendix B.



Figure 5-15. LDA results, different axial velocity profiles isothermal and combustion conditions at X= 0mm Y= 5 mm.



Figure 5-16 LDA results, different in axial velocity profile between isothermal and combustion at X= 29 mm Y= 5 mm.



Figure 5-17 LDA results, the effect of axial air injection on the defect of axial velocity downstream burner mouth X = 150 mm, Y = 5 mm (Combustion condition).

5.3.2 Effect of Central Air Injection

Although injecting air axially from the burner baseplate enhances the stability map significantly, the position of the central injector inside the burner plenum with respect to the tangential inlets and burner nozzle has a particular importance regarding the flow regime inside the combustor and hence the overall stability operation. Figure 5.18 shows the flame flashback trends for six different positions of the central air injector with respect to the burner baseplate. When air is injected directly from the burner baseplate X = 0, the flame flashback margin is around 0.7 equivalence ratio over inlet tangential velocities ranging from 2.5 to 7.5 m/s.

However, when the position of the axial air injector opening is parallel to the bottom edge of the tangential inlets X=29 mm, flashback trends are affected significantly and shifted to the leaner region. This is mainly because the axial air jet is subjected directly to the high momentum tangential flow which in turns produces a considerable pressure fluctuation, hence high turbulence intensity in this region as will be explained in section 5.3.3 leading to flashback occurring at lower Φ .

Further downstream at X=48 mm the flashback behaviour is similar to that of X=0 case at low flow rates of velocity at 2-3 m/s. However, upon increasing flowrates the overall turbulence generated at the bottom of the burner sleeve produces an aerodynamic flow fluctuation leading to earlier flashback compared with that of X=0, and little bit richer than X=29 mm case.

When the air injection is located further downstream inside the burner sleeve at X=75 mm, the flame flashback trends recover and simulate the initial case at X=0 mm despite some difference in the tangential velocity in the range of 3-4 m/s. This behaviour is linked to the position of the air injector as it is protected from fluctuations generated at the sleeve bottom.

Enhancement is observed at X=110 mm and stable operation became wider whereas in this region both the air injection and burner physical body (central body) promote flame stability. These stability limits are almost the same when the air injector opens at X=150 mm. However, flashback trends move to a richer region at higher flow rates (5 m/s tangential velocity), and no flashback is observed at higher flow rates, as can be seen from figure 5.19. This outcome is of particular importance regarding possibility of switching to a higher power operation.



Figure 5-18 flame flashback trends at different positions of central air injection S = 0.9.



Figure 5-19 Burner stability map, X=150 mm.

At a slightly higher swirl strength, S = 1.12 (when the Quarl nozzle is used) the effect of air injector position is nearly the same as that of, S = 0.9 case. However, the flashback trends of X = 0 mm and X = 29 mm have improved, moving to higher equivalence ratios compare with that of low-swirl case. The sudden expansion of burner exit nozzle reduces the local pressure at the exit plane. Consequently the region of high turbulence fluctuations shifted downstream the bottom of tangential moving from X = 29 mm to X = 48 mm, figure 5.20 illustrates the burner stability map when a Quarl nozzle is used. Using Quarl nozzle (sudden divergent) leads to decrease local pressure downstream burner mouth. Consequently the position of high turbulence region is shifted from X=29 mm to X=48 mm.



Figure 5-20. Flame flashback trends at different positions of central air injection S = 1.12.

5.3.3 Effect of Turbulence Intensity on Flame Flashback Propensity

The mutual relation between combustion, heat transfer and turbulence is of high importance in determining operation stability figures of any combustion system. Thus understanding the interaction between those parameters is considered a crucial factor in solving many combustion instability issues. In turbulent flames, the high level of turbulence can increases the flame speed, consequently the possibility of flashback initiation [152, 195, 196]. The relation between turbulence intensity and turbulent flame speed is described as follows [35];

$$S_T \propto S_L + u' \tag{5-1}$$

Where

 S_T : turbulent flame speed [m/s]

 S_L : laminar flame speed [m/s]

$$u'$$
: velocity fluctuations [m/s]

According to equation 5-1, any increase in turbulence intensity will consequently be followed by an increase in turbulent flame speed. Hence if this increment occurs at some weak regions inside the swirling flow, especially in the tip of recirculation bubble or the recirculation zone CRZ, there is a strong possibility of upstream flame propagation caused by turbulence effects in the flame. Thus determining and correlating turbulence intensity with combustion instabilities, especially flame flashback has significant potentials. However, methods of measuring turbulence intensity mainly depend on the flow characteristics.

Swirl flows are three-dimensional time-dependent in nature. Thus they reveal dramatic changes in their characteristics from one to another adjacent position inside the flow. Therefore, such flow is considered as anisotropic [197]. Thus, in this study two methods have been used to determine the turbulence intensity based on one - dimensional LDA axial velocity measurements at different distances downstream the burner mouth.

The first method can describe the instantaneous changes in turbulence intensity across the swirl flows, especially in the shear layers, high momentum flow region and CRZ, and consequently, correlate those changes with combustion instabilities. Turbulence intensity in this method is calculated according to the following equation:

$$T_u = \frac{u_{rms}}{\bar{u}} \tag{5-2}$$

Where;

 u_{rms} : the root mean square value in the axial direction measured by LDA

 \overline{u} : mean axial velocity measured by LDA

This method will be denoted by method A.

While for the second method turbulence intensity is determined by the following equation:

$$T_u = \frac{u_{rms}}{U_b} \tag{5-3}$$

 U_b : bulk flow velocity calculated by dividing the volume flow rate of air-fuel mixture at the burner exit nozzle area. This method will be denoted by method B.

5.3.3.1 Effect of Turbulence Intensity on Flow Field Characteristics, (Isothermal Test, Method A)

Swirling flows undergo more complexity under combustion conditions. Thus the first set of measurements has been done under isothermal conditions and then compared with the combustion case. Experiments have been implemented to investigate different flow rates at different conditions (with and without central air injection, different air injector positions X and different distances downstream the burner mouth Y). Figure 5.21 shows the effect of central air injection on turbulence intensity values

(Method A).



Figure 5-21. LDA results, the effect of central air injection on turbulence intensity values under isothermal conditions (Air flow rate 400 LPM, X = 150 mm, Y = 5 mm, method A).

As can be seen from figure 5.21, central air injection decreases the level of turbulent intensity considerably. The turbulent intensity values fall by almost 45% than its original values. This effect is of great interest in terms of overall flow stability. The peaks of turbulence intensity values indicate the presence of shear layers that are very close to the burner centre [152], which suggests that this testing position (Y = 5 mm) is almost at the tip of the recirculation zone where shear layers became close to each other around the burner centre. Hence this represents a proper position for investigation.

Although this high turbulence can improve mixing and promote stability especially at low swirl numbers, it can initiate flame flashback under certain conditions, as the adverse pressure gradient generated by the swirl leads to significant variations in the level of turbulence especially at the burner centerline which in turn provokes CIVB flashback. However, introducing the central air injection eliminates turbulence variations producing more uniformity in the velocity gradients. This effect is observed at different flowrates as can be seen from figures 5.22 and 5.23.



Figure 5-22. LDA results, the effect of central air injection on turbulence intensity values under isothermal conditions (Air flow rate 600 LPM, X = 150 mm, Y = 5 mm, method A).



Figure 5-23. LDA results, the effect of central air injection on turbulence intensity values under isothermal conditions (Air flow rate 800 LPM, X = 150 mm, Y = 5 mm, method A).

However, this effect in the reduction of turbulence is less acute or disappears at slightly higher flow rates as can be seen from figures 5.24 and 5.25. It seems to be a consequence of the elongation of the CRZ, as in Figures 5.12, 5.13 and 5.14.

For low tangential flow rates, the ratio of central air injection is around 9%, while at higher flow rates it is reduced to about 5%. Thus injection of 10 % of air axially seems to be reasonable to achieve the movement of the CRZ further away from the outlet of the burner.



Figure 5-24 LDA results, the effect of central air injection on turbulence intensity values under isothermal conditions (Air flow rate 1000 LPM, X = 150 mm, Y = 5 mm, method A) Central air injection effect reduces at high flow rates.



Figure 5-25 LDA results, effect of central air injection on turbulence intensity values under isothermal conditions (Air flow rate 1200 LPM, X = 150 mm, Y = 5 mm, method A) Central air injection effect reduces at high flow rates.

The effect of high tangential flow rates on the ability of using central air injection has been observed at almost all conditions at different air injector positions (X = 110 mm, X= 29 mm and X = 0 mm), as can be seen from Figures 5.26, 5.27,5.28, 5.29, 5.30, 5.31, 5.32, 5.33, 5.34 and 5.35.

However, for X= 29 mm and X = 0 mm the effect almost disappears at slightly moderate tangential flow rate values (800 LPM). This due to high turbulence generated in this region which arises when the air injector opening is set to be parallel to the tangential inlets.

However, turbulence at 1000 LPM at the centre shows a reduction. Since shear layers show high turbulence, it is likely that the CRZ has already entered the sleeve. The same happens with X=0. The results suggest that the resistance of the incoming flows to the movement of the CRZ has been considerably reduced as a consequence of high turbulence produced at the tip of the central injector, or that the turbulent speed at the tip of the CRZ has been increased. Thus promoting the movement of the coherent structure upstream the measuring point, beyond which the air injection start effect again. At x= 150 and 110 mm the air injector yields less turbulence because it is protected by the sleeve.



Figure 5-26.LDA results, the effect of central air injection on turbulence intensity values under isothermal conditions (Air flow rate 800 LPM, X = 110 mm, Y = 5 mm, method A).



Figure 5-27. LDA results, the effect of central air injection on turbulence intensity values under isothermal conditions (Air flow rate 1000 LPM, X = 110 mm, Y = 5 mm, method A).











Figure 5-30. LDA results, the effect of central air injection on turbulence intensity values under isothermal conditions (Air flow rate 800 LPM, X = 29 mm, Y = 5 mm, method A), Central air injection effect reduces at moderate flow rates.



Figure 5-31. LDA results, the effect of central air injection on turbulence intensity values under isothermal conditions (Air flow rate 1000 LPM, X = 29 mm, Y = 5 mm, method A), Central air injection effect recovers at high flow rates.



Figure 5-32. LDA results, the effect of central air injection on turbulence intensity values under isothermal conditions (Air flow rate 400 LPM, X = 0 mm, Y = 5 mm, method A).







Figure 5-34. LDA results, the effect of central air injection on turbulence intensity values under isothermal conditions (Air flow rate 800 LPM, X = 0 mm, Y = 5 mm, method A), Central air injection effect reduces at moderate flow rates.



Figure 5-35. LDA results, the effect of central air injection on turbulence intensity values under isothermal conditions (Air flow rate 1000 LPM, X = 0 mm, Y = 5 mm, method A), Central air injection effect recovers at high flow rates.

Turbulence intensity values vary considerably upon changing the air injector position inside the plenum; this variation has crucial effects regarding burner stability and the mechanism of upstream flame propagation during flashback propensity. At low tangential flow rates (400 LPM), injecting air directly from the burner baseplate reveals higher turbulence intensity than the other two positions as can be seen from Figure 5.36.



Figure 5-36. LDA results, the effect of air injector position on turbulence figures, 400 LPM tangential flow rate with central air injection, method A.

However, at a slightly higher tangential flow rate (600 and 800 LPM) the effect has changed. X=29 mm position reveals higher turbulence intensity among other positions as can be seen from figure 5.37 and figure 5.38. Moving to more higher flow rate (1000 LPM) turbulent intensity values became higher at X=110 and 150 mm than other two positions as can be seen from figure 5.39. The results show how the CRZ moves at different flow rates with different injectors, a consequence of different turbulence at the tip of the CRZ structure.

Although determining turbulence intensity according to method A (based on local axial velocity) can give a reasonable prediction about CRZ propagation, there are some drawbacks of using this method. As the turbulence intensity could become infinity as local axial velocity tends to be zero, thus correlation the results based on turbulence intensity determined based on bulk flow velocity seems to be crucial.



Figure 5-37. LDA results, the effect of air injector position on turbulence figures, 600 LPM tangential flow rate with central air injection, method A.



Figure 5-38. LDA results, the effect of air injector position on turbulence figures, 800 LPM tangential flow rate with central air injection, method A.



Figure 5-39. LDA results, the effect of air injector position on turbulence figures, 1000 LPM tangential flow rate with central air injection, method A.

5.3.3.2 Effect of Turbulence Intensity on Flow Field Characteristics, (Isothermal Test, Method B)

Turbulence intensity measured according to bulk flow velocity gives the overall turbulence profile, correlating root mean square velocity fluctuations of swirling flows with bulk flow velocity. Results can be compared with other combustors regardless differences in configuration or power output. Thus turbulence intensity figures of this method have been represented according to the relation between (u_{rms}/U_b) vs. (r/R_0) .

Figure 5.40, shows the effect of air injection on the turbulence intensity values at low flow rates (Method B). It appears that the central air injection reduces the amount of turbulent intensity by about 25%. Moreover, a variation of turbulence intensity in the shear layers became more uniform and less intense. Higher values of turbulence intensity is located at the boundary of the central recirculation zone. Those values may have crucial effects on the propagation of the CRZ previous to the CIVB flashback. Thus by tackling this high turbulence fluctuation around the CRZ, it is possible to govern some flashback mechanisms specially CIVB.



Figure 5-40. LDA results, the effect of central air injection on turbulence intensity values under isothermal conditions (Air flow rate 400 LPM, X = 150 mm, Y = 5 mm, method B).

This effect is also acute at slightly high flow rates, as can be seen from Figures 5.41, 5.42, 5.43 and 5.44. Nevertheless, this effect is reduced with increasing tangential flow rates as the CRZ move to the measuring position.



Figure 5-41. LDA results, the effect of central air injection on turbulence intensity values under isothermal conditions (Air flow rate 600 LPM, X = 150 mm, Y = 5 mm, method B).



Figure 5-42. LDA results, the effect of central air injection on turbulence intensity values under isothermal conditions (Air flow rate 800 LPM, X = 150 mm, Y = 5 mm, method B).



Figure 5-43. LDA results, the effect of central air injection on turbulence intensity values under isothermal conditions (Air flow rate 1000 LPM, X = 150 mm, Y = 5 mm, method B).



Figure 5-44. LDA results, the effect of central air injection on turbulence intensity values under isothermal conditions (Air flow rate 1200 LPM, X = 150 mm, Y = 5 mm, method B).

The effect of the position of air the injector on turbulence intensity figures is also evident when using this method. Figures 5.45 and 5.46 show this effect at low flow rates 400 LPM and 600 LPM respectively.



Figure 5-45 Effect of air injector position on turbulence figures, 400 LPM tangential flow rate with central air injection, isothermal conditions, method B.

It appears that as the air injector opens closer to the burner baseplate e.g. X = 0 and X=29 mm, turbulence increases considerably. This observation has also been evinced in figure 5.36 method A. For this reason at these positions no stable flame was obtained at 400 LPM flow rate. However, they have almost the same flashback trends at 600 LPM (air tangential flowrate of 2.8 m/s at $\Phi = 0.65$).



Figure 5-46. LDA results, effect of air injector position on turbulence figures, 600 LPM tangential flow rate with central air injection, isothermal conditions method B.

For the third position, since it has lower turbulence intensity values, it enables the burner lit at low flow rates at almost 1.6 m/s inlet tangential velocity and increases the stability margin at slightly higher flowrates (600 LPM, inlet tangential velocity 2.8 m/s). The flame flashback is delayed to $\Phi = 0.7$ compared to the previous two positions as can be seen from figure 5.18. For higher flowrates 800 and 1000 LPM, figures 5.47 and 5.48 respectively. Despite some increase in turbulence intensity at X= 110 & 150 mm positions compared with those of X=0 & 29 mm, they still produce high flashback resistance due to a combined effect of both air injection and bluff body effects, which is not available for X= 0 and X = 29 mm positions.



Figure 5-47. LDA results, the effect of air injector position on turbulence figures, 800 LPM tangential flow rate with central air injection, isothermal conditions, method B.



Figure 5-48. LDA results, the effect of air injector position on turbulence figures, 1000 LPM tangential flow rate with central air injection, isothermal conditions, method B.

5.3.3.3 Turbulence Intensity Combustion Tests (Method A)

Although swirling flows yield dramatic changes under combustion condition as mentioned previously, turbulence intensity profiles revealed the same trends as in most isothermal conditions. Axial air injection also has a high impact on turbulence intensity values as can be seen from figure 5.49. However, turbulence intensity values under combustion conditions are much lower than isothermal cases.



Figure 5-49. LDA results, the effect of central air injection on turbulence intensity values under combustion conditions (Air flow rate 400 LPM, X = 150 mm, Y = 5 mm).

This effect has been noticed at different flow rates as can be seen from Figures 5.50, 5.51 and 5.52.



Figure 5-50 Effect of central air injection on turbulence intensity values under combustion conditions (Air flow rate 600 LPM, X = 150 mm, Y = 5 mm).



Figure 5-51. LDA results, the effect of central air injection on turbulence intensity values under combustion conditions (Air flow rate 800 LPM, X = 150 mm, Y = 5 mm, method A).



Figure 5-52. LDA results, the effect of central air injection on turbulence intensity values under isothermal conditions (Air flow rate 1000 LPM, X = 150 mm, Y = 5 mm, method A).

Correlation of turbulence intensity values between air and without air injection cases is available just for X = 150 mm, because it is not possible to lit the burner for the other three positions without air injection. However, data acquired for this position reveal the same effect of axial air injection on turbulence intensity as that of the cases under isothermal conditions.

The effect of the air injector position is also clearly pronounced during the combustion tests. Figure 5.53 illustrates the effects of changing air injector positions on turbulence intensity values. The first and second positions, X = 0 and X = 29 mm, undergo high turbulence values. Thus early flame flashback occurs at low tangential flow rates at $\Phi = 0.65$ as mentioned before.



Figure 5-53. LDA results, the effect of air injector position on turbulence figures, 600 LPM tangential flow rate with central air injection, Y = 5 mm, method A, combustion conditions.

However, at X=150 mm position turbulence intensity values are reduced considerably. Hence flashback resistance can be achieved until Φ =0.7 at the same tangential velocity level 2.8 m/s.

Moving to higher flow rates 800 LPM, tangential velocity 3.74 m/s figure 5.54 some changes on turbulence intensity are observed. Nevertheless, X=29 mm still has the highest value, while for X=0 mm, the strength of turbulence decreases, showing the movement of the CRZ at different conditions.



Figure 5-54. LDA results, the effect of air injector position on turbulence figures, 800 LPM tangential flow rate with central air injection, method A, combustion conditions.

A significant enhancement has also proven for X=110 and 150 mm flashback figures as they moved from Φ = 0.65 to Φ =0.77. At 1000 and 1200 LPM, figures 5.55 and 5.56 show the variation of turbulence as the CRZ moves across the flow field. It is a consequence of an improvement of flashback resistance by pushing the CRZ out of the sleeve, reducing turbulence and having a geometry capable of annihilating the CIVB.



Figure 5-55. LDA results, the effect of air injector position on turbulence figures, 1000 LPM tangential flow rate with central air injection, method A, combustion conditions.



Figure 5-56. LDA results, the effect of air injector position on turbulence figures, 1200 LPM tangential flow rate with central air injection, method A, combustion conditions.

5.3.3.4 Turbulence Intensity Combustion Tests (Method B)

Turbulence intensity measured based on bulk flow velocity shown the same trends as that of isothermal conditions regarding the effect of central air injection when changing the air injector position and the tangential mass flow rates.

Figure 5.57, shows the effect of the central air injection on turbulence intensity profiles at a low tangential flow rate (400 LPM). Turbulence intensity has decreased considerably upon injecting air axially at these tangential flow rate values. The maximum amount of reduction is mainly around the burner central axis with approximately 50% than that without air injection. However, near to nozzle boundary, almost no effect has been perceived.

This effect still obvious at even higher flow rates as can be seen from figures 5.58 and 5.59. Although the central air injection increases the turbulence intensity at the boundary of the shear layers almost 12 %, it still affects in producing considerable reduction at the burner central axis. This decrease is nearly 35%.



Figure 5-57. Effect of central air injection on turbulence intensity values under combustion conditions (Air flow rate 400 LPM, X = 150 mm, Y = 5 mm, method B).

Thus despite small increments of turbulence intensity values, the huge reduction in the centre can be effectively observed as the CRZ has been pushed downstream. However, there is a less effect on the central turbulence of the system. Figures 5.57 to 5.59. There is an acute reduction of the turbulence at the tip of the recirculation bubble, possibly consequence of a leaner combustion and axial incoming flows.



Figure 5-58 Effect of central air injection on turbulence intensity values under combustion conditions (Air flow rate 800 LPM, X = 150 mm, Y = 5 mm, method B).



Figure 5-59 Effect of central air injection on turbulence intensity values under combustion conditions (Air flow rate 1000 LPM, X = 150 mm, Y = 5 mm, method B).

For both cases, isothermal and combustion it appears that central air injection pushes the tip of the central recirculation zone (CRZ) downstream. Introducing non-swirling flow in the burner central axis leads to increase the radius of the vortex core, whereas wider vortex core results in lower radial pressure gradient, consequently reduce the baroclinic torque and thus prevent the conditions of formation recirculation bubble at this region which is one of the main reason that provokes CIVB, consequently promote better flame flashback resistance, this effect has also proved before in previous works [35, 87, 160].

It has also previously elucidated that keeping the vortex core radius as constant as possible in the axial direction or at least decreasing with streamwise direction is recommended to achieve good stability conditions [145]. Thus, to achieve constant vortex core radius, the central air injection should still effective at a certain distance downstream burner centre, as can be seen from Figures 5.10 and 5.11 that central air injection still effective at Y=15 and 25 mm respectively, hence it can ensure large vortex core radius.

Nevertheless, the degree of axial air injection effect on axial velocity defects is less than that of Y = 5 mm, which explains why it is important to have constant (axial/tangential) flow rate ratio. In other words, the amount of central air injection must be proportional with increasing the amount of tangential flowrate as mentioned

before. For instance, at low flowrates, the amount of axial air injection is around (10%) to achieve the required effect regarding flame flashback resistance. However, at high flowrates, the ratio is decreased to almost (4 %). Thus keeping axial to tangential flow ratio at about 8-10 % is essential to achieve the desired stability operation.

5.3.4 Flame re-Stabilisation

The high potential flame flashback resistance of axial air injection can provide an excellent flame re-stabilisation technique. Thus this effect can be considered as a great outcome regarding avoiding the consequences arises when combustion system subjected to different flame flashback mechanisms, such systems especially that fuelled by high turbulent flame speed fuel like pure hydrogen or high hydrogen blends can operate more safely if central air injection is used as flame stabiliser technique.

Figure 5.60 illustrates this effect, upon increasing tangential gas flow rates at a constant central air injection, the flame starts propagating upstream with the collapse of coherent structures (CRZs, PVC, shear layers and HMFR) as can be seen from figures 5-60 -a, b, c. However, high momentum of the axial air injection at the central axis allows the flame to stay for a period inside the nozzle, figure 5-60 -d, and e. noticeably, at this stage the central air flow prevents flame flashback into the burner. Consequently, central core flame propagation is prevented probably by the distortion of the combustion induced vortex breakdown flashback (CIVB). Upon a small increment of axial air flowrate, the flame re-stabilises again and swirl coherent structures re-formed, figure 5.60 f-i.



Figure 5-60 effects of axial air injection on flame re-stabilisation.

5.3.5 Summary

The demands for having more reliable combustion systems regarding operation at a wide range of stability margins and avoids undesirable flame flashback mechanisms still represent urgent priority importance for gas turbines sector.

Although traditional techniques that were used previously and even now to protect combustors from combustion instabilities can ensure safe operation under different conditions, they also have some undesirable consequences, using central fuel injectors or bluff bodies as flame stabilisers increase the maintenance coast and the need of replacing such injectors after a period due to degradation. Moreover using central fuel injection can increase the pollutant level significantly. Thus this chapter proposed the validity of using axial air injection as a flame stabilisation mechanism and the effect of this technique on flame characteristics. The key results can be summarised as follows:

- 1. Wider stability operation region can be achieved when using axial air injection compared with the case when central fuel or bluff body is used to promote flame stability. This finding is very important in terms of switching to another fuel, especially hydrogen or high hydrogen blends that have the higher flame speed that makes the blow off and flashback limits very close to each other which in turns impose some changes in system combustor hardware to make it able for switching [59].
- 2. Central air injection can significantly reduce the defect in the axial velocity near the burner centerline, this in turn lead to produce wider vortex core. Whereas wider vortex core results in the lower pressure gradient, consequently reduce the baroclinic torque which is one of the parameters that initiate upstream flame propagation, hence promote better flame flashback resistance.
- 3. Central air injection can significantly affect the flow field characteristics downstream of the burner dump plane, this effect is mainly represented by the considerable reduction of turbulence intensity values, and since turbulent flame speed is proportional to the turbulence fluctuations, its values will, in turn, reduce to a considerable level, hence decrease in the flame flashback propensity. This variation in turbulence intensity is also important in switching to different fuels, whereas the switching process requires a change of mass flow rate at constant output power due to the difference in fuel heating values. Thus increasing or decreasing mixture flow rate can lead to some variation in turbulence intensity, this variation can be avoided by using central air injection through affecting turbulence intensity level or even compensate the change in mass flowrate.
- 4. The effect of central air injection on flame characteristics and hence flame flashback depends on the (axial/tangential ratio), at low flow rate axial air injection effects are significantly evident. However upon increasing swirl tangential flowrate this effect start to decrease, thus keeping constant ratio is
important. The amount of reduction in swirl number should be taken into consideration, a higher amount of axial injection can degrade the swirl strength considerably.



Effect of Change in Downstream Velocity Gradient on Flame Flashback Mechanisms

Chapter 6

The true sign of intelligence is not knowledge but imagination Albert Einstein

6.1 Introduction

In most practical combustion systems the achievement of high flame flashback resistance depends on flame stabilisation downstream burner nozzle, which depends on the equilibrium between flame speed and incoming flow velocity at the reaction zone both in magnitude and direction. This equilibrium, in turn, is a function of different parameters such as burner configuration or geometry, the degree of mixing (premixed, partial premixed or diffusive), fuel type, initial conditions of the mixture (pressure and temperature) and working conditions inside the combustion chamber.

This flow field balancing is quite simple to attain in diffusion flows. However, in swirl flows the situation becomes so complicated because of the complex nature of these premixed flows. This complexity stems from the generation of different coherent structures and mixing parameters. The structures reveal considerable variation between each other in terms of having significant differences in magnitude and direction to the flow field velocity. For instance, the central recirculation zone CRZ may have high-velocity values in flow field in the opposite direction of the flow (negative direction) which is one of the important features of swirl flows. Moreover, other structures such as shear layers can undergo flow stagnation (zero velocity region) due to the interaction between the high momentum flow region and the CRZ.

According to Lewis and von Elbe [147] when velocity gradients in the boundary layer sublayer became very low or lower than the flame speed in this region, a flame flashback occurs via the sub-layer. Although, this hypothesis is for wall boundary layer, moreover, the velocity gradient approach seems to be a useful tool to describe flame flashback via central core or combustion induced vortex breakdown (CIVB) and how can these types of flashback be affected by the variation in velocity gradient.

6.2 Results and Discussion

6.2.1 Effect of Air injection Position on the CRZ Propagation

The effect of air injector position on axial velocity and turbulence generated at the burner exit plane has been discussed in chapter 5. However, it seems that the air injection position does not affect the flow field characteristics close to the burner mouth, but also it can considerably change the shape and size of the swirl coherent structures, especially the CRZ. Figure 6-1 shows this effect under isothermal conditions.



Figure 6-1. LDA results, the effect of air injection on the position and characteristics of the CRZ, isothermal conditions, X= 150 mm.

From figure 6-1 it is apparent that the axial air jet is pushing the CRZ, as the CRZ without air injection is very close to burner mouth, and its tip is almost parallel to the inside nozzle wall. Under combustion conditions, the existence of the CRZ in this position can easily affect the flow field characteristics, the balancing between heat release generated at the tip of the CRZ and the cooling due to the coming flow can

promote the baroclinic torque which can finally lead to CIVB. Contrary, the air jet is pushing the tip of the CRZ downstream to Y/D = 0.2. This distance is enough to achieve the required balancing between volume expansion and baroclinic torque at the tip of the CRZ hence reducing the possibility of initiation of the CIVB.

As previously mentioned in Chapter 5, high turbulence near burner mouth can lead to the increase of turbulent flame speed and as a consequence initiate flame flashback. The root mean square of velocity fluctuation, which reflects the turbulence level, seems to be affected by the axial air injection not just close to the burner exit plane, but also in all downstream flow domain, figure 6-2.



Figure 6-2. LDA results, the effect of air injection on the amount of turbulence generated downstream the burner exit plane, X=150 mm.

It obviously appears that axial air injection reduces the turbulence level considerably in all the flow field domain. Without air injection, the maximum turbulence level is located close to burner mouth and on the boundary of the bottom of the CRZ, at the shear layer. Root mean square velocity fluctuation is about 1.4-1.7 m/s. However, RMS values at the same region using axial air injection is around 1.0 m/s.

This change in turbulence profile could also attribute to the effect of air injection on the type of the generated vortex breakdown; axial air injection can lead to suppress the flow field, consequently change the type of the vortex breakdown from bubble to cone type [198]. It has proved that bubble vortex breakdown exhibits a local defect in axial velocity values in its upstream plane, while cone type vortex breakdown does not, hence the latter is preferred for flashback safety requirements [70]

The position of the air injector also has a high impact on the CRZ position and the turbulence profile downstream. However, this depends on the ratio between the tangential to axial flux momentums. Figure 6-3 illustrates the difference in CRZ characteristics at two different air injection positions, X=150 and 29 mm.



Figure 6-3. LDA results, the effect of air injection position on the position of the CRZ, isothermal conditions.

From figure 6-3 it can be seen that when air injection opens at X=29 mm downstream the burner baseplate, the size of the CRZ becomes smaller. This reduction of the CRZ size results in a slightly wider shear layer or high momentum flow region than that of the X=150 mm.

The existence of a wider shear layer or high momentum flow region close to the burner mouth can lead to an increased turbulence level at this region, consequently increasing the local turbulent flame speed at a small area, and hence provoke CIVB. The effect of air injector position became more obvious when considering the change in turbulence level downstream and close to the burner mouth. Figure 6-4 shows the RMS figures for two air injector positions, X=150 and X=29 mm.



Figure 6-4. LDA results, the effect of air injection position on the turbulence level downstream, isothermal conditions.

From figure 6-4 it can be observed that when air injector is positioned at X=150 mm, turbulence intensity levels are noticeably lower in all the flow field domain than those

at X=29 mm. Moreover, when considering the region between r/Ro= 0.3 to 0.9 at X=29 mm position, which is very close to the burner mouth, it appears that the RMS values are about 1.7 m/s compared to RMS= 1.0 m/s in the same region for the X=150 mm position. Although the difference in RMS values between the two cases is not too high, this can produce a considerable change in turbulent flame speed values. Additionally, it is important to mention here that this measuring position is at Y=3 mm downstream the burner mouth, as it was practically difficult to get closer to the burner mouth due to the reflection of one of the laser beams. Thus it is hypothesised that this region of high turbulence extended towards the burner exit plane or even more inside the nozzle at X=29 mm.

6.2.2 Effect of Air Injection on the Axial Velocity Gradient

According to the isothermal findings, it appears that air injection can affect the CRZ position and all the flow field characteristics downstream the burner exit plane. Thus to get more insight the investigations were also carried out under combustion conditions to determine the effect of change of velocity gradients inside the CRZ close to flame flashback. The comparisons and correlation are based on the measurements of the flow field velocity downstream the burner mouth by using LDA with and without central air injection. These measurements have been done at almost the same equivalence ratios, $\Phi = 0.59$ for the case without air injection and $\Phi = 0.60$ for the case with air injection, while the air inlet tangential flow rate input at 800 LPM. The central air injection used was 50 LPM which represent about 5% of the total flow rate. This ratio ensured the required effect of flame stabilisation, and in the same time, it avoids reduced swirl strength as mentioned in chapter 5.

The matrix of the measurement domain downstream the burner mouth has been divided into 21 points in the axial direction (r/Ro) and 8 points in a vertical position (Y/D). This domain ensures the total inclusion of the flame region, as can be seen from figure 6-5. The test were repeated five times until the deviation was not greater than 5% in the axial velocity values, to ensure high accuracy results and avoid uncertainties.

Figure 6-6 shows the effect of axial air injection on the CRZ position downstream the burner exit plane. As discussed before in the isothermal investigations, under combustion conditions the axial air injection also has a considerable effect on the position and size of the CRZ. It can be observed from figure 6-6 when no axial air injection is used that the CRZ propagates upstream and almost enters to the nozzle. This position leads to critical stability conditions. The heat release distribution with respect to the flame position in this region can easily initiate the formation of the recirculation bubble at the tip of the CRZ, as a consequence onset of the CIVB [163].



Figure 6-5. Measurements matrix.

While under the effect of axial air injection the CRZ is shifted downstream, thus reducing the possibility of quick formation of the recirculation bubble under the effect of the interaction between flame and heat release from the burner wall.

Although axial air jets can push the CRZ to a quite safe stability position regarding good flame flashback resistance, they reduces the size of the CRZ, hence producing weak swirl strength. Thus, to avoid this obstacle a proper balance between the amount of axial injection and swirl strength is highly important to achieve both good stability limits and swirl numbers. This effect will be discussed briefly in section 6-3.

Figure 6-6 also illustrates that axial air injection can significantly reduce the amount of the defect in the axial velocity values downstream the burner mouth. Thus it is evident that the negative velocity region with air injection is less pronounced with axial air injection. This lead to conclude that axial air injection not only pushes the central recirculation zone but also reduces negative axial velocity values in the vortex core that can initiate upstream flame flashback or CIVB.



Figure 6-6. LDA results, the effect of air injection on the CRZ position downstream burner mouth, combustion conditions.

Since the velocity values change when the air is injected axially, it is predicted that velocity gradients change as well. Although higher variation in velocity gradient values is one of the important features of swirling flows due to the existence of different coherent structures that have different velocity values and directions, this significant difference may promote sudden upstream flame propagation, especially CIVB. It was reported before that moderate increasing of axial velocity in the streamwise direction characterises the optimum velocity distribution for flashback resistance [145]. In other words, moderate axial velocity gradient is beneficial to improve CIVB flashback resistance. Moreover, the axial air jet at the burner axis can affect the mean axial pressure gradient, this in turn enables suitable conditions for

axial velocity to match the turbulent flame speed and hence allows stable flame. Figure 6-7 shows the difference in velocity gradient profiles when axial air injection is used.



Figure 6-7. LDA results, the effect of air injection on the velocity gradient downstream burner mouth (combustion conditions).

From figure 6-7 it can notice that the change in velocity gradient when air injection is used is smoother than the case when no axial air injection is introduced to the flow. Thus this variation in axial velocity gradient under the effect of axial air injection is crucial to avoid upstream flame propagation, especially at the burner axis. The positive or moderate velocity gradient in this region can play an important role in pushing downstream the recirculation bubble generated at the tip of the central recirculation zone. Herein this variation in velocity gradient provides considerable resistance against the opposing induced force in the main flow, leading to a positive change in azimuthal vorticity, consequently tackling the CIVB. The velocity gradient values are depicted to have a clear comparison for each (r/Ro) position at different (Y/D) positions. The main focus here is on velocity gradients inside the CRZ.

Figure 6-8 shows the difference in velocity gradients for two cases, with and without air injection for r/Ro= 0.0. It is obvious that the use of air injection leads to almost constant or step velocity gradient. On the contrary, without using central air injection, the variation in velocity gradient values is noticeably sharp. Thus this dramatic change in flow field makes the system less reluctant against upstream flame propagation and this region, in particular, less resistance to CIVB.



Figure 6-8. Comparison of change in velocity gradient values in the axial direction (with and without air injection) at r/Ro = 0.

However, at r/Ro= 0.13 and 0.26 the effect of air injection is reduced as can be seen from figures 6-9 and 6-10.



Figure 6-9. Comparison of change in velocity gradient values in the axial direction (with and without air injection) at r/Ro = 0.13.



Figure 6-10. Comparison of change in velocity gradient values in the axial direction (with and without air injection) at r/Ro = 0.26.

At r/Ro = 0.39 and r/Ro = 0.78, figures 6-11 and 6-12, the effect almost disappears and becomes similar to that without air injection. This is attributed to flow stagnation at the boundary of the CRZ due to the existence of the shear layer.



Figure 6-11. Comparison of change in velocity gradient values in the axial direction (with and without air injection) at r/Ro = 0.39.



Figure 6-12. Comparison of change in velocity gradient values in the axial direction (with and without air injection) at r/Ro = 0.78.

At r/Ro= 0.90 and 1.04, which lie outside the CRZ region, the variation has changed. Air injection cases reveal slightly higher velocity gradient values, figures 6-13, and 6-14. This effect could be the consequence of the increase in positive axial velocity at the burner centre that enforces flame to propagate via the nozzle wall, i. e. boundary layer flashback (BLF) as mentioned in chapter 5 section 5.3.4, and illustrated in figure 5.60.



Figure 6-13. Comparison of change in velocity gradient values in the axial direction (with and without air injection) at r/Ro = 0.9.



Figure 6-14. Comparison of change in velocity gradient values in the axial direction (with and without air injection) at r/Ro = 1.04.

6.2.3 Effect of Velocity Gradient on Axial Velocity Values Close to Flashback

The equilibrium between local flame speed and local flow velocity is the crucial factor in maintaining flame stabilisation. Velocity gradients is the key features that promote this equilibrium, Cheng, et al. [199] proposed a good correlation between velocity gradient in the axial direction and turbulent flame speed; this correlation is suitable for low swirl numbers

$$U_b - \frac{dU}{dY}(x_f - x_0) = S_T \tag{6-1}$$

Where:

- U_b : Bulk flow velocity [m/s]
- $\frac{dU}{dV}$: Velocity gradient in streamwise direction [1/s]
- X_f : Leading edge position of the flame brush [m]
- X_0 : the reference point of the divergent flow or the stagnation point at the tip of recirculation bubble [m]
- S_T : Turbulent flame speed

The validity of this equation depends on the accuracy of measurements of velocity gradient (dU/dY), the measuring positions need to be as close as possible to each other. However, the size of the probe control volume must be taken into consideration. According to the above equation by controlling velocity gradients, it is possible to control turbulent flame speed, and hence central air injection that can affect velocity gradients. As a result, it can also affect local turbulent flame speed values. Considering figure 6-8, at positions Y/D = 0.24 to 0.60, the case without air injection reveals negative velocity gradients that according to equation 6-1 will lead to an increase in turbulent flame speed. Furthermore, the sharp increase in velocity gradient values downstream cause a significant increase in turbulent velocities. The higher turbulent velocity values and sudden change along the downstream axis can promote sudden upstream propagation of the recirculation bubble, consequently, increasing the possibility of provoking CIVB flashback.

Contrary the case with air injection shows velocity gradient values at the same positions being positive, consequently, decreasing turbulent flame speed. The variation in these values along the downstream axis is almost constant. Hence no considerable change in turbulent velocity values is observed. This fact explains why central air injection can provide high flame flashback resistance and enable wider equivalence ratios. This significant effect of central air injection on velocity gradient has a direct effect on the upstream movement of the CRZ and hence the onset of the CIVB.

The effects mentioned above of change in axial velocity gradient values can be explained when considering the amount of change in axial velocity values when increasing equivalence ratios from stable to flashback conditions. Figure 6-15 shows the change in axial velocity values when increasing equivalence ratio from stable operation (Φ_{stable}) to that close to flashback conditions (Φ_{FB}) without air injection, while figure 6-16 illustrates the case when air injection is used.

From figure 6-15 it can be seen that the axial velocity has changed by approximately 0.5 m/s when the equivalence ratio is increased towards flashback conditions. This change in velocity refers to the change of the CRZ position with respect to the measurement level. In other words, the high-velocity gradient at the central axis leads to a sudden jump of the CRZ towards the incoming flow.



Figure 6-15. Change in axial velocity values during the transition from stable to flashback conditions, X= 150 mm, Y = 5mm, no central air injection, air inlet tangential flow rate 800 LPM.

However, when the axial air is injected at the central axis the amount of change of axial velocity values is very low, although the change in equivalence ratio from (Φ_{stable}) to (Φ_{FB}) in both cases is almost the same, figure 6-16. This finding is attributed to the moderate velocity gradient under the effect of axial air injection.



Figure 6-16. Change in axial velocity values during the transition from stable to flashback conditions, X= 150 mm, Y = 5mm, with central air injection, air inlet tangential flow rate 800 LPM.

The profiles of kinetic energy in an axial-tangential plane also demonstrate this effect. The two dimensional kinetic energy has been determined according to the following equation;

$$Ke = 0.5 \cdot (u^{2} + w^{2}) \tag{6-2}$$

Where:

Ke: the turbulent kinetic energy (m^2/s^2)

u': fluctuating axial velocity, RMS in axial direction [m/s]

w': fluctuating tangential velocity, RMS in tangential direction [m/s]

Figure 6-17 shows that when no central air injection exists, increasing equivalence ratio to Φ_{FB} results in an increase in turbulent kinetic energy values, from 0.4 m²/s² at stable operation to 1.7 m²/s². This change in turbulent kinetic energy at the nozzle dump plane indicates the propagation of the CRZ, whereas higher turbulent kinetic energy in swirls flows associated with boundaries or interaction regions of swirl structures.



Figure 6-17. Change of turbulent kinetic energy during the transition from stable to flashback conditions, X= 150 mm, Y = 5mm, no central air injection, air inlet tangential flow rate 800 LPM.

Contrary, for the case using central air injection, although the values of turbulent kinetic energy are higher than those with no air injection, there is no change in turbulent kinetic energy values when equivalence ratio is increased towards (Φ_{FB}), figure 6-18, which reveals that the CRZ is still downstream the burner dump plane.

The effect of axial velocity gradient is also observed at higher inlet air tangential flow rates. Figure 6-19 shows a comparison between the axial velocity values at stable and close to flashback conditions under high inlet tangential flow rate of 1000 LPM with central air injection. It is quite clear that variation in the axial velocity values at the tip of the CRZ are not very high. Hence this proves that central air injection still prevents sudden jumps of the CRZ. Turbulent kinetic energy profiles at this flow rate also demonstrate this fact as can be seen from figure 6-20.



Figure 6-18. Change of turbulent kinetic energy during the transition from stable to flashback conditions, X= 150 mm, Y = 5mm, with central air injection, air inlet tangential flow rate 800 LPM.



Figure 6-19. Change in axial velocity values during the transition from stable to flashback conditions, X= 150 mm, Y = 5mm, no central air injection, air inlet tangential flow rate 1000 LPM.



Figure 6-20. Change of turbulent kinetic energy during the transition from stable to flashback conditions, X= 150 mm, Y = 5mm, with central air injection, air inlet tangential flow rate 1000 LPM.

Although there is the difference in some values, the effect of axial air injection on the axial velocity gradient is evident. The change in axial velocity during the transition to flashback conditions has been observed at the other air injector positions, i.e. X=0, 29 and 110 mm which are almost the same of those of X = 150 mm.

However, results without air injection at these positions are not available because it was very difficult to start the burner without air injection. Figures 6-21, 6-22 and 6-23 show the change in axial velocity values during the transition from stable to flashback conditions.

It is important mentioning here that as the position of the central air injection gets closer to the nozzle exit, the velocity gradient becomes more moderate, hence the change in axial velocity values is less pronounced, consequently reducing the distance that the CRZ moves in the upstream direction.

The change in turbulent kinetic energy when moving from the X = 0 mm to the X=29 mm, X=110mm and X=150 mm positions elucidate this variation in the location of the CRZ. At X=0 mm the figures of turbulent kinetic energy reveal that the measurement is at the tip of the CRZ, while, at close to flashback conditions the lower kinetic energy values indicate that measurement levels were positioned inside the CRZ. Further pilot positions, i.e. X=29 mm and X=110 mm showed that the axial

velocity gets closer indicating further pushing of the CRZ, and finally at X=150 mm is the optimum position for produce the required axial velocity gradient, consequently best CIVB resistance.





Figure 6-21. (a) Change in axial velocity values (b) Change of turbulent kinetic energy, during the transition from stable to flashback conditions, X= 0 mm, Y = 5mm, central air injection, air inlet tangential flow rate 800 LPM.





Figure 6-22. (a) Change in axial velocity values (b) Change of turbulent kinetic energy, during the transition from stable to flashback conditions, X= 29 mm, Y = 5mm, central air injection, air inlet tangential flow rate 800 LPM.



Figure 6-23. (a) Change in axial velocity values (b) Change of turbulent kinetic energy, during the transition from stable to flashback conditions, X= 110 mm, Y = 5mm, central air injection, air inlet tangential flow rate 800 LPM.

6.3 Effect of Axial Air Injection on Swirl Strength

Swirl strength linked to swirl number is a crucial factor in the formation and characteristics of swirling coherent structures. This criterion depends on the ratio of tangential to the axial flux of swirling flows. Introducing axial jets into the vortex core can significantly degrade the swirl strength, high amount of axial injection could also lead to a decrease in swirl numbers to a lower degree below the required for vortex breakdown generation which is necessary for the formation of the swirl coherent structures. Moreover, such axial jets can suppress the PVC in such a way that affects its interaction with other swirl structures. Thus, to achieve the desired benefits of using

central air injection with a high swirl to promote flame stability requirements, it is highly important to make an appropriate balance between axial and tangential flow rates of swirling flows.

Although geometric swirl numbers are used to characterise the swirl strength, they can change locally according to the ratio of (tangential/axial flow rates). Thus this numbers can vary when a considerable amount of axial injection is used. Therefore the local swirl number for isothermal conditions where density is assumed to be constant was determined by equation (2-8) based on burner geometry, inlet conditions and neglecting pressure variations [102]. Table 6-1 shows the local swirl number at different flow rates with axial air injection. Figure 6-24 shows the effect of increasing the inlet tangential flow rates at constant axial air injection on the size and characteristics of the CRZ.

Tangential flow rate (l/min)	Axial flow rate (l/min)	Qta (m ³ /s)	Qto (m ³ /s)	Local swirl number
400	50	0.006	0.007	0.725
500	50	0.008	0.009	0.758
600	50	0.01	0.010	0.782
700	50	0.011	0.012	0.799
800	50	0.013	0.014	0.813
900	50	0.015	0.015	0.823
1000	50	0.016	0.017	0.832
1100	50	0.018	0.019	0.839
1200	50	0.02	0.020	0.846
1300	50	0.021	0.022	0.851
1400	50	0.023	0.024	0.855
1500	50	0.025	0.025	0.859
1600	50	0.026	0.027	0.863
1700	50	0.028	0.029	0.866

Table 6-1. Local swirl numbers.

As can be seen from the figure 6-24, at low inlet tangential flowrates 800 LPM, the axial air injection considerably affects the axial velocity values in the CRZ. However,

very low-velocity gradients could have an impact on the formation of the vortex breakdown. Upon increasing the inlet tangential flowrate to 1200 and 1600 LPM, the CRZ becomes more robust and coherent and significant variation in axial velocity values is achieved.



Figure 6-24. Effect of axial air injection on the CRZ size and axial velocity defects. Isothermal conditions.

6.4 Summary

Central air injection has proved to be much advantageous regarding enhancement of flow field structures in such way that promotes highly flame flashback resistance. The existence of central air injection not only affects the turbulence intensity values downstream the burner nozzle, but also it extends further downstream, and significantly affects the axial velocity gradient especially across the CRZ. This downstream velocity gradient even for small values is extremely important in terms of controlling the CRZ position and hence to avoid the formation of the CIVB. The main results and outcomes of the effect of changes in the downstream axial velocity gradient can be summarised as follow:

- 1. Injecting small fractions of air axially can significantly affect the downstream flow field axial velocity gradient. This was mostly observed inside the CRZ.
- The maximum change in axial velocity gradient was observed at the central axis of the burner. The effect of air injection in this context is producing considerably moderate axial velocity gradients compared to those for the no air injection case.
- 3. The variation in axial velocity gradient when air injection is used is reduced when moving outwards radially. It almost disappears at r/Ro= 0.26, 0.39 respectively and became almost the same of that without air injection. These positions represent the region of CRZ boundary and where the shear layer exists.
- 4. Outside the CRZ the situation is changed inversely. The air injection case reveals a velocity gradient higher than that with air injection. This effect was assumed as the high momentum axial air injection enforce the flame to propagate mainly in the annular region close to nozzle boundary, as shown previously in figure 5-60.
- 5. Since the axial air injection decreases or moderates the axial velocity gradient it affects the axial velocity values and hence it promotes the propagation of the CRZ when increasing equivalence ratios to rich conditions. The axial velocity measurements at Y/D = 0.08 proved this fact.
- 6. Turbulent kinetic energy in an axial-radial plane also demonstrated the effect of axial air injection on the CRZ position.
- 7. The position of the axial air injector is very important regarding the amount of change in the axial velocity gradient, hence the CRZ propagation, consequently the formation of the CIVB. X=150 mm position is the optimum position that can change axial velocity gradients considerably.
- 8. Velocity gradient model is not only suitable for investigating boundary layer flashback, but also it seems to be vital in detecting CRZ propagation, consequently determining the CIVB conditions.

9. The ratio of axial to tangential inlet flowrates is very important in determining local swirl numbers, consequently the shape and size of the CRZ, which plays an important role in terms of stability limits.



Enhancement of Flashback Resistance against Boundary Layer Flashback

Chapter 7

An experiment is a question which science poses to Nature, and a Measurement is the recording of Nature's answer

Max Planck

7.1 Introduction

Although using the appropriate diameter of central fuel injector or central air injection can considerably tackle upstream flame propagation through the central core, especially CIVB, some drawbacks can arise, the system could be more likely subjected to wall boundary layer flashback (BLF), especially at higher tangential flowrates.

Large diameter fuel injectors enforce the outer shear layer to propagate radially outwards and to stabilise on the nozzle lip, Figure 7-1a. In most practical applications this contact between flame and nozzle lip have some consequences represented by life degradation of nozzle material due to continuous high temperatures in addition to the possibility of increasing pollutants level. At high equivalence ratios, the outer boundary layer between the shear layer and nozzle wall starts an upstream propagation via low-velocity sublayer leading to a flame flashback.

High momentum axial air injection also performs the same function of the central fuel injector, i.e. the high resistance against central upstream flow breakdown propagation imposes radial propagation of the flow. Hence the outer shear layer locates on the burner rim or inside the nozzle, Figure 7-1 b. Despite using central air injection can maintain a stable flame slightly longer than that of using central fuel injectors, boundary layer flashback can onset when increasing equivalence ratios.

Thus this chapter describes the appropriate techniques to reduce the possibility of boundary layer flashback when central fuel injectors or central air injection are used as flame stabilisation techniques.





(b) Central air injection

Figure 7-1. Effect of using central fuel injector and central air injection on the outer boundary layer propagation.

7.2 Nozzle Surface liner for Enhancement of Flame Flashback Resistance

Flame flashback via wall boundary layer depends on many parameters such as the flow field characteristics, equivalence ratio, pressure, temperature, wall temperature, confinement type, the state of the boundary layer and the geometry of interior liners in the burner nozzle [69, 148]. The geometry of the nozzle wall plays an important role in upstream flame propagation during boundary layer flashback, i.e., the interaction between nozzle wall and flame can affect directly the amount of heat flux which consequently changes the wall quenching distance [200].

Furthermore, the interaction between nozzle wall surface and the parallel flow generates a viscous drag which produces an adverse pressure gradient, consequently promoting velocity gradient. The degree of wall roughness is of particular importance in this context as it promotes the amount of heat transfer, hence decreasing or increasing the shear wall stress. Consequently, the velocity distribution mode and the difference in Reynold stresses, especially in the sublayer. The contribution of the effect of the surface type on the mean velocity profile and hence the wall turbulent boundary layer usually is described by a roughness function which represents the difference in normalised velocity distribution between smooth and rough surfaces[201].

The direction of flame movement with respect to the wall is also important when considering its effect on flame flashback mechanism and its quenching position. From one hand, when the flame front is parallel to the wall, the quenching mechanism is called Head-on quenching (HOQ); on the contrary when the flame is perpendicular to the wall the mechanism is called side-wall quenching (SWQ)[202].

However, in swirling flows the flow directions are very complex compared to other flow types, as the existence of coherent structures and their interaction with each other make the analyses and prediction of wall boundary layer characteristics a seriously challenging task. These flows are extremely affected by the variation of flow streamlines and the adverse pressure gradient downstream all directions (axial, azimuthal and radial). Therefore the flame direction with respect to the wall during boundary layer flashback (BLF) cannot be simply studied for a certain direction, and the predicted movement will depend on the inside linear geometry of the nozzle.

Although most of the swirling flows are turbulent, the flame-wall interaction is considered laminar flow due to the considerable reduction of flow velocity in the adjacent region to the wall that results from viscous friction. Thus, most hypotheses about boundary layer flashback are based on laminar flame speed. However, this correlation can be adopted for turbulent flows as well. Microsurfaces of different geometries can positively increase the boundary layer flashback resistance. Those surfaces have high potentials in reducing drag effect in the wall adjacent region. However, their ability in dealing with drag forces depends mainly on their configuration and the drag force.

Two different configurations have been used to determine their effect on boundary layer flashback (Two woven steel grids or micromesh of size 50 μ m, 150 μ m). The grid makes a liner; the liner thickness was scanned by Shared Labs Europe LTD [203]. Figure 7-2 shows the geometry and dimensions of the liner, and figure 7-3 shows the position of the micromesh inside the burner nozzle.





Figure 7-2. The geometry of the Microsurfaces.



(a)



(b)

Figure 7-3. Position of the microstructures inside the burner nozzle, (a) 50 µm micromesh (b) 150 µm micromesh.

7.3 Results and Discussion

The experimental procedure is similar to that of chapter 5. The difference here is based on investigating the effect of an inside surface into the burner nozzle for studies on boundary layer flame flashback. Two different geometrical grid configurations were used, figure 7-3, those grids can change the flow field characteristics and hence boundary layer adjacent to nozzle wall. This change can provide the required flashback resistance. The effect of the geometrical shape of those grids has been analysed using numerical simulation by other authors, the details of the numerical approach and results are illustrated in [204]. Stability limits for different grids were determined under different configurations with the central fuel injector, and with and without central air injection at various inlet tangential flowrates. Central air injection was kept at 50 LPM which accounts for 10% of the total mass flow rate as mentioned before in chapter 5 to avoid swirl strength degradation. The burner inlet tangential velocity has been used to illustrates the flame flashback trends, as this allows a fair comparison between cases without considering the change of Re at the burner mouth and enables comparison with different size burners. Blowoff and flashback limits were determined by the same procedure used in chapter 4 and chapter 5. Experiments were repeated five times until a deviation not greater than 5% in tangential velocity was achieved. The stable operation region is to the right of the blowoff points, and to the left of the flashback points.

During LDA measurements the main concern with using different grid sizes was the possibility of seeding can block the grids during the test, thus after each test a new piece of grid material is used. The velocity profile and turbulence intensity measurements were carried out at 5 mm downstream the burner mouth.

To achieve a complete understanding of the effect of wall microstructures surfaces on flame stability it is vital to ensure that flame flashback mechanism is kept as mainly BLF. Thus two techniques have been used to enforce flame propagation via wall boundary layer, i.e. central fuel injector and axial air injection.

7.3.1 Central Fuel Injector Method

For a central fuel injector, the outside diameter was 23 mm, this ratio of this diameter to the nozzle inside diameter is $\chi = 0.378$, this ratio is enough to ensure BLF as proved before [205] and explained in figure 5-3. While air injection is proved to be a more efficient in preventing upstream flame propagation through the central core, hence high resistance against CIVB consequently the flashback mechanism is mainly BLF as explained in details in chapter 5. Figure 7-4 shows the difference in flame flashback trends when using central fuel injector as a flame stabiliser or bluff body with no liner and that when the 50 µm grid used as linear for nozzle wall.



Figure 7-4. The effect of using the 50 µm Grid microsurfaces on the stability of operation.

As can be seen from figure 7-4, when no grid is used, the stability operation region is located at $\Phi \sim 0.45$ -0.65, and the maximum tangential velocity that can be achieved is almost 4 m/s. This could be attributed to the limitation of nozzle surface resistance to boundary layer flashback at high tangential velocities. Higher values of tangential velocity lead to outward radial propagation of the CRZ which in turn pushes the high momentum flow region closer to the wall. Due to the lack of damping mechanism against BLF, flashback can occur earlier.

Contrary when the micromesh is imposed (50 μ m grid) the change in drag force, and hence downstream velocity gradient will positively support the BLF resistance, for this reason, a flame flashback occurs at slightly higher equivalence ratios. Moreover, at higher flow rates the potential flame flashback resistance provided by the micromesh consolidates the flame stability, and no flashback is observed.

With the existence of the micro-grid surface, the sudden variation from high axial velocity values at the central axis of the nozzle to the low-velocity region near the wall is reduced, consequently lower gradients in the velocity values at the boundary layer region near to the nozzle wall are achieved. Thus the velocity gradient on the immediate vicinity of the wall became lower enough than the critical value.

Moreover, this change in velocity gradient will directly reduce the pressure gradient generated at the tip of the flame. As previously documented the adverse pressure gradient created by the flame backpressure causes deflection and retardation of the axial approaching flow, which in turn accelerates upstream flame propagation [53].

It is difficult to directly measure the flow field characteristics in the adjacent region close to burner nozzle wall because of the difficulty of laser access. Thus the turbulence profile and hence the velocity gradient measured at 5 mm downstream this region at the burner exit plane, is a useful tool that can also describe the effect of the micromesh properties on flame flashback resistance. Figure 7-5 shows turbulent intensity figures for two configurations, without using microsurfaces and with the grid. The change in drag force in the adjacent to nozzle leads to a decrease in the turbulence fluctuation compared with that when no grid being used. Consequently, this reduction in turbulence decreases the local turbulent flame speed, hence promotes good flashback resistance.

The effect mentioned above of the micromesh on turbulence intensity values revealed the considerable variation in flow field velocity and hence the axial velocity gradient.



Figure 7-5. The effect of using 50 µm Grid microsurfaces on turbulence intensity.
The effect of using another micromesh configuration (150 μ m) on flame flashback trends is almost the same for that of 50 μ m regarding equivalence ratios range. However, this configuration led to flame flashback late up to 5 m/s of tangential velocity range, figure 7-6.

This difference in flashback trends suggested that wider micromesh configuration has less effect on velocity gradient and other flow field characteristics close to the nozzle wall.



Figure 7-6. The effect of using 150 µm Grid microsurfaces on flame flashback trends.

7.3.2 Air Injection Method

Axial air injection has proved its potentials for flame flashback resistance, as explained in details in chapter 5. However, this resistance is mainly against CIVB, as it has been visually observed that the high momentum axial jet at the central axis enforces the flame to propagate via the nozzle wall boundary. The phenomenon experimentally demonstrated in chapter 6 by investigating the variation of axial velocity gradient downstream nozzle wall under the effect of axial air injection. Thus another flashback damping mechanism for wall boundary layer is also required. Figure 7-7 shows a comparison between two configurations. From figure 7-7a it can be observed that at a tangential velocity of 1.8 m/s, blow off and flashback curves are very close to each other leading to a narrow stability region. Upon increasing equivalence ratio and tangential velocity, the stability map became wider. However, there is still some limits over which the flame cannot be still stable, and flame flashback is observed. The last conditions where flashback is observed were, tangential velocity = 6 m/s, Φ =0.83, beyond these values no flashback is observed.



Figure 7-7. Effect of using; a) air injection and; b) air injection with micromesh grid for flame stabilisation purposes, respectively.

When the grid is used for additional stability support figure 7-7b, significant improvement in flashback resistance is achieved. The flame flashback is observed for inlet tangential velocity ranging from 2 to 2.8 m/s; beyond this limit, no flame flashback occurs, and operation is stable for higher equivalence ratios and high tangential flowrates.

This finding is interesting for switching to higher power operations at constant equivalence ratio or switching to another fuel blend which might have stability operation regions overlapping with the original fuel stability margin. Moreover, the high BLF resistance at high flowrates is very important in protecting the nozzle inside walls due to harsh environment conditions represented by high temperatures when the flame comes in touch with the nozzle walls.

7.3.3 Flame Flashback Resistance Technique for both CIVB and BLF Simultaneously

Based on the previous results, it appears that the two flame flashback resistance mechanisms, i.e., air injection and microsurfaces are working together to achieve highly flame flashback resistance for BLF and CIVB. Figure 7-8 illustrates flashback resistance scenarios when increasing equivalence ratio from stable operation (Φ_{stable}) to flashback conditions (Φ_{FB}).

Figure 7-8a represents stable operation where the flame is anchored downstream the nozzle. The equivalence ratio is increased by increasing the tangential gas flowrate at constant tangential and axial air flowrates; the flame is still stable for a period at constant equivalence ratio, then it starts to move forward and backwards, mainly via burner central axis which refers to the CIVB conditions, figure 7-8b.

However, the axial air injection pushes the flame downstream to its initial stable position. Further increase of the equivalence ratio initiates more flame propagation as can be seen from figure 7-8c. Nevertheless, the high momentum axial air injection is still coherent enough to prevent upstream propagation via central core.

The process of outer boundary layer propagation is now observed from one side of the flame, figure 7-8d, and 7-8e. Finally, all the annular shear layer became totally in contact with nozzle surface or the grid, figure 7-8f.



(a)









(d)



(e)



(f)

Figure 7-8. Flashback resistance scenario for CIVB and BLFB simultaneously.

7.4 Summary

Achievement of overall flame flashback resistance is considered a highly difficult task especially for combustors fed by different fuel blends, as different blends have a different stability limit and hence different flame flashback trends. Increasing flame flashback resistance against one flashback mechanism could result in some drawbacks for other mechanisms impedance.

Thus enhancement system performance to have high flashback resistance for both CIVB and BLF requires proper management for the interaction between combustor geometry and flow field aerodynamics. However, controlling such interaction is such a complicated process and needs accurate experimental and numerical investigations to obtain the most practical and reasonable configurations that can fulfil flame stability requirements.

The main outcomes of this chapter can be summarised as follows:

- 1. Using central fuel injectors or axial air injection to improve flame flashback resistance against CIVB could lead to increase BLF, whereas such stabilisation techniques enforce the flame to propagation via wall boundary layer.
- 2. Increasing resistance against boundary layer flame flashback needs enhancement of the surface characteristics of the nozzle interior wall in such a way that those geometrical modifications lead to increase the velocity gradient close to the nozzle wall to be higher than critical velocity gradient, hence avoid being overcome by turbulent flame speed in this region.
- 3. Two woven steel micromesh configurations 50 µm and 150 µm have used to investigate their effect on velocity gradient and hence flame flashback trends; those surface modifications have been analysed and investigated numerically by other work.
- 4. Results showed that micromesh configurations have very good potentials in terms of increasing flame flashback resistance, this observed visually and through increasing operation stability margins compared with nozzle base case (nozzle wall without micromesh).
- Using both air injection and micromesh configurations increase the stability limits considerably, stability margins increased both in equivalence ratio and inlet tangential velocity, this observed for both micromesh 50 and 150 μm.

However, 50 μ m reveals slightly better performance. This outcome is vital for combustor systems that require increasing their power while keeping constant equivalence ratio. Moreover, such technique provides the possibility of switching from one fuel to another.

6. From an industrial point of view using such techniques is considered as a commercial solution to protect swirl combustors from CIVB and BLF at the same time.



Summary of Discussions

Chapter 8

For every fact, there is an infinity of hypotheses Robert Maynard Pirsig

8.1 Summary of Discussions

Flame flashback is one of the combustion instabilities that occurs when flame propagates from its stable desired position downstream burner nozzle towards the premixing zone. This undesirable phenomenon represents a serious threaten for the operation stability of swirl combustors, and it can cause severe damages to systems hardware, consequently degrading the commercial performance of such systems.

Fuel blends characterised by high flame speeds such as high hydrogen content blends are more likely to be subjected to different flame flashback mechanisms. This problem is considered a serious obstacle in the way of development of gas turbine combustors to switch to a different fuel, whereas using low pollutant blends becomes an urgent desire for low emission requirements. Therefore, to address and understand different flame flashback mechanisms, intensive research and investigations are required. The studies should focus on both the system hardware modifications and flowfield manipulation.

This study has investigated and demonstrated the feasibility of various techniques for flame flashback resistance in gas turbines swirl combustors. The interaction between combustion system hardware and highly complex swirling flows is the key feature that should be taken into consideration for the proper methodologies to tackle various combustion instabilities.

Thus, this study targeted the effects of fuel central fuel injector geometries on flame flashback mechanisms, starting with the combustion-induced vortex breakdown flashback (CIVB). Then the research was expanded to investigate the effect of axial air injection when it physically simulates the central fuel injector. The effect was investigated and demonstrated via characterisation of changing velocity profiles and turbulence intensity which were measured with the aid of a 1D - LDA system. Finally, the effects of the nozzle inside surface or geometry on boundary layer flashback were also characterised.

The most important findings and conclusions of this study are listed below.

- Although they have high flame stability potentials based on the formation of the so-called coherent structures, swirl combustors are frequently subjected to different combustion instabilities. Employment of such combustors to work with lean premixed mode could increase the possibility of these undesirable instabilities.
- Flame flashback which is the upstream flame propagation from a presumably stable operation position towards the premixing zone is considered an inherent instability problem in swirl combustors working on premixed conditions. The main consequences of this phenomenon are the possibility of the system being subjected to severe damages in addition to increasing pollution.
- Flashback can occur in four mechanisms: wall boundary layer flashback (BLF), turbulent core flashback, acoustic combustion instabilities and combustion induce vortex breakdown (CIVB).
- CIVB is considered as one source of sudden flame transition and flashback in swirl combustors, especially in the absence of a centre-body or axial injection. This type of upstream flame propagation can initiate even if the fresh mixture velocity is higher than flame speed. The phenomenon is attributed to small changes in the flow field and interaction between the turbulence and heat release by chemical reaction at the tip of the CRZ.
- Despite intensive studies about flame flashback mechanisms. Investigations on the CIVB are still limited. Whereas this type of flashback represents state of the art research compared to the other three types, the first time that this phenomenon was explicitly observed and

documented was in 2003 by [65, 66]. Since then many hypothesis and theories have been implemented experimentally and numerically in order to achieve more insight understanding on this characteristic flame flashback mechanism. Nevertheless, there is still some ambiguity surrounding the factors governing the CIVB and more investigations in this area are still important and urgent.

- In this study, two techniques have been used to determine their effect on the CIVB, central fuel injectors or centre bodies and axial air injection.
- Central fuel injectors, when used as centre bodies, can affect the flow field characteristics downstream the burner nozzle and suppress the CIVB providing a considerable resistance against it.
- Using the dimensionless number (χ) which represents the ratio of the outside injector diameter to the inside nozzle diameter, and the burner inlet tangential velocity may make the results and correlations more global and capable of being applied on any swirl burner regardless its size.
- > The injector outside diameter can significantly affect flame flashback limits; i.e., at a geometric swirl number Sg=1.12, the reduction of (χ) moves the operation stability limits to the leaner region. However, there is a limit for this reduction, any reduction beyond (χ = 0.09) does not change in flashback limits. At small (χ) or small injector outside diameter, the thin vortex core extends over the fuel injector to the baseplate for all tangential inlet velocities. Flashback can be seen to occur as a radial flashback from the outer boundary of the CRZ, to the radial aligned tangential inlets. Flashback resistance here is weak because it depends on the relatively low radial velocity in the swirl chamber. The CRZ formation is less affected by the injector diameter.
- > Increasing injector diameter, i.e. increasing (χ) up to 0.253 improves flashback limits and provides wider stability operation. At these (χ) values the CRZ/burning vortex region still extends around the fuel injector but not to the baseplate.

- Bigger injector diameters (χ) 0.280 and 0.306 enhances flashback resistance further. The flame front position and stabilisation mechanism are considerably changed. For these configurations, CRZ is significantly affected by inlet tangential velocities. Low tangential velocity leads to unstable CRZ, hence early flashback. However, high tangential velocities result in better flow stabilisation, with the flashback mechanism now occurring via the outer boundary layer between the shear layer and the nozzle wall.
- Reducing the geometric swirl number to Sg=0.9 does not change the trends significantly. The effect of the injector diameter is still almost the same.
- Visual observations revealed that increasing (χ) lead to change the type of flashback mechanism. At low values i.e. low injector diameters the flashback mechanism is CIVB, while with higher values enforce flame propagation via wall i.e. BLF. Thus, there is a critical value of (χ) for flashback transition from CIVB to BLF. This value is $\chi = 0.280$ for Sg=1.12 and $\chi = 0.320$ for Sg= 0.9.
- The proper selection of a central fuel injector diameter can improve flashback resistance in conjunction with optimisation of swirl number at a level beyond initial vortex breakdown, i.e. (S>0.7). This finding can be used for design purposes. When selecting an injector diameter, it must be taken into consideration the required central fuel injection to start burner operation. Moreover, its effect on flame flashback transition from CIVB to BLF is important too. Good geometry selection is the one that can make the desired balance between those flashback mechanisms [87].
- Although central fuel injectors or bluff-bodies have a good potential regarding flame flashback resistance, they also have some drawbacks in this context. Even with the appropriate selection of injector diameter, it is difficult to achieve the entire safe operation, as the risk of being subjected to different flashback mechanisms, especially CIVB, still exists. Operation

stability limits for swirl burners with certain injector diameter are different from those when using the same injector but with a different fuel. Certain injector diameter could have good resistance against the CIVB for certain fuel type. However, the same injector may have porn resistance when different fuel is used.

- Central fuel injectors could undergo life degradation due to significant high-temperature levels as they will be in contact with the flame under the effect of the CIVB. This effect could be increased when high hydrogen content blends are used. Consequently, the maintenance cost will rise.
- Central fuel injection for flame stability demands may lead to increase pollutants level and degradation of the mixing degree.
- Axial air injection, when used for flame stabilisation requirements instead of central fuel injectors, can provide many solutions for problems that can arise when central fuel injectors are used.
- Axial air injection jets simulate the physical shape of the central fuel injector can lead to a considerable variation in the stability map. They produce wider stability operation, and the stability limits increase both regarding equivalence ratio and inlet tangential velocity.
- The possibility of change the air jet diameter can give more flexibility of operation. It is worth mentioning here that central fuel injectors do not have this feature.
- The obtained wide operation region enables the burner to work at higher power than that when the central fuel injector is used, hence it is possible to increase power to higher levels at constant equivalence ratio, whereas this is difficult for burners using central injectors due to the limitation of their stability map.
- The high resistance against flame propagation by the axial air injection promotes the possibility of flame re-stabilisation under flashback conditions. The slightly increasing of axial air injection pushes the CRZ

downstream and prevents the flame from propagating towards the burner plenum, while this is impossible to achieve when the central fuel injector is used.

- The defects in axial velocity at the central vortex core that leads to CIVB can be reduced significantly when axial air injection is used. This jet also leads to a wider vortex core, consequently lower pressure gradient, hence reducing baroclinic torque effects which are considered as one of the factors the provoke CIVB.
- Axial air injection has a tangible effect on all flowfield characteristics downstream burner mouth. This is represented by the considerable reduction of turbulence levels, hence reducing of local turbulent flame speed. Thus this reduction of turbulent flame speed, especially close to the burner exit plane, will enable the upcoming fresh mixture to compensate the velocity defect at the tip of the CRZ providing the necessary velocity balance at this region.
- The air injector position with respect to burner base plate is crucial regarding stability operation and turbulence intensity.
- ► Injecting air directly from the burner baseplate (X=0) keeps the coherence of the flow and is less affected by the inlet tangential flow. This allows the burner to work under reasonable conditions, with the flashback retarded to $\Phi = 0.7$ over inlet tangential velocities ranging from 2.5 to 7.5 m/s.
- > Injecting air at further downstream position (X=29 mm), the air stream became directly subjected to the high momentum tangential flow which in turn produces pressure fluctuations and high turbulence intensity in this region, leading to flashback occurring at lower Φ .
- Injecting air from other downstream positions (X=110 and 150 mm) leads to significant decrease in the turbulence level and hence better flashback resistance, consequently wider stability operation margins. The reason for that is at those positions the air injector became protected by the burner sleeve and hence is less affected by the inlet tangential flow. Moreover,

the flashback resistance for the last two positions is promoted by both the air injection and the injector body.

- Using air injection instead of central fuel injectors can significantly reduce the maintenance cost required when central injectors are used. Moreover, for fuel switching requirements, the wider stability operation map allows wider areas for fuel switching.
- Axial velocity and turbulence intensity measurements in the axial direction by LDA based on instantaneous and bulk flow velocity seem to be useful tools to investigate the effect of different air injection positions on flashback limits. Both techniques give approximately the same results.
- Flow field measurements under the effect of air injection by using 1D LDA reveal a considerable change in turbulence level and velocity gradient not just close to the burner mouth but also for a significant distance downstream. This was evident when pushing the CRZ away from the burner exit plane with interest reduction of turbulence intensity.
- The amount of central air injection with respect to inlet tangential flow is very important and should be chosen carefully. This amount has a direct effect on the size and velocity values of the CRZ and other coherent structures. Thus, the ratio of axial injection should be proportional to the tangential flow according to the required swirl structure. From one hand, high axial injection could reduce the swirl strength to be lower than S=0.6 which is the minimum swirl number required for vortex breakdown. On the other hand, very low axial injection could reduce its potential for flame stabilisation.
- The PVC is affected by the amount of axial injection. The high axial injection can suppress the PVC, and hence reduce the frequency of oscillation which in turn reduce the acoustic combustion instabilities. Moreover, axial air injection can alter the generated vortex breakdown from bubble to cone type which is better for stability.

- Axial air injection at the burner central axis produces a positive axial velocity gradient that promotes CIVB resistance. However, the positive gradient at the centre leads to negative velocity gradients near the nozzle wall, which in turn result in BLF. This effect has been visually observed whereas axial injection and big diameter fuel injectors enforce flame propagation via wall boundary layer.
- In order to achieve the required benefits from using axial air injection, its consequences on BLF must be taken into consideration.
- Using micromeshs proved their potentials for BLF resistance. The structure changes the behaviour of the boundary sublayer in such a way that the fresh mixture velocity is still high even near the wall which in turn increases the velocity gradient greater than a critical value that can overcome by turbulent flame speed.
- The geometry and size of the micromesh are important. Fine micromeshs with grid size 50 μm give better results than coarse ones with a grid size of 150 μm.
- The effect of using a micromesh can be observed visually. The effect leads to an increase of the operation stability map in both cases, with and without air injection. The optimum conditions were when using the two techniques together, the axial air injection and micromesh surfaces.
- For industrial requirements the use of air injection and micromeshes for flame stabilisation requirements seems to be useful from the commercial point of view, the availability of air and possibility of commercial production of such microsurfaces consolidates that. However, there is still considerable research required for prediction of the manufacturing cost of same advanced micromeshs.



Conclusions and Future Work

Chapter 9

Success is a science; if you have the conditions, you will get the results

Oscar Wilde

9.1 Conclusions

The main conclusions of this study can be summarised as follows;

- Central fuel injectors, when used as centre bodies, can affect the flow field characteristics downstream the burner nozzle and suppress the CIVB, providing a considerable resistance against it. Increasing injector outside diameter alter flame flashback mechanism from CIVB to BLF.
- Axial air injection jets, when used for flame stabilisation requirements instead of central fuel injectors, can provide many solutions for problems that can arise when central fuel injectors or bluff bodies are used. They produce wider stability operation regions than that of central fuel injectors. The main effect is pushing the CRZ downstream the nozzle and hence avoid high turbulence that can lead to flame flashback.
- Using both air injection and micromesh configurations increase the stability limits considerably. They provide good flame flashback resistance against BLF and CIVB.

9.2 Future Work

The outcomes and findings of this study could open new horizons of flame flashback resistance techniques. Some other relevant investigations that could not be covered by this study need to be taken into consideration.

9.2.1 Central Fuel Injectors

Despite many studies on the effect of central fuel injectors or bluff-bodies on the performance of swirl burners, this area of study still needs more investigations.

- The geometry of central fuel injectors may have an effect on flame stability different from that of circular types. Thus it is important to investigate other geometry configurations such as square, rectangular and triangular, whereas such geometries could change flow field and the interaction between swirl coherent structures, consequently alter the stability operation map.
- The type of material of central fuel injector may have a considerable effect on the flame stability. The heat exchange between flame and injector material can play an important role in this context. The amount of heat absorbed by the injector and transferred to the downstream coherent structures may affect the mechanism of the bubble formation at the tip of the CRZ that leads finally to CIVB. It has previously demonstrated that volume expansion due to heat release at the tip of the CRZ is crucial in flame propagation and the onset of the CIVB [163].

9.2.2 Axial Air Injection

Research on the use of axial air injection for flame stability are very limited. Thus this technique and its potentials can be considered as state of the art. Therefore more investigations are still important and urgent to achieve a clear view of it.

- > Determine the effect of variable air jet diameter on stability limits;
- Investigate the effect of axial air injection using different fuels, especially high hydrogen blends;
- Study the effect of preheating or cooling, pressurised or atmospheric air, on flame stability and swirl structures.
- Use of other gases to inject axially instead of air or mixes of those gases with air at different mixing rates investigating the impact of those on flame stability.

- Determine experimentally and numerically the difference in turbulent flame speed values with and without air injection to see the effect of variation in velocity gradient generated by axial air injection on flame speed.
- Study the effect of air injection on the concentration of combustion products and correlate that with pollutant products.
- Since this study was based mainly on one-dimensional measurements of swirling flow structures, three-dimensional studies using 3D LDA or PIV are encouraged.
- Study of the PVC at different turbulence values (X positions) to define the impact of turbulence intensity on the structure. Similarly, the impact on the former using micromeshs is required to acknowledge the impact of the stabilisation mechanisms on the PVC.

Appendix A

Test conditions, chapter 5 (isothermal)

Tangential air flowrate (l/min)	Axial air flowrate (l/min)	Average bulk exit axial velocity (m/s)
400	50	2.45
600	50	3.54
800	50	4.63
1000	50	5.72
1200	50	6.81
400	0	2.18
600	0	3.27
800	0	4.36
1000	0	5.45
1200	0	6.54

Appendix B

Test conditions, chapter 5 (combustion)

	Tangential	Tangential	Axial air	Χ	Ф	Average	Power (Kw)
No	flowrate	flowrate	(l/min)	(mm)		exit	
	(l/min)	(l/min)	(1/11111)	(11111)		axial	
	()	()				velocity	
						(m/s)	
1	600	41	50	0	0.59	3.77	22.78
2	800	52	50	=	0.57	4.92	28.90
3	1000	68	50	=	0.60	6.10	37.79
4	800	56	50	=	0.61	4.94	31.12
5	1000	72	50	=	0.64	6.12	40.02
6	800	52	50	=	0.57	4.92	28.90
7	800	55	50	=	0.60	4.94	30.57
8	1000	72	50	=	0.64	6.12	40.02
9	1200	78	50	=	0.58	7.24	43.35
10	1200	82	50	=	0.61	7.26	45.57
11	1200	80	50	=	0.59	7.25	44.46
12	1200	85	50	=	0.63	7.28	47.24
13	600	41	50	29	0.59	3.77	22.78
14	800	52	50	=	0.57	4.92	28.90
15	800	55	50	=	0.60	4.94	30.57
16	1000	66	50	=	0.58	6.09	36.68
17	1000	70	50	=	0.62	6.11	38.90
18	1200	77	50	=	0.57	7.24	42.79
19	1200	80	50	=	0.59	7.25	44.46
20	800	53	50	=	0.58	4.92	29.45
21	800	55	50	=	0.60	4.94	30.57
22	600	45	50	110	0.64	3.79	25.01
23	800	53	50	=	0.58	4.92	29.45
24	800	66	50	=	0.72	5.00	36.68
25	1000	71	50	=	0.63	6.11	39.46
26	1000	79	50	=	0.70	6.16	43.91
27	1200	86	50	=	0.64	7.28	47.80
28	1200	94	50	=	0.70	7.33	52.24
29	400	27	0	150	0.63	2.32	15.00
30	600	36	0	=	0.56	3.46	20.01
31	800	50	0	=	0.58	4.63	27.79
32	800	54	0	=	0.63	4.65	30.01
33	1000	73	0	=	0.68	5.84	40.57

34	400	31	50	=	0.64	2.63	17.23
35	600	45	50	=	0.64	3.79	25.01
36	800	53	50	=	0.58	4.92	29.45
37	800	60	50	=	0.66	4.96	33.35
38	800	66	50	=	0.72	5.00	36.68
39	1000	69	50	=	0.61	6.10	38.35
40	1000	78	50	=	0.69	6.15	43.35
41	1000	83	50	=	0.74	6.18	46.13
42	1200	80	50	=	0.59	7.25	44.46
43	1200	86	50	=	0.64	7.28	47.80
44	1200	100	50	=	0.74	7.36	55.58
45	800	50	66	=	0.72	5.00	36.68

References:

- [1] The Solar Spark. (2016). *Global Warming*. Available: <u>http://www.thesolarspark.co.uk/the-science/the-energy-crisis/global-warming/</u>
- [2] NOAA. (2009). *Climate Indicators* Available: <u>http://www.ncdc.noaa.gov/bams-state-of-the-climate/2009-time-series/land</u>
- [3] NOAA. (2016). *Basics of the Carbon Cycle and the Greenhouse Effect,* . Available: <u>https://www.esrl.noaa.gov/gmd/ccgg/basics.html</u>
- [4] EPA. (2015). *Overview of Greenhouse Gases*. Available: <u>http://www.epa.gov/climatechange/ghgemissions/gases/co2.html</u>
- [5] D. Reay and C. Hogan. (2012). *Greenhouse gas*. Available: <u>http://editors.eol.org/eoearth/wiki/Greenhouse_gas</u>
- [6] NOAA. Global Greenhouse Gas Reference Network [Online]. Available: <u>http://www.esrl.noaa.gov/gmd/ccgg/trends/global.html</u>
- [7] BP, "British petroleum Annual Report and Form " 2014.
- [8] F. Catapano, S. Di Iorio, A. Magno, P. Sementa, and B. M. Vaglieco, "A comprehensive analysis of the effect of ethanol, methane and methane-hydrogen blend on the combustion process in a PFI (port fuel injection) engine," *Energy*, vol. 88, pp. 101-110, 8// 2015.
- [9] F. A. D. Oliveira, J. A. Carvalho, P. M. Sobrinho, and A. de Castro, "Analysis of oxyfuel combustion as an alternative to combustion with air in metal reheating furnaces," *Energy*, vol. 78, pp. 290-297, 2014.
- [10] M. Lo Faro, V. Antonucci, P. L. Antonucci, and A. S. Aricò, "Fuel flexibility: A key challenge for SOFC technology," *Fuel*, vol. 102, pp. 554-559, 2012.
- [11] J. Lewis, A. Valera-Medina, R. Marsh, and S. Morris, "Augmenting the Structures in a Swirling Flame via Diffusive Injection," *Journal of Combustion*, vol. 2014, pp. 1-16, 2014.
- [12] H. S. Zhen, C. W. Leung, C. S. Cheung, and Z. H. Huang, "Characterization of biogashydrogen premixed flames using Bunsen burner," *International Journal of Hydrogen Energy*, vol. 39, pp. 13292-13299, 2014.
- [13] B. J. M. de Vries, D. P. van Vuuren, and M. M. Hoogwijk, "Renewable energy sources: Their global potential for the first-half of the 21st century at a global level: An integrated approach," *Energy Policy*, vol. 35, pp. 2590-2610, 2007.
- [14] VIASPACE. (2015). *Biomass Versus Fossil Fuels, Solar and Wind*. Available: http://www.viaspace.com/biomass_versus_alternatives.php
- [15] A. E. E. Khalil and A. K. Gupta, "Clean combustion in gas turbine engines using Butyl Nonanoate biofuel," *Fuel*, vol. 116, pp. 522-528, 2014.
- [16] R. Jones, J. Goldmeer, and B. Monetti, "Addressing Gas Turbine Fuel Flexibility," 2011.
- [17] L. Rye, S. Blakey, and C. W. Wilson, "Sustainability of supply or the planet: a review of potential drop-in alternative aviation fuels," *Energy & Environmental Science*, vol. 3, pp. 17-27, 2010.
- [18] P. Bajpai and V. Dash, "Hybrid renewable energy systems for power generation in stand-alone applications: A review," *Renewable and Sustainable Energy Reviews*, vol. 16, pp. 2926-2939, 2012.
- [19] P. Sayad, A. Schönborn, and J. Klingmann, "Experimental Investigations of the Lean Blowout Limit of Different Syngas Mixtures in an Atmospheric, Premixed, Variable-Swirl Burner," *Energy & Fuels*, vol. 27, pp. 2783-2793, 2013.
- [20] F. Catapano, S. Di Iorio a, A. Magno, P. Sementa, and B. M. Vaglieco, "A comprehensive analysis of the effect of ethanol, methane and methane-hydrogen blend on the combustion process in a PFI (port fuel injection) engine," *Energy*, 2015.

- [21] E. Ahmadi Moghaddam, S. Ahlgren, C. Hulteberg, and Å. Nordberg, "Energy balance and global warming potential of biogas-based fuels from a life cycle perspective," *Fuel Processing Technology*, vol. 132, pp. 74-82, 2015.
- [22] J. U. Nef, "An Early Energy Crisis and its Consequences," *Scientific American*, vol. 237, 1977.
- [23] BP, "Statistical Review of World Energy," 2013.
- [24] IEA. (2015). *Renewable energy*. Available: <u>http://en.wikipedia.org/wiki/Renewable_energy</u>
- [25] M. Braun-Unkhoff, J. Dembowski, J. Herzler, J. Karle, C. Naumann, and U. Riedel, "Alternative Fuels Based on Biomass: An Experimental and Modeling Study of Ethanol Cofiring to Natural Gas," *Journal of Engineering for Gas Turbines and Power*, vol. 137, p. 091503, 2015.
- [26] K. K. Gupta, A. Rehman, and R. M. Sarviya, "Bio-fuels for the gas turbine: A review," *Renewable and Sustainable Energy Reviews*, vol. 14, pp. 2946-2955, 2010.
- [27] IEA. (2007). *Energy Technology Essentials*. Available: <u>https://www.iea.org/publications/freepublications/publication/iea-energy-</u> technology-essentials-biofuel-production.html
- [28] A. Lantz, R. Collin, M. Aldén, A. Lindholm, J. Larfeldt, and D. Lörstad, "Investigation of Hydrogen Enriched Natural Gas Flames in a SGT-700/800 Burner Using OH PLIF and Chemiluminescence Imaging," *Journal of Engineering for Gas Turbines and Power*, vol. 137, p. 031505, 2014.
- [29] Energy.gov. (2017). *Liquefied Natural Gas*. Available: <u>https://energy.gov/fe/science-innovation/oil-gas/liquefied-natural-gas</u>
- [30] C. Zamfirescu and I. Dincer, "Ammonia as a green fuel and hydrogen source for vehicular applications," *Fuel Processing Technology*, vol. 90, pp. 729-737, 2009.
- [31] Balakrishnan R and V. S. Murugan, "P Series Fuel- The Future Fuel for a Green Environment," *Journal of Chemical and Pharmaceutical Sciences*, pp. 252-254, 2016.
- [32] T. A. Semelsberger, R. L. Borup, and H. L. Greene, "Dimethyl ether (DME) as an alternative fuel," *Journal of Power Sources*, vol. 156, pp. 497-511, 2006.
- [33] T. Lieuwen, V. McDonell, E. Petersen, and D. Santavicca, "Fuel Flexibility Influences on Premixed Combustor Blowout, Flashback, Autoignition, and Stability," *Journal of Engineering for Gas Turbines and Power*, vol. 130, p. 011506, 2008.
- [34] A. Kalantari, E. Sullivan-Lewis, and V. McDonell, "Application of a Turbulent Jet Flame Flashback Propensity Model to a Commercial Gas Turbine Combustor," *Journal of Engineering for Gas Turbines and Power*, vol. 139, pp. 041506-041506-8, 2016.
- [35] T. Lieuwen, V. McDonell, D. Santavicca, and T. Sattelmayer, "Burner Development and Operability Issues Associated with Steady Flowing Syngas Fired Combustors," *Combustion Science and Technology*, vol. 180, pp. 1169-1192, 2008.
- [36] Patrick Tebbe. Gas Turbines [Online]. Available: <u>http://cset.mnsu.edu/engagethermo/components_gasturbine.html</u>
- [37] D. G. Wlison, The Design of High Effeciency Turbomachinary and Gas Turbines, Second Edition ed. A Vicom Company Upper Saddle River, NJ 07458: MIT Press Simon&Schuster, 1998.
- [38] A. MATSUOKA. (2013). Design and Development of A Multistage Axial-Flow Compressor for Industrial Gas Turbine. Available: <u>http://www.jsme-fed.org/newsletters-e/2013_4/no3.html</u>
- [39] A. H.Lefevber and D. R.Ballal, Gas turbine combustion alternative fuels and emissions, Third ed. 6000 Broken Sound Parkway NW, Suite 300 /Boca Raton, FL 33487-2742: CRC Press Taylor & Francis Group, 2010.
- [40] L. K. Smith, Hasan; Etemad, Shahrokh; and Pfefferle, William C.,, *The Gas Turbine Handbook*: U.S. Department of Energy-National Energy Technology Laboratory (NETL), 2006.
- [41] NASA. Combustion Burner [Online]. Available: <u>http://www.grc.nasa.gov/WWW/k-12/airplane/burner.html</u>

- [42] I. V. Kolmanovsky, L. C. Jaw, W. Merrill, and H. T. Van, "Robust control and limit protection in aircraft gas turbine engines," in *2012 IEEE International Conference on Control Applications*, Dubrovnik, Croatia, 2012, pp. 812-819.
- [43] A. Leiserowitz, "International Public Opinion, Perception, and Understanding of Global Climate Change," Yale University2007.
- [44] R. Revelle and H. E. Suess, "Carbon Dioxide Exchange Between Atmosphere and Ocean and the Question of an Increase of Atmospheric CO2 during the Past Decades," *Tellus*, vol. 9, pp. 18-27, 1957.
- [45] D. Clark, "Has the Kyoto protocol made any difference to carbon emissions? ," The Guardian2012.
- [46] European Commission. (2016). *Climate Action*. Available: <u>https://ec.europa.eu/clima/policies/ets_en</u>
- [47] D. G. Wilson, The Design of High Effeciency Turbomachinary and Gas Turbines, Second Edition ed. A Vicom Company Upper Saddle River, NJ 07458: MIT Press Simon&Schuster, 1998.
- [48] R. E. Hayes and S. T. Kolaczkowski, *Introduction to Catalytic Combustion*. Amsterdam B.V: OPA (Overseas Publishers Association), 1997.
- [49] T. Fujii, Y. Ozawa, S. Kikumoto, M. Sato, Y. Yuasa, and H. Inoue, "High Pressure Test Results of a Catalytic Combustor for Gas Turbine," *Journal of Engineering for Gas Turbines and Power*, vol. 120, pp. 509-513, 1996.
- [50] Y. Ozawa, Y. Tochihara, N. Mori, I. Yuri, T. Kanazawa, and K. Sagimori, "High Pressure Test Results of a Catalytically Assisted Ceramic Combustor for a Gas Turbine," *Transactions of the ASME*, vol. 121, pp. 422-428, 1999.
- [51] S. Wang, L. Chen, F. Niu, D. Chen, L. Qin, X. Sun, *et al.*, "Catalytic combustion of hydrogen for residential heat supply application," *International Journal of Energy Research*, vol. 40, pp. 1979-1985, 2016.
- [52] Y. Huang and V. Yang, "Dynamics and stability of lean-premixed swirl-stabilized combustion," *Progress in Energy and Combustion Science*, vol. 35, pp. 293-364, 2009.
- [53] G. Baumgartner, L. R. Boeck, and T. Sattelmayer, "Experimental Investigation of the Transition Mechanism From Stable Flame to Flashback in a Generic Premixed Combustion System With High-Speed Micro-Particle Image Velocimetry and Micro-PLIF Combined With Chemiluminescence Imaging," *Journal of Engineering for Gas Turbines and Power*, vol. 138, pp. 021501-10, 2016.
- [54] M. Utschick and T. Sattelmayer, "Flame Holding in the Premixing Zone of a Gas Turbine Model Combustor After Forced Ignition of H2–Natural Gas–Air Mixtures," *Journal of Engineering for Gas Turbines and Power*, vol. 139, pp. 041504-041504-10, 2016.
- [55] J. A. Wagner, M. W. Renfro, and B. M. Cetegen, "Premixed jet flame behavior in a hot vitiated crossflow of lean combustion products," *Combustion and Flame*, vol. 176, pp. 521-533, 2// 2017.
- [56] D. Dunn-Rankin, M. M. Miyasato, and T. K. Pham, "Chapter 1 Introduction and Perspectives," in *Lean Combustion*, ed Burlington: Academic Press, 2008, pp. 1-18.
- [57] L. Rosentsvit, Y. Levy, V. Erenburg, V. Sherbaum, V. Ovcharenko, B. Chudnovsky, *et al.*, "Extension of the Combustion Stability Range in Dry Low NOxLean Premixed Gas Turbine Combustor Using a Fuel Rich Annular Pilot Burner," *Journal of Engineering for Gas Turbines and Power*, vol. 136, p. 051509, 2014.
- [58] M. P. Boyce, *Gas Turbine Engineering Handbook*, Second ed. Houston Texas: Buttewarth-Heinemann, 2002.
- [59] M. Abdulsada, N. Syred, P. Bowen, T. O'Doherty, A. Griffiths, R. Marsh, *et al.*, "Effect of exhaust confinement and fuel type upon the blowoff limits and fuel switching ability of swirl combustors," *Applied Thermal Engineering*, vol. 48, pp. 426-435, 2012.

- [60] A. Valera-Medina, N. Syred, and A. Griffiths, "Visualisation of isothermal large coherent structures in a swirl burner," *Combustion and Flame*, vol. 156, pp. 1723-1734, 2009.
- [61] N. Syred, C. Wong, R.-M. v, J. Dawson, and R. Kelso, "Characterisation of the Occurrence of the Precessing Vortex Core in Partially Premixed and Non-Premixed Swirling Flow," presented at the 12th International Symposium on Applications of Laser Techniques to Fluid Mechanics, (12th : 2004 : Lisbon, Portugal), 2004.
- [62] J. O'Connor, V. Acharya, and T. Lieuwen, "Transverse combustion instabilities: Acoustic, fluid mechanic, and flame processes," *Progress in Energy and Combustion Science*, vol. 49, pp. 1-39, 2015.
- [63] S. Candel, D. Durox, T. Schuller, J.-F. Bourgouin, and J. P. Moeck, "Dynamics of Swirling Flames," *Annual Review of Fluid Mechanics*, vol. 46, pp. 147-173, 2014.
- [64] S. Zhu and S. Acharya, "Flame Dynamics With Hydrogen Addition at Lean Blowout Limits," *Journal of Engineering for Gas Turbines and Power*, vol. 136, p. 051506, 2014.
- [65] J. Fritz, M. Kröner, and T. Sattelmayer, "Flashback in a Swirl Burner With Cylindrical Premixing Zone," *Journal of Engineering for Gas Turbines and Power*, vol. 126, p. 276, 2004.
- [66] M. Kröner, J. Fritz, and T. Sattelmayer, "Flashback Limits for Combustion Induced Vortex Breakdown in a Swirl Burner," *Journal of Engineering for Gas Turbines and Power*, vol. 125, p. 693, 2003.
- [67] C. Eichler and T. Sattelmayer, "Premixed flame flashback in wall boundary layers studied by long-distance micro-PIV," *Experiments in Fluids*, vol. 52, pp. 347-360, 2011.
- [68] V.-Z. M. Osvaldo and V.-M. Agustín, "Flashback Avoidance in Swirling Flow Burners," Ingeniería Investigación y Tecnología, volumen XV (número 4), octubrediciembre 2014: 603-614, vol. XV, 2014.
- [69] D. Ebi and N. T. Clemens, "Experimental investigation of upstream flame propagation during boundary layer flashback of swirl flames," *Combustion and Flame*, vol. 168, pp. 39-52, 2016.
- [70] T. G. Reichel, S. Terhaar, and O. Paschereit, "Increasing Flashback Resistance in Lean Premixed Swirl-Stabilized Hydrogen Combustion by Axial Air Injection," *Journal of Engineering for Gas Turbines and Power*, vol. 137, pp. 071503-071503-9, 2015.
- [71] B. Dam, G. Corona, M. Hayder, and A. Choudhuri, "Effects of syngas composition on combustion induced vortex breakdown (CIVB) flashback in a swirl stabilized combustor," *Fuel*, vol. 90, pp. 3274-3284, 2011.
- [72] P. Sayad, A. Schönborn, M. Li, and J. Klingmann, "Visualization of Different Flashback Mechanisms for H2/CH4Mixtures in a Variable-Swirl Burner," *Journal of Engineering for Gas Turbines and Power*, vol. 137, p. 031507, 2015.
- [73] N. Syred, M. Abdulsada, A. Griffiths, T. O'Doherty, and P. Bowen, "The effect of hydrogen containing fuel blends upon flashback in swirl burners," *Applied Energy*, vol. 89, pp. 106-110, 1// 2012.
- [74] N. Syred, "A review of oscillation mechanisms and the role of the precessing vortex core (PVC) in swirl combustion systems," *Progress in Energy and Combustion Science*, vol. 32, pp. 93-161, 2006.
- [75] V. Acharya and T. Lieuwen, "Premixed Flame Response to Helical Flow Disturbances: Non-linear Effects," presented at the 9th U. S. National Combustion Meeting organized by the Cincinnati, Ohio, 2015.
- [76] C. Jeong, J. Bae, T. Kim, J. Yoon, S. Joo, and Y. Yoon, "Investigation of flashback characteristics coupled with combustion instability in turbulent premixed bluff body flames using high-speed OH-PLIF and PIV," *Proceedings of the Combustion Institute*, 2016.
- [77] A. J. De Rosa, S. J. Peluso, B. D. Quay, and D. A. Santavicca, "The Effect of Confinement on the Structure and Dynamic Response of Lean-Premixed, Swirl-

Stabilized Flames," *Journal of Engineering for Gas Turbines and Power*, vol. 138, pp. 061507-061507-10, 2016.

- [78] M. S. Klassen, "Fuel Flexibility for Dry Low Emission Gas Turbines Cleanly Burning Biofuels, Coal Liquids and Petroleum Fuels," presented at the PowerGen International 2007 December 10-14, 2007, New Orleans, LA, 2007.
- [79] J. E. Temme, P. M. Allison, and J. F. Driscoll, "Combustion instability of a lean premixed prevaporized gas turbine combustor studied using phase-averaged PIV," *Combustion and Flame*, vol. 161, pp. 958-970, 2014.
- [80] Y. Yan, L. Dang, Y. Deng, J. Li, and J. Zhao, "Experimental study of flow dynamics and fuel spray characteristics in Lean Premixed Prevaporized Combustor," *Fuel*, vol. 144, pp. 197-204, 2015.
- [81] J. C. Abanades, E. J. Anthony, J. Wang, and J. E. Oakey, "Fluidized Bed Combustion Systems Integrating CO2 Capture with CaO," *Environmental Science & Technology*, vol. 39, pp. 2861-2866, 2005/04/01 2005.
- [82] L. Jia, Y. Tan, C. Wang, and E. J. Anthony, "Experimental Study of Oxy-Fuel Combustion and Sulfur Capture in a Mini-CFBC," *Energy & Fuels*, vol. 21, pp. 3160-3164, 2007.
- [83] S. Li, W. Li, M. Xu, X. Wang, H. Li, and Q. Lu, "The experimental study on nitrogen oxides and SO2 emission for oxy-fuel circulation fluidized bed combustion with high oxygen concentration," *Fuel*, vol. 146, pp. 81-87, 2015.
- [84] D. Kunii and O. Levenspiel, *Fluidization Engineering*, Second Edition ed. USA: Butterworth-Heineman, 1991.
- [85] E. J. Anthony, "Fluidized bed combustion of alternative solid fuels; status, successes and problems of the technology," *Progress in Energy and Combustion Science*, vol. Volume 21, Issue 3, 1995, Pages, pp. 239–268, 1995.
- [86] H. D. Ross, *Microgravity Combustion: Fire in Free Fall*. London, UK: Academic Press, London, U.K, 2001.
- [87] T. Sattelmayer, C. Mayer, and J. Sangl, "Interaction of Flame Flashback Mechanisms in Premixed Hydrogen–Air Swirl Flames," *Journal of Engineering for Gas Turbines and Power*, vol. 138, p. 011503, 2016.
- [88] N. Peters, *Turbulent Combustion*. Cambridge: Cambridge University Press, 2000.
- [89] T. García-Armingol, A. Sobrino, and J. Ballester, "Stability ranges of fully and partially premixed syngas flames," in *Proceedings of the European Combustion Meeting 2013*, 2013.
- [90] Y. Nada, K. Matsumoto, and S. Noda, "Liftoff heights of turbulent non-premixed flames in co-flows diluted by CO2/N2," *Combustion and Flame*, vol. 161, pp. 2890-2903, 2014.
- [91] V. K. Arghode, A. K. Gupta, and K. M. Bryden, "High intensity colorless distributed combustion for ultra low emissions and enhanced performance," *Applied Energy*, vol. 92, pp. 822-830, 2012.
- [92] T. Kim and Y. Kim, "Interactive transient flamelet modeling for soot formation and oxidation processes in laminar non-premixed jet flames," *Combustion and Flame*, vol. 162, pp. 1660-1678, 2015.
- [93] K. Oh and H. Shin, "The effect of oxygen and carbon dioxide concentration on soot formation in non-premixed flames," *Fuel*, vol. 85, pp. 615-624, 2006.
- [94] F. LIU, H. GUO, G. J. SMALLWOOD, and O"MERL.GU"LDER, "The Chemical Effects of Carbon Dioxide as an Additive in an Ethylene Diffusion Flame: Implications for Soot and NOx Formation," *Combustion and Flame*, vol. 125, pp. 778–787, 2001.
- [95] S. P. Burke and T. E. W. Schumann, "Diffusion Flames," *Industrial & Engineering Chemistry*, vol. 20, pp. 998-1004, 1928.
- [96] A. Valera-Medina, N. Syred, P. Kay, and A. Griffiths, "Central recirculation zone analysis in an unconfined tangential swirl burner with varying degrees of premixing," *Experiments in Fluids*, vol. 50, pp. 1611-1623, 2011.

- [97] O. Lucca-Negro and T. O'Doherty, "Vortex breakdown: a review," *Progress in Energy and Combustion Science*, vol. 27, pp. 431-481, // 2001.
- [98] D. Durox, J. P. Moeck, J.-F. Bourgouin, P. Morenton, M. Viallon, T. Schuller, *et al.*, "Flame dynamics of a variable swirl number system and instability control," *Combustion and Flame*, vol. 160, pp. 1729-1742, 2013.
- [99] A. Coghe, G. Solero, and G. Scribano, "Recirculation phenomena in a natural gas swirl combustor," *Experimental Thermal and Fluid Science*, vol. 28, pp. 709-714, 2004.
- [100] A. Valera-Medina, " Coherent Structures and their Effects on Processes Occurring in Swirl Combustors," PhD, School of Engineering / Cardiff University /UK, 2009.
- [101] A. J. Griffiths, N. Syred, and W. Fick, "A review of Biomass and Associated Work at Cardiff Relating to Small Scale Heat and Power Systems," *IFRF Combustion Journal*, pp. 1-40, 2000.
- [102] N. Syred and J. M. Beér, "Combustion in swirling flows: A review," *Combustion and Flame*, vol. 23, pp. 143-201, 1974/10/01 1974.
- [103] S. Leibovich, "The Structure of Vortex Breakdown," Annual Review of Fluid Mechanics, vol. 10, pp. 221-246, 1978.
- [104] A. M. Steinberg, C. M. Arndt, and W. Meier, "Parametric study of vortex structures and their dynamics in swirl-stabilized combustion," *Proceedings of the Combustion Institute*, vol. 34, pp. 3117-3125, 2013.
- [105] A. K. Gupta, D. G. Lilley, and N. Syred, *Swirl flows* Tunbridge Wells : Abacus, 1984.
- [106] D. M. Wicksall, A. K. Agrawal, R. W. Schefer, and J. O. Keller, "The interaction of flame and flow field in a lean premixed swirl-stabilized combustor operated on H2/CH4/air," *Proceedings of the Combustion Institute*, vol. 30, pp. 2875-2883, 2005.
- [107] X. Lu, S. Wang, H.-G. Sung, S.-Y. Hsieh, and V. Yang, "Large-eddy simulations of turbulent swirling flows injected into a dump chamber," *Journal of Fluid Mechanics*, vol. 527, pp. 171-195, 2005.
- [108] C. W. Foley, I. Chterev, J. Seitzman, and T.Lieuwen, "Flame Configurations in a Lean Premixed Dump Combustor with an Annular Swirling Flow," in US National Combustion Meeting Organized by the Eastern States Section of the Combustion Institute and Hosted by the Georgia Institute of Technology, Atlanta, GA, 2011.
- [109] T. Sarpkaya, "On stationary and travelling vortex breakdowns," *J. Fluid Mech*, vol. vol. 45, pp. 545-559, 1971.
- [110] P. Yazdabadi, A. J. Griffiths, and N. Syred, "Investigations into the Precessing Vortex Core Phenomenon in Cyclone Dust Separators," *Proceedings of the Institution of Mechanical Engineers, Part E: Journal of Process Mechanical Engineering*, vol. 208, pp. 147-154, 1994/08/01 1994.
- [111] W. Fick, "Characterisation and Effects of the Precessing Vortex Core.," PhD, Cardiff University: Wales, UK., Cardiff University: Wales, UK, Cardiff University: Wales, UK., 1998.
- [112] S. Terhaar, "Identification and Modeling of Coherent Structures in Swirl-Stabilized Combustors at Dray and SteamDiluted Conditions," PhD, Technische Universität Berlin 2015.
- [113] D. Galley, S. Ducruix, F. Lacas, and D. Veynante, "Mixing and stabilization study of a partially premixed swirling flame using laser induced fluorescence," *Combustion and Flame*, vol. 158, pp. 155-171, 2011.
- [114] J. M. Beer and N. A. Chigier, *Combustion Aerodynamics*. London: Applied Science Publishers LTD, 1972.
- [115] M. H. Abdulsada, "Flashback and Blowoff Characteristics of Gas Turbine Swirl Combustor," PhD, Institute of Energy, Cardiff School of Engineering, Cardiff University, Cardiff, UK, Cardiff University / Cardiff / UK, 2011.
- [116] P. Palies, D. Durox, T. Schuller, and S. Candel, "The combined dynamics of swirler and turbulent premixed swirling flames," *Combustion and Flame*, vol. 157, pp. 1698-1717, 2010.

- [117] Y. M. Al-Abdeli and A. R. Masri, "Recirculation and flowfield regimes of unconfined non-reacting swirling flows," *Experimental Thermal and Fluid Science*, vol. 27, pp. 655-665, 2003.
- [118] A. C. Benim and K. J. Syed, *Flashback Mechanisms in Lean Premixed Gas Turbine Combustion*. Boston: Academic Press, 2015.
- [119] M. D. Durbin and D. R. Ballal, "Studies of Lean Blowout in a Step Swirl Combustor," *Journal of Engineering for Gas Turbines and Power*, vol. 118, pp. 72-77, 1996.
- [120] S. J. Shanbhogue, S. Husain, and T. Lieuwen, "Lean blowoff of bluff body stabilized flames: Scaling and dynamics," *Progress in Energy and Combustion Science*, vol. 35, pp. 98-120, 2009.
- [121] C. W. Foley, J. Seitzman, and T. Lieuwen, "Analysis and Scalings of Blowoff Limits of 2D and Axisymmetric Bluff Body Stabilized Flames," presented at the Proceedings of ASME Turbo Expo, June 11-15, 2012, Copenhagen, Denmark, 2012.
- [122] A. A. Chaparro and B. M. Cetegen, "Blowoff characteristics of bluff-body stabilized conical premixed flames under upstream velocity modulation," *Combustion and Flame*, vol. 144, pp. 318-335, 2006.
- [123] S. Chaudhuri and B. M. Cetegen, "Blowoff characteristics of bluff-body stabilized conical premixed flames with upstream spatial mixture gradients and velocity oscillations," *Combustion and Flame*, vol. 153, pp. 616-633, 2008.
- [124] S. G. Tuttle, S. Chaudhuri, S. Kostka, K. M. Kopp-Vaughan, T. R. Jensen, B. M. Cetegen, *et al.*, "Time-resolved blowoff transition measurements for two-dimensional bluff body-stabilized flames in vitiated flow," *Combustion and Flame*, vol. 159, pp. 291-305, 2012.
- [125] N. Syred, A. Giles, J. Lewis, A. Valera-Medina, P. Bowen, and A. Griffiths, "Tangential Velocity Effects and Correlations for Blow-Off and Flashback in a Generic Swirl Burner and the effect of a Hydrogen containing Fuel," presented at the 51st AIAA Aerospace Science Meeting, Texas, USA, 2013.
- [126] D. G. Pugh, P. J. Bowen, R. Marsh, A. P. Crayford, J. Runyon, S. Morris, *et al.*, "Dissociative influence of H2O vapour/spray on lean blowoff and NOx reduction for heavily carbonaceous syngas swirling flames," *Combustion and Flame*, vol. 177, pp. 37-48, 3// 2017.
- [127] K. S. Kedia and A. F. Ghoniem, "The blow-off mechanism of a bluff-body stabilized laminar premixed flame," *Combustion and Flame*, vol. 162, pp. 1304-1315, 4// 2015.
- [128] Q. Zhang, "Lean Blowoff Characteristics of Swirling H2/CO/CH4 flames," PhD, Georgia Institute of Technology, Georgia Institute of Technology, Georgia Institute of Technology, 2008.
- [129] D. R. Noble, Q. Zhang, and T. Lieuwen, "Blowout Measurements in a Syngas-Fired Gas Turbine Combustor," presented at the The 22nd Annual International Pittsburgh Coal Conference Pittsburgh, PA, 2005.
- [130] M. Stöhr, C. M. Arndt, and W. Meier, "Effects of Damköhler number on vortex-flame interaction in a gas turbine model combustor," *Proceedings of the Combustion Institute*, vol. 34, pp. 3107-3115, 2013.
- [131] Q. Zhang, S. J. Shanbhogue, and T. Lieuwen, "Dynamics of Premixed H2/CH4 Flames Under Near Blowoff Conditions," *Journal of Engineering for Gas Turbines* and Power, vol. 132, pp. 111502-111502-8, 2010.
- [132] T. Lieuwen, S. Shanbhogue, S. Khosla, and C. Smith, "Dynamics of Bluff Body Flames Near Blowoff," in *45th AIAA Aerospace Sciences Meeting and Exhibit*, ed: American Institute of Aeronautics and Astronautics, 2007.
- [133] B. Roy Chowdhury and B. M. Cetegen, "Experimental study of the effects of free stream turbulence on characteristics and flame structure of bluff-body stabilized conical lean premixed flames," *Combustion and Flame*, vol. 178, pp. 311-328, 4// 2017.
- [134] Q. Zhang, D. R. Noble, S. J. Shanbhogue, and T. Lieuwen, "PIV Measurements in H2/CH4 Swirling Flames under near Blowoff Conditions," presented at the 5th US

Combustion Meeting, Organized by the Western States Section of the Combustion Institute and Hosted by the University of California at San Diego 2007.

- [135] R. Schefer, "Hydrogen enrichment for improved lean flame stability," *International Journal of Hydrogen Energy*, vol. 28, pp. 1131-1141, 2003.
- [136] H. Baej, A. Valera-Medina, N. Syred, R. Marsh, and P. Bowen, "Blowoff propensity, CRZs and Flow Turbulent structure using a range of Syngas compositions for Gas Turbines," presented at the SusTem international conference, Newcastel University / UK, 7-8 July, 2015.
- [137] J. de Vries and E. L. Petersen, "Autoignition of methane-based fuel blends under gas turbine conditions," *Proceedings of the Combustion Institute*, vol. 31, pp. 3163-3171, 2007.
- [138] A. Lantz, R. Collin, M. Aldén, A. Lindholm, J. Larfeldt, and D. Lörstad, "Investigation of Hydrogen Enriched Natural Gas Flames in a SGT-700/800 Burner Using OH PLIF and Chemiluminescence Imaging," *Journal of Engineering for Gas Turbines and Power*, vol. 137, p. 031505, 2015.
- [139] Y. Zhang, X. Jiang, L. Wei, J. Zhang, C. Tang, and Z. Huang, "Experimental and modeling study on auto-ignition characteristics of methane/hydrogen blends under engine relevant pressure," *International Journal of Hydrogen Energy*, vol. 37, pp. 19168-19176, 2012.
- [140] B. K. Dam, G. Corona, M. M. Hayder, and A. R. Choudhuri, "An Experimental Investigation of Combustion Induced Vortex Breakdown Flashback in a Swirl Stabilized Burner," presented at the 49th AIAA Aerospace Sciences Meeting including the New Horizons Forum and Aerospace Exposition, 4 - 7 January 2011, Orlando, Florida, 2011.
- [141] F. Kiesewetter, M. Konle, and T. Sattelmayer, "Analysis of Combustion Induced Vortex Breakdown Driven Flame Flashback in a Premix Burner With Cylindrical Mixing Zone," *Journal of Engineering for Gas Turbines and Power*, vol. 129, p. 929, 2007.
- [142] M. Abdulsada, N. Syred, A. Griffiths, and P. Bowen, "Effects of Swirl Number and Fuel Type Upon the Flashback in Swirl Combustors," ed: the American Institute of Aeronautics and Astronautics, Inc., with permission., 2011.
- [143] B. Shaffer, Z. Duan, and V. McDonell, "Study of Fuel Composition Effects on Flashback Using a Confined Jet Flame Burner," *Journal of Engineering for Gas Turbines and Power*, vol. 135, p. 011502, 2013.
- [144] A. Nauert, P. Petersson, M. Linne, and A. Dreizler, "Experimental analysis of flashback in lean premixed swirling flames: conditions close to flashback," *Experiments in Fluids*, vol. 43, pp. 89-100, 2007.
- [145] S. Burmberger and T. Sattelmayer, "Optimization of the Aerodynamic Flame Stabilization for Fuel Flexible Gas Turbine Premix Burners," *Journal of Engineering* for Gas Turbines and Power, vol. 133, p. 101501, 2011.
- [146] C. Heeger, R. L. Gordon, M. J. Tummers, T. Sattelmayer, and A. Dreizler, "Experimental analysis of flashback in lean premixed swirling flames: upstream flame propagation," *Experiments in Fluids*, vol. 49, pp. 853-863, 2010.
- [147] B. Lewis and G. von Elbe, "Stability and Structure of Burner Flames," *The Journal of Chemical Physics*, vol. 11, p. 75, 1943.
- [148] C. Eichler and T. Sattelmayer, "Experiments on Flame Flashback in a Quasi-2D Turbulent Wall Boundary Layer for Premixed Methane-Hydrogen-Air Mixtures," *Journal of Engineering for Gas Turbines and Power*, vol. 133, p. 011503, 2011.
- [149] Y.-C. Lin, S. Daniele, P. Jansohn, and K. Boulouchos, "Turbulent Flame Speed as an Indicator for Flashback Propensity of Hydrogen-Rich Fuel Gases," *Journal of Engineering for Gas Turbines and Power*, vol. 135, p. 111503, 2013.
- [150] A. N. Lipatnikov and J. Chomiak, "Molecular transport effects on turbulent flame propagation and structure," *Progress in Energy and Combustion Science*, vol. 31, pp. 1-73, 2005.

- [151] H. Kobayashi, Y. Otawara, J. Wang, F. Matsuno, Y. Ogami, M. Okuyama, *et al.*, "Turbulent premixed flame characteristics of a CO/H2/O2 mixture highly diluted with CO2 in a high-pressure environment," *Proceedings of the Combustion Institute*, vol. 34, pp. 1437-1445, 2013.
- [152] D. Beerer, V. McDonell, P. Therkelsen, and R. K. Cheng, "Flashback and Turbulent Flame Speed Measurements in Hydrogen/Methane Flames Stabilized by a Low-Swirl Injector at Elevated Pressures and Temperatures," *Journal of Engineering for Gas Turbines and Power*, vol. 136, p. 031502, 2014.
- [153] C. M. Arndt, M. Severin, C. Dem, M. Stöhr, A. M. Steinberg, and W. Meier, "Experimental analysis of thermo-acoustic instabilities in a generic gas turbine combustor by phase-correlated PIV, chemiluminescence, and laser Raman scattering measurements," *Experiments in Fluids*, vol. 56, 2015.
- [154] L. Lei, G. Zhihui, Z. Chengyu, and S. Xiaofeng, "A Passive Method to Control Combustion Instabilities with Perforated Liner," *Chinese Journal of Aeronautics*, vol. 23, pp. 623-630, 2010.
- [155] J. H. Cho and T. Lieuwen, "Laminar premixed flame response to equivalence ratio oscillations," *Combustion and Flame*, vol. 140, pp. 116-129, 2005.
- [156] R. Balachandran, A. P. Dowling, and E. Mastorakos, "Dynamics of bluff-body stabilised flames subjected to equivalence ratio oscillations," in *Proceedings of the European Combustion Meeting 2011*, Cardiff, Great Britain, 2011.
- [157] R. Balachandran, B. O. Ayoola, C. F. Kaminski, A. P. Dowling, and E. Mastorakos, "Experimental investigation of the nonlinear response of turbulent premixed flames to imposed inlet velocity oscillations," *Combustion and Flame*, vol. 143, pp. 37-55, 10// 2005.
- [158] B. T. Zinn and T. C. Lieuwen, "Combustion Instabilities: Basic Concepts," in Combustion Instabilities In Gas Turbine Engines: Operational Experience, Fundamental Mechanisms, and Modeling, ed: American Institute of Aeronautics and Astronautics, 2005, pp. 3-26.
- [159] A. Coker, Y. Neumeier, B. T. Zinn, S. Menon, and T. I. M. Lieuwen, "Active Instability Control Effectiveness in a Liquid Fueled Combustor," *Combustion Science* and Technology, vol. 178, pp. 1251-1261, 2006.
- [160] M. Konle, F. Kiesewetter, and T. Sattelmayer, "Simultaneous high repetition rate PIV–LIF-measurements of CIVB driven flashback," *Experiments in Fluids*, vol. 44, pp. 529-538, 2008.
- [161] E. Tangermann, M. Pfitzner, M. Konle, and T. Sattelmayer, "Large-Eddy Simulation and Experimental Observation of Combustion-Induced Vortex Breakdown," *Combustion Science and Technology*, vol. 182, pp. 505-516, 2010.
- [162] M. Konle, A. Winkler, F. Kiesewetter, J. Wäsle, and T. Sattelmayer, "CIVB Flashback Analysis with Simultaneous and Time Resolved PIV-LIF Measurements," presented at the 13th Int Symp on Applications of Laser Techniques to Fluid Mechanics, Lisbon, Portugal, 2006.
- [163] M. Konle and T. Sattelmayer, "Interaction of heat release and vortex breakdown during flame flashback driven by combustion induced vortex breakdown," *Experiments in Fluids*, vol. 47, pp. 627-635, 2009.
- [164] M. Kroner, T. Sattelmayer, J. Fritz, F. Keisewetter, and C. Hirsch, "Flame propagation in Swirling Flows- Effects of Local extinction on the Combustion Induced Vortex Breakdown," *Combustion Science and Technology* vol. 179, pp. pp. 1385-1416, 2007.
- [165] N. Syred, A. Giles, J. Lewis, M. Abdulsada, A. Valera Medina, R. Marsh, et al., "Tangential Velocity Effects and Correlations for Blow-Off and Flashback in A Generic Swirl Burner and the effect of a Hydrogen containing Fuel," *Applied Energy*, vol. 116, pp. 288-296, 2014.
- [166] B. Dam, N. Love, and A. Choudhuri, "Flashback propensity of syngas fuels," *Fuel*, vol. 90, pp. 618-625, 2011.
- [167] G. Baumgartner and T. Sattelmayer, "Experimental Investigation of the Flashback Limits and Flame Propagation Mechanisms for Premixed Hydrogen-Air Flames in

Non-Swirling and Swirling Flow," in *Proceedings of ASME Turbo Expo*, San Antonio, Texas, USA, 2013, p. V01AT04A010.

- [168] C. Mayer, J. Sangl, T. Sattelmayer, T. Lachaux, and S. Bernero, "Study on the Operational Window of a Swirl Stabilized Syngas Burner Under Atmospheric and High Pressure Conditions," *Journal of Engineering for Gas Turbines and Power*, vol. 134, p. 031506, 2012.
- [169] EIA, "Natural Gas Issues and Trends," Office of Oil and Gas U.S. Department of Energy ,Washington, DC 205851999.
- [170] C. Park, S. Won, C. Kim, and Y. Choi, "Effect of mixing CO2 with natural gashydrogen blends on combustion in heavy-duty spark ignition engine," *Fuel*, vol. 102, pp. 299-304, 2012.
- [171] E. Moniz, "The future of Natural Gas," 2011.
- [172] EIA, " International Energy Outlook 2011. Natural Gas: World Natural Gas Consumption by region, Reference Case.," 2011.
- [173] K. Kim, H. Kim, B. Kim, K. Lee, and K. Lee, "Effect of Natural Gas Composition on the Performance of a CNG Engine," *Oil & Gas Science and Technology - Revue de l'IFP*, vol. 64, pp. 199-206, 2008.
- [174] A. Demirbas, Methane Gas Hydrate /Chpter 2 Natural Gas: Springer, 2010.
- [175] J. Lewis, R. Marsh, A. Valera-Medina, S. Morris, and H. Baej, "The Use of CO2 to Improve Stability and Emissions of an IGCC Combustor," in ASME Turbo Expo, 2014, p. V04AT04A029.
- [176] R. Slefarski, J. Sacha1, and P. Grzymislawski1, "Combustion of mixtures of biogases and syngases with methane in strong swirl flow," in *Proceedings of the European Combustion Meeting*, 2013.
- [177] P. Chiesa, G. Lozza, and L. Mazzocchi, "Using Hydrogen as Gas Turbine Fuel," *Journal of Engineering for Gas Turbines and Power*, vol. 127, p. 73, 2005.
- [178] A. E. E. Khalil and A. K. Gupta, "Hydrogen addition effects on high intensity distributed combustion," *Applied Energy*, vol. 104, pp. 71-78, 2013.
- [179] F. Wang, Y. Cao, and J. Zhou, "Thermodynamic analysis of high-temperature helium heated fuel reforming for hydrogen production," *International Journal of Energy Research*, vol. 39, pp. 418-432, 2015.
- [180] T. Gül, S. Kypreos, H. Turton, and L. Barreto, "An energy-economic scenario analysis of alternative fuels for personal transport using the Global Multi-regional MARKAL model (GMM)," *Energy*, vol. 34, pp. 1423-1437, 2009.
- [181] M. F. Orhan, I. Dincer, and M. A. Rosen, "Investigation of an integrated hydrogen production system based on nuclear and renewable energy sources: a new approach for sustainable hydrogen production via copper-chlorine thermochemical cycles," *International Journal of Energy Research*, vol. 36, pp. 1388-1394, 2012.
- [182] K. Kaufman, M. Emadi, and A. Ratner, "Effect of Hydrogen Addition to Methane Fuel in a Low Swirl Burner," presented at the 8th U. S. National Combustion Meeting, University of Utah, 2013.
- [183] T. García-Armingol and J. Ballester, "Operational issues in premixed combustion of hydrogen-enriched and syngas fuels," *International Journal of Hydrogen Energy*, vol. 40, pp. 1229-1243, 2014.
- [184] T. Wind, F. Güthe, and K. Syed, "Co-Firing of Hydrogen and Natural Gases in Lean Premixed Conventional and Reheat Burners (Alstom GT26)," in *Proceedings of ASME Turbo Expo*, Düsseldorf, Germany, 2014, p. V04AT04A053.
- [185] N. Syred, A. Giles, J. Lewis, M. Abdulsada, A. Valera Medina, R. Marsh, *et al.*, " Effect of inlet and outlet configurations on blow-off and flashback with premixed combustion for methane and a high hydrogen content fuel in a generic swirl burner," *Applied Energy*, vol. 116, pp. 288-296, 2014.
- [186] DANTEC. (2015). Laser Doppler Anemometry [LDA]. Available: http://www.dantecdynamics.com/measurement-principles-of-lda
- [187] C. T. Eichler, "Flame flashback in wall boundary layer of premixed combustion system," PhD, Technische Universität München Institue für Engineering, Technische

Universität München Institue für Engineering, Technische Universität München Institue für Engineering, 2011.

- [188] J. Sangl, C. Mayer, and T. Sattelmayer, "Dynamic Adaptation of Aerodynamic Flame Stabilization of a Premix Swirl Burner to Fuel Reactivity Using Fuel Momentum," *Journal of Engineering for Gas Turbines and Power*, vol. 133, p. 071501, 2011.
- [189] M. Konle and T. Sattelmayer, "Time Scale Model for the Prediction of the Onset of Flame Flashback Driven by Combustion Induced Vortex Breakdown," *Journal of Engineering for Gas Turbines and Power*, vol. 132, p. 041503, 2010.
- [190] N. Shelil, A. Bagdanavicius, N. Syred, A. Griffiths, and P. Bowen, "Premixed Swirl Combustion and Flashback Analysis with Hydrogen/Methane Mixtures," presented at the 48th AIAA Aerospace Sciences Meeting Including the New Horizons Forum and Aerospace Exposition, 2010.
- [191] N. Karimi, S. McGrath, P. Brown, J. Weinkauff, and A. Dreizler, "Generation of Adverse Pressure Gradient in the Circumferential Flashback of a Premixed Flame," *Flow, Turbulence and Combustion*, vol. 97, pp. 663-687, 2016.
- [192] T. C. Lieuwen, "Flame Stabilization, Flashback, Flameholding, and Blowoff," in Unsteady Combustor Physics, ed Cambridge: Cambridge University Press, 2012, pp. 293-316.
- [193] D. R. Noble, Q. Zhang, A. Shareef, J. Tootle, A. Meyers, and T. Lieuwen, "Syngas Mixture Composition Effects Upon Flashback and Blowout," in *Proceedings of GT2006 ASME Turbo Expo 2006: Power for Land, Sea and Air*, Barcelona, Spain, 2006, pp. 357-368.
- [194] T. G. Reichel and C. O. Paschereit, "Experimental Investigation of Flame Stability of Swirl-Stabilized, Lean Pre-mixed Hydrogen Flames," presented at the 8th International Seminar on Flame structure, Berlin Institute of Technology, Fasanenstr. 89 • 10623 Berlin, 2014.
- [195] R. Szasz, A. A. Subash, A. Lantz, R. Collin, L. Fuchs, and E. Gutmark, "Hysteretic Dynamics of Flashback in a Low-Swirl Stabilized Combustor," *Combustion Science* and Technology, vol. 189, pp. 266-289, 2017.
- [196] S. Taamallah, K. Vogiatzaki, F. M. Alzahrani, E. M. A. Mokheimer, M. A. Habib, and A. F. Ghoniem, "Fuel flexibility, stability and emissions in premixed hydrogenrich gas turbine combustion: Technology, fundamentals, and numerical simulations," *Applied Energy*, vol. 154, pp. 1020-1047, 2015.
- [197] Z. Zhengji, "LDA Application Methods: Laser Doppler Anemometry for Fluid Dynamics (Experimental Fluid Mechanics)," 2010 edition ed. Berlin: Springer, 2010.
- [198] P. Billant, J.-M. Chomaz, and P. Huerre, "Experimental study of vortex breakdown in swirling jets," *Journal of Fluid Mechanics*, vol. 376, pp. 183-219, 1998/12/10 1998.
- [199] R. K. Cheng, D. Littlejohn, W. A. Nazeer, and K. O. Smith, "Laboratory Studies of the Flow Field Characteristics of Low-Swirl Injectors for Adaptation to Fuel-Flexible Turbines," *Journal of Engineering for Gas Turbines and Power*, vol. 130, p. 021501, 2008.
- [200] A. Gruber, J. H. Chen, D. Valiev, and C. K. Law, "Direct numerical simulation of premixed flame boundary layer flashback in turbulent channel flow," *Journal of Fluid Mechanics*, vol. 709, pp. 516-542, 2012.
- [201] R. A. Antonia and P. Å. Krogstad, "Turbulence structure in boundary layers over different types of surface roughness," *Fluid Dynamics Research*, vol. 28, pp. 139-157, 2001.
- [202] F. Dabireau, B. Cuenot, O. Vermorel, and T. Poinsot, "Interaction of flames of H2 + O2 with inert walls," *Combustion and Flame*, vol. 135, pp. 123-133, 2003.
- [203] Shared Labs Europe LTD. (2017). Available: <u>http://www.sharedlabseurope.com/</u>
- [204] M. Al-Fahham, F. A. Hatem, A. S. Alsaegh, A. V. Medina, S. Bigot, and R. Marsh, "Experimental study to enhance resistancefor boundary layer flashback in swirl burners using microsurfaces.," in *Proceedings of ASME Turbo Expo 2017: Turbomachinery Technical Conference and Expansion*, Charlotte, NC, USA, 2017.

[205] F. A. Hatem, N. Syred, A. Valera-Medina, R. Marsh, and P. J. Bowen, "Experimental Investigation of the Effects of Fuel Diffusive Injectors on Premixed Swirling Flames," in *53rd AIAA Aerospace Sciences Meeting*, ed: American Institute of Aeronautics and Astronautics, 2015.