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An amplified dusty star-forming galaxy at $z=6$: unveiling an elusive population of galaxies

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Since their discovery, submillimeter-selected galaxies^{1,2} (SMGs) have revolutionized the field of galaxy formation and evolution. Hundreds of square degrees have been mapped at submillimeter wavelengths³⁻⁵ and notwithstanding the negative K -correction in the submm bands⁶, where there is no significant loss of sensitivity to the detection of these sources up to $z \sim 10$, only a handful of sources have been confirmed to lie at $z > 5$ (ref.⁷⁻¹¹) and only two at $z \geq 6$ (ref.^{12,13}). All of these SMGs are rare examples of extreme starburst galaxies with star formation rates (SFRs) of $\gtrsim 1000 M_{\odot} \text{ yr}^{-1}$ and therefore are not representative of the general population of dusty star-forming galaxies. Consequently, our understanding of the nature of these sources, at the earliest epochs, is still incomplete. Here we report the spectroscopic identification of a gravitation-

ally amplified ($\mu = 9.3 \pm 1.0$) dusty star-forming galaxy at $z = 6.027$. After correcting for gravitational lensing we derive an intrinsic SFR of $380 \pm 50 M_{\odot} \text{ yr}^{-1}$ for this source, and find that its gas and dust properties are similar to those measured for local Ultra Luminous Infrared Galaxies (ULIRGs), extending the local trends up to an unexplored territory at high redshift. This ULIRG-like galaxy at $z = 6$ suggests a universal star-formation efficiency during the last 12.8 Gyr for dusty star-forming galaxies.

HATLAS J090045.4+004125 ($\alpha = 09^{\text{h}}00^{\text{m}}45.8$, $\delta = +00^{\circ}41'23''$; hereafter G09 83808, since it was detected in the GAMA 09hrs field) is part of a sub-sample of the *Herschel* ATLAS ‘500 μm -riser’ galaxies¹⁴ with ultra-red far-infrared (FIR) colours of $S_{500\mu\text{m}}/S_{250\mu\text{m}} > 2$ and $S_{500\mu\text{m}}/S_{350\mu\text{m}} > 1$,

with a flux density threshold of $S_{500\mu\text{m}} < 80$ mJy. The FIR colours of this source are consistent with thermal dust emission redshifted to $z > 4$ and represent a relatively simple selection criterium to find high-redshift galaxies. A similar selection allowed the identification of HFLS3¹², an extreme starburst galaxy (even after corrected for gravitational amplification¹⁵) at $z = 6.3$, in the HerMES blank field survey³.

G09 83808 was observed, among other ultrared-*Herschel* dusty star-forming galaxies, as part of a follow-up program with the Large Millimeter Telescope *Alfonso Serrano* (LMT) using the AzTEC camera, in order to obtain higher angular resolution (~ 8.5 arcsec) continuum observations at 1.1 mm. A sub-sample of those galaxies detected as a single source in the AzTEC images (i.e. with no evidence of multiple components) and with photometric redshifts of > 4 , was selected for spectroscopic observations in the 3 mm band using the Redshift Search Receiver (RSR) on the LMT. In the LMT/RSR spectrum of G09 83808 we identify three emission lines corresponding to $^{12}\text{CO}(6 - 5)$, $^{12}\text{CO}(5 - 4)$, and $\text{H}_2\text{O}(2_{11} - 2_{02})$ (see Fig. 1). Based on these lines we unambiguously determine the galaxy redshift to be $z = 6.0269 \pm 0.0006$ (i.e. when the Universe was just 900 million years old). Follow-up observations with the SMA telescope confirm this solution through the detection of the redshifted [CII] ionized carbon line at 270.35 GHz (Fig. 1).

High-angular resolution observations (0.24 arcsec \times 0.13 arcsec, corresponding to a physical scale of ~ 1 kpc at this redshift) taken with the Atacama Large Millimeter/submillimeter Array (ALMA) at ~ 890 μm reveal a double arc structure (in a partial Einstein ring configuration of radius ~ 1.4 arcsec) around a foreground galaxy at $z = 0.776$ (see Fig. 2), implying strong gravitational amplification of the high-redshift background galaxy. Using these ALMA continuum observations to constrain the effects of gravitational lensing, modelling directly the visibilities in the uv plane (see Methods section for additional details), we derive a gravitational amplifi-

cation factor of $\mu = 9.3 \pm 1.0$. This amplification factor is used to derive the intrinsic physical properties of G09 83808.

Using the *Herschel* 250, 350, and 500 μm photometry¹⁴, combined with the SCUBA-2 850 μm ¹⁴ imaging and our AzTEC 1.1 mm observations (see Table 1), we model the continuum spectral energy distribution (SED; see Figure 3). We estimate an infrared (IR, 8 – 1000 μm) luminosity, L_{IR} , of $3.8 \pm 0.5 \times 10^{12}$ L_{\odot} (corrected for gravitational magnification) which implies a SFR^a of 380 ± 50 $M_{\odot}\text{yr}^{-1}$ (see Methods section for more information). This implies that G09 83808 is a member of the Ultra Luminous Infrared Galaxy (ULIRGs¹⁶) population. This is the only SMG with an unambiguous spectroscopic redshift in this luminosity range at $z \gtrsim 5$, lying between the extreme obscured starbursts^{7-9,12,13} ($\gtrsim 1000$ $M_{\odot}\text{yr}^{-1}$) discovered at submm wavelengths and the UV/optical selected star-forming galaxies with follow-up detections at submm wavelengths¹⁷⁻¹⁹ ($\lesssim 100$ $M_{\odot}\text{yr}^{-1}$).

Although these galaxies are unreachable with the current generation of submm wide-area surveys^{3,4} without the benefit of gravitational amplification, they can be found in the deepest surveys recently achieved with ground-based telescopes, such as the James Clerk Maxwell Telescope (JCMT) SCUBA-2 Cosmology Legacy Survey (S2CLS). However, none of them has yet been spectroscopically confirmed. With the caveat of using the position of the dust SED peak as an estimation of redshift, a study based on S2CLS observations⁵ has derived a comoving space density of 3.2×10^{-6} Mpc^{-3} for sources with $300 < \text{SFR} < 1000$ $M_{\odot}\text{yr}^{-1}$ at $5 < z \lesssim 6$ (i.e. in the range probed by our galaxy). With a duty-cycle correction of ≈ 40 Myr, as the gas depletion time scale measured for G09 83808 (see below) and other galaxies¹², we estimate the corrected comoving space density of this population of galaxies to be $\approx 2 \times 10^{-5}$ Mpc^{-3} , which perfectly matches that of massive quiescent galaxies at $z \approx 3 - 4$ (refs.²⁰). This suggests, that these ULIRG-type galaxies at $5 \lesssim z \lesssim 6$ are the progenitors of these quiescent

^aHere we will use SFR to refer to the dust-obscured SFR

galaxies, which cannot be explained only by the rare extreme starburst (like HFLS3), since they are an order of magnitude less abundant¹⁴.

Based on the CO lines detected in the LMT/RSR spectrum we derive a molecular gas mass of $M(\text{H}_2) = 1.6 \pm 0.6 \times 10^{10} M_\odot$ (see Methods section for details). This implies a gas depletion timescale of $M(\text{H}_2)/\text{SFR} \approx 40$ Myr, consistent with the value found for other SMGs at lower redshifts with ULIRG-luminosity²². G09 83808 shows a remarkable large gas mass fraction of $f_{\text{gas}} = M_{\text{H}_2}/M_{\text{dyn}} \sim 60\%$ (see Methods section), among the largest measured for star forming galaxies at $z \approx 2 - 3$ (ref. 21). The CO(6-5)/CO(5-4) line luminosity ratio of 0.4 ± 0.1 is in agreement with local ULIRGs (although lower than the average²⁴), and implies a CO ladder peaking at $J \leq 5$ (i.e. less excited than AGN-dominated galaxies²³). These two CO transitions, as well as the H₂O line, lie (within the error bars) on their respective FIR/IR-line luminosity relations ($L_{\text{FIR}} \propto L_{\text{CO}(6-5)}^{0.93}$, $L_{\text{FIR}} \propto L_{\text{CO}(5-4)}^{0.97}$, and $L_{\text{H}_2\text{O}} \propto L_{\text{IR}}^{1.16}$) found for local ULIRGs and lower redshifts SMGs^{24,25}. The star-formation efficiency (SFE) of our galaxy, estimated through the $L'_{\text{CO}} - L_{\text{IR}}$ relation (which describes the relationship between the luminosity due to star formation and the gas content), is similar to local (U)LIRGs (see Fig. 4). This suggests a universal SFE across several decades of molecular gas masses from $z = 6$ to $z \sim 0$ (i.e., during the last 12.8 Gyr of the Universe) for this kind of galaxies (although some works²⁶ have reported a slight evolution of the SFE with redshift). In addition, the estimated dust mass of $M_d = 1.9 \pm 0.4 \times 10^8 M_\odot$ results in a gas-to-dust ratio, δ_{GDR} , of 80 ± 30 . This is in agreement with the value estimated for HFLS3¹² and also with local (U)LIRGs²⁷ ($\delta_{\text{GDR}} = 120 \pm 28$).

The luminosity of the [CII] ionized carbon line detected with the SMA is $1.3 \pm 0.4 \times 10^9 L_\odot$ which corresponds to a [CII]/FIR ratio of $3.4 \pm 1.1 \times 10^{-4}$, a value that is among the lowest measured for local (U)LIRGs and SMGs. As shown in Figure 4, our source follows the same [CII] deficiency trend measured for local LIRGs²⁸ extending it to $L_{\text{FIR}} \gtrsim 10^{12} L_\odot$ and up to $z = 6$. The [CII]/FIR ratio of G09

83808 is also consistent with the lowest values measured for lower-redshift SMGs and lies on a region where SMGs and AGN-host galaxies converge (Fig. 4). It may be the case that other SMGs suffer from gravitational amplification, which could help to reduce the large scatter since many of these galaxies should fall along the LIRG relation when corrected for magnification. However, the intrinsic scatter in the relation is high²⁸, even for the local sample, and therefore, larger samples of SMGs are required to derive conclusions about the origin of the [CII] deficiency.

We confirm the existence of ULIRG-like galaxies within the first billion years of Universe's history. These sources may be more representative of the dusty star-forming galaxy population at these epochs than the extreme starbursts previously discovered. Four emission-line-selected galaxies with similar luminosities and redshifts have been recently found around quasars²⁹ (with the caveat of using just one line for redshift determination), however, the properties of these sources may be affected by the companion quasar and therefore not representative of the whole population. Although G09 83808 shows similar properties to those measured in lower-redshift SMGs, its higher dust temperature ($T_d = 49 \pm 3$ K) and compact morphology ($R_{1/2} = 0.6 \pm 0.1$ kpc) resemble that of local ULIRGs. For comparison, typical UV/optically-selected star-forming galaxies at $z \sim 6$ have SFRs ~ 10 times lower and radii ~ 1.7 times larger than G09 83808³⁰. This study is hence crucial for understanding the evolutionary path of SMGs and their link with local galaxies. Although a larger sample is needed to statistically estimate the properties of these sources and their contribution to the cosmic star formation history, this galaxy suggests that star formation in dusty star-forming galaxies has been driven by similar physical processes during the last ~ 12.8 Gyr.

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Author Contributions JAZ led the scientific analysis and the writing of the paper, as well as the SMA follow-up proposal. RJI, EV, SE, AC, HD, JSD, LD, MJM, SS, MWLS, and PW have contributed to the original *Herschel* proposals and source selection of the red sources, where this source was originally identify. AM, DHH, EV, IA, VAR, MC, DRG, ET, and OV performed the selection of the sample for the LMT observations and lead the LMT proposals. MSY, GN, GW, DSA, AV, and MZ carried out LMT data reduction and interpretation. DW, MY, and AIGR assisted with the SMA observations and data reduction. JS, IO, HN have contributed to the data analysis and to fitting and modeling the results. All the authors have

discussed and contributed to this manuscript.

Competing Interests The authors declare that they have no competing financial interests.

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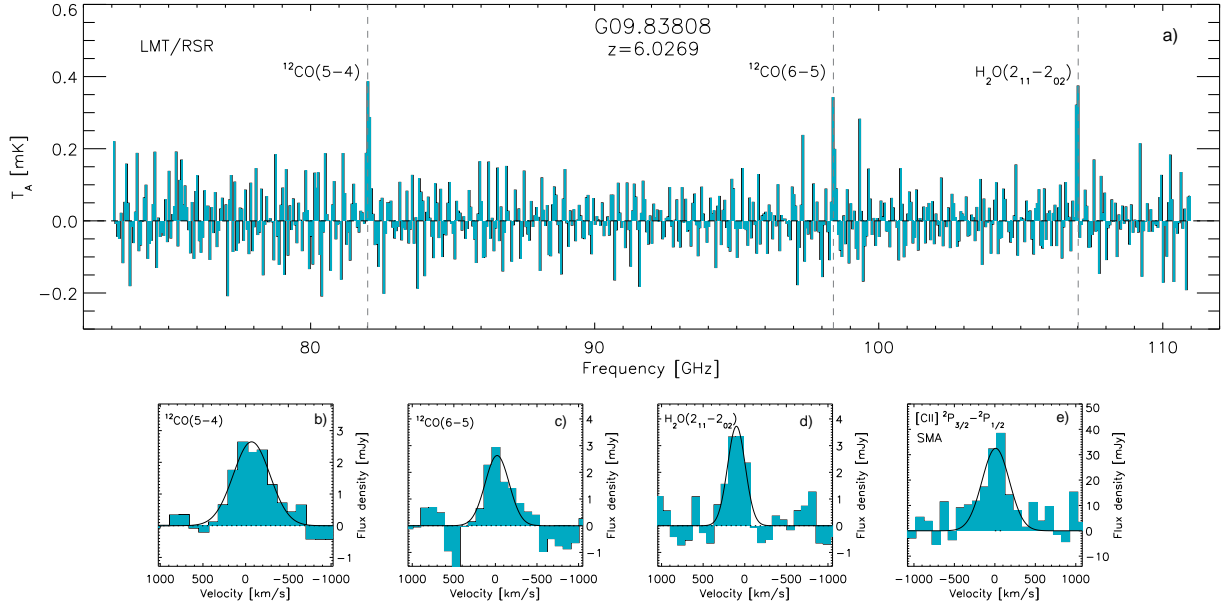


Figure 1: Identification of molecular emission lines and redshift derivation. **a)**, Wide-band Redshift Search Receiver (RSR) 3 mm spectrum of G09 83808 taken with the Large Millimeter Telescope (LMT). The transitions detected above $S/N = 5$ are marked with vertical dashed lines, and correspond to $^{12}\text{CO}(5-4)$, $^{12}\text{CO}(6-5)$, and $\text{H}_2\text{O}(2_{11}-2_{02})$ at $z = 6.0269 \pm 0.0006$. The spectrum has been rebinned into 2 pixels bins ($\sim 200\text{km/s}$) for better visualization. **b)**, **c)**, **d)**, LMT/RSR raw spectra at the position of the detected lines along with the best-fitting Gaussian profiles. **e)**, SMA spectrum centered at the position of the detected line. The x -axis is in velocity offset with respect to the derived redshift of $z = 6.0269$. The derived properties of the lines are reported in Table 1.

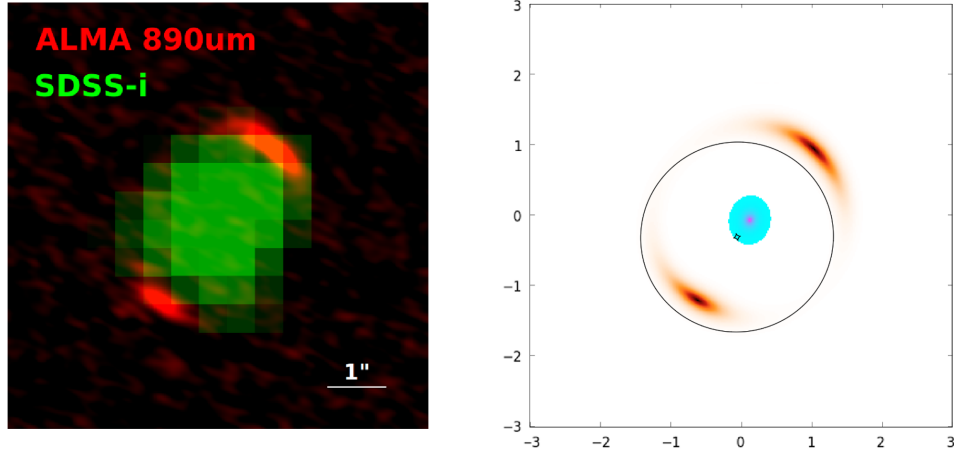


Figure 2: ALMA high-angular resolution continuum observations and lensing model. *Left:* Color composite image of G09 83808. The green channel represents the *i*-band data from SDSS and the red channel the ALMA 890 μm observations. An Einstein ring-like structure of radius ≈ 1.4 arcsec in the ALMA image is clearly seen around a foreground galaxy at $z = 0.776$, which confirms that our high-redshift galaxy is strongly amplified. *Right:* Best-fit lensing model based on the visibilities of ALMA observations, from which we derived a gravitational amplification of $\mu = 9.3 \pm 1.0$.

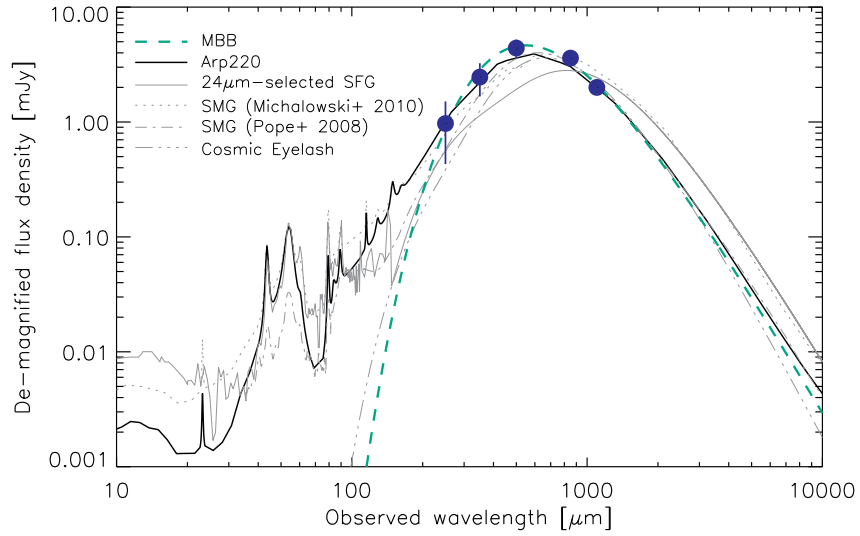


Figure 3: Photometry and spectral energy distribution (SED). De-magnified (with $\mu = 9.3 \pm 1.0$) flux densities at 250, 350, 500, 850 and 1100 μm from *Herschel/SPIRE*, JCMT/SCUBA-2, and AzTEC/LMT are represented by the blue circles. These flux densities were fitted with different SED templates, including: Arp220, Cosmic Eyelash, two average SMG templates, an average 24 μm -selected star-forming galaxy template, and a modified black body (MBB, see Methods section for details). We achieve the lowest χ^2 with the Arp220 template, from which we derive an IR luminosity of $3.8 \pm 0.5 \times 10^{12}$ (corrected for magnification). From the best-fit modified black body distribution we derive a dust temperature of 49 ± 3 K. As discussed in the Methods section, the CMB effects are not significant.

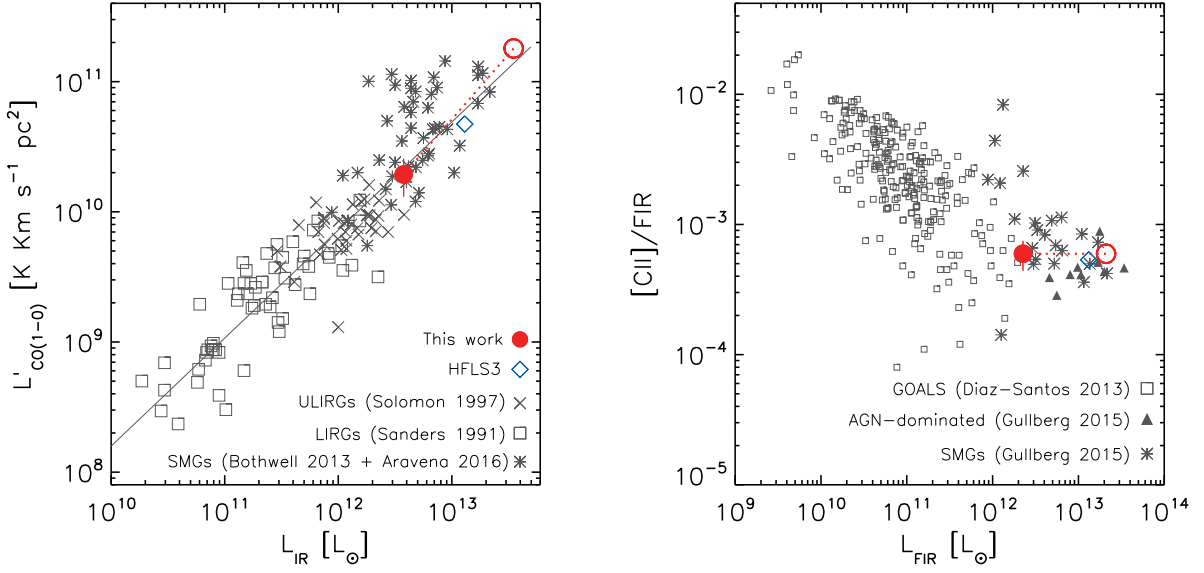


Figure 4: Star formation efficiency and [CII] deficiency. *Left:* Lens-corrected CO(1-0) luminosity versus IR luminosity ($L'_{\text{CO}(1-0)} - L_{\text{IR}}$) as a proxy for the star-formation efficiency of G09 83808. For comparison, local LIRGs³¹, ULIRGs³², and lower-redshift SMGs^{22,33} are plotted along with the best-fit relation to the three samples²². As can be seen, G09 83808 falls on the same relation (as well as HFLS3¹² after correcting for magnification¹⁵), which suggests that the same star formation efficiency holds from $z \sim 0$ to $z = 6$ (i.e. during the last ~ 12.8 Gyr). The empty circle represents the position of our source if no lensing amplification correction is applied. *Right:* [CII]/FIR versus de-magnified (filled circle) and amplified (empty circle) FIR luminosity for G09 83808. For comparison, we also plot a sample of (U)LIRG galaxies from the Great Observatories All-sky Survey (GOALS²⁸), and a compilation of high-redshift sources⁴⁵ that includes SMGs and AGN-dominated sources. Our source follows the same trend as local (U)LIRGs and lies in a region between lower-redshift SMGs and AGNs.

Table 1: Measured spectral line and continuum properties (not corrected for gravitational amplification).

| | Transition | | | [CII] | Photometry ^a | |
|--|--------------------|------------------|--|-------------------|-------------------------|----------------|
| | CO(5-4) | CO(6-5) | H ₂ O(2 ₁₁ - 2 ₀₂) | | [μm] | [mJy] |
| ν_{obs} [GHz] | 82.031 ± 0.007 | 98.41 ± 0.01 | 106.993 ± 0.007 | 270.35 ± 0.03 | 250 | 9.7 ± 5.4 |
| FWHM [km s ⁻¹] | 490 ± 60 | 320 ± 70 | 240 ± 40 | 400 ± 70 | 350 | 24.6 ± 7.9 |
| S_{int} [Jy km s ⁻¹] | 1.6 ± 0.3 | 0.9 ± 0.3 | 0.8 ± 0.2 | 13.8 ± 3.0 | 500 | 44.0 ± 8.2 |
| L' [10^{10} K km s ⁻¹ pc ⁻²] | 7.6 ± 1.2 | 2.9 ± 0.8 | 2.3 ± 0.5 | 6.1 ± 1.3 | 850 | 36.0 ± 3.1 |
| | | | | | 1100 | 20.0 ± 1.0 |

^aThe flux densities at 250, 350, 500, and 850 μm were taken from ref.¹⁴

Methods

1 Observations and data reduction

1.1 LMT observations

Continuum and spectroscopic observations were obtained using the Large Millimeter Telescope (LMT³⁴, PI: D. Hughes), located on the summit of Volcán Sierra Negra (*Tliltépetl*), Mexico, at ~ 4600 m.a.s.l. Observations were carried out during the Early Science Phase of the telescope using the 1.1 mm continuum camera, AzTEC³⁵, and the 3 mm spectrograph, Redshift Search Receiver (RSR³⁶). During these observations only the inner 32-m diameter region of the telescope active surface was illuminated, which provided an effective beam size of ≈ 8.5 arcsec at 1.1 mm and between 20 – 28 arcsec in the RSR 3 mm window (75 GHz - 110 GHz).

AzTEC observations were performed on 2014 November 10 with an opacity of $\tau_{225} = 0.07$ and total on-source integration time of 11 min. Data reduction were done following the AzTEC Standard Pipeline³⁷. G09 83808 was detected with a S/N ≈ 20 with a flux density of $S_{1.1\text{mm}} = 20.0 \pm 1.0$ mJy. RSR observations were subsequently taken at the AzTEC position in two different periods: February 2016 and February 2017, along five different nights with an opacity range of $\tau_{225} = 0.05 - 0.15$ and a total integration time of 8 hrs. Pointing observations on bright millimetre sources were done every hour. Data reduction was performed using the Data Reduction and Analysis Methods in Python (DREAMPY). The final spectra were obtained by averaging all scans using $1/\sigma^2$ weights after flagging bad data. Finally, to convert from antenna temperature units to flux, a factor of 7 Jy K^{-1} was used³⁸. The final spectrum shows three lines detected at S/N $\gtrsim 5$ associated to CO(6-5), CO(5-4) and H₂O(2₁₁ - 2₀₂) at $z = 6.0269$. A cross-correlation template analysis³⁸ also identifies this redshift as the best solution with a S/N = 9.1. Figure 1 shows the final spectrum after a Savitzky-Golay filter³⁹ has been applied for better visualization (the filter does not modified any of the properties of the detected lines).

At the redshift of our source the [CII] 158 μm line

(see below) falls within the AzTEC band pass and then contributes to the total flux density measured at 1.1 mm. However, the contamination from the line is measured to be less than 2 per cent. Even if the [CII] line luminosity was as high as 1 per cent of the total IR luminosity, the contamination to the AzTEC measurement would be only ~ 6 per cent, which is similar to the absolute flux calibration uncertainty. Therefore, and at least for this source, the contamination of the emission line to the 1.1 mm continuum flux density is less important than anticipated⁴⁰.

1.2 SMA observations

G09 83808 was observed with the Submillimeter Array (SMA, PI: J. Zavala) on Mauna Kea, Hawaii, on 2017 April 03. The weather conditions were good, with an average atmospheric opacity of $\tau_{225} = 0.07$ and stable phase. The seven available array antennas were in a compact configuration that provided baseline lengths from 8 to 77 meters. The ‘345’ receiver set was tuned to provide spectral coverage $\pm(4 - 12)$ GHz from a LO frequency of 277.5 GHz, specifically to span a broad range around the estimated (redshifted) [CII] line frequency of ~ 270.5 GHz in the lower sideband. The SWARM correlator provided uniform channel spacing of 140 kHz ($\sim 0.16 \text{ km s}^{-1}$) over the full bandwidth. The usable field of view is set by the FWHM primary beam size of ~ 47 arcsec at this frequency.

The basic observing sequence consisted of a loop of 2 minutes each on the gain calibrators J0909+013 (1.57 Jy) and J0831+044 (0.47 Jy) and 17.5 minutes on G09 83808. The track spanned an hour angle range of -0.8 to 4.8 for the target source. Pass-band calibration was obtained with observations of the strong quasar 3C279. The absolute flux scale was set using observations of Callisto, with an estimated accuracy of 20%. All of the basic calibration steps were performed using standard procedures in the MIR software package. The calibrated visibilities were exported to the MIRIAD software package for imaging and deconvolution. Within MIRIAD, the task `uvaver` was used to combine the 4 correlator windows of the lower sideband and to resam-

ple the visibilities to 50 km s^{-1} spectral resolution. The task `uvlin` was used for continuum subtraction, using a linear fit to line-free channel ranges in the band. The task `invert` provided Fourier inversion for both continuum and spectral line imaging, followed by `clean` for deconvolution. The synthesized beam size obtained with natural weighting was $2.5'' \times 2.3''$, p.a. 82° for the spectral line data cube, with rms noise 7.1 mJy per 50 km s^{-1} bin. The final spectrum (see Fig. 1) was then extracted from a rectangular region that comprise all the continuum emission. We measured the continuum flux density of the source to be $21.5 \pm 3 \text{ mJy}$, in very good agreement with the AzTEC photometry.

2 Lensing model

The lens model was created using the publicly-available `visilens` code⁴¹; details of the code are given in that work. Briefly, the lens mass profile is parameterized as a Singular Isothermal Ellipsoid, and the background source is modeled with a single elliptical Sérsic profile. The parameter space is explored using a Markov Chain Monte Carlo sampling method, generating a model lensed image at each proposed combination of lens and source parameters. The redshift of both sources is fixed at $z = 0.776$ (based on X-Shooter/VLT observations⁴²) and $z = 6.027$, respectively. Because pixel values in interferometric images are correlated and subject to difficult-to-model residual calibration errors, the proposed model image is inverted to the visibility domain and sampled at the uv coordinates of the ALMA data (Project code: 2013.1.00001.S; PI: R. Ivison; Oteo et al. in preparation). We also allow for residual antenna-based phase calibration errors in the model which could be due to, for example, uncompensated atmospheric delays. The phase shifts of all antennas are $< 10 \text{ deg}$, indicating that no significant antenna-based calibration problems remain.

The lensed emission is reasonably well-fit by a single background Sérsic component, leaving peak residuals of $\sim 4\sigma$ (the source is detected at peak significance $\sim 20\sigma$). These residuals may indicate that either the lens, source, or both are more complex than the simple parametric forms we have assumed. We

have verified that an additional background source component is not statistically motivated. The best-fit magnification of the source is $\mu_{890\mu\text{m}} = 9.3 \pm 1.0$, with an intrinsic flux density $S_{890\mu\text{m}} = 4.3 \pm 0.5 \text{ mJy}$ and half-light radius $0.10 \pm 0.01''$ ($= 0.6 \pm 0.1 \text{ kpc}$). This compact morphology resembles the sizes found for local ULIRGs⁴³ ($\sim 0.5 \text{ kpc}$), which are smaller than the typical values in SMGs ($\sim 1.8 \text{ kpc}$, ref:⁴⁴).

3 SED fitting and dust properties

We fit different galaxy SED templates to the photometry of G09 83808 through a χ^2 minimization method. We include the SED template of Arp220⁴⁶, Cosmic Eyelash⁴⁷ (SMM J2135-0102), two average SMGs templates^{48,49}, and finally a composite SED of $24 \mu\text{m}$ -selected star-forming galaxy⁵⁰. All the SED templates were fixed at $z = 6.027$. The Arp220 SED template gives us the best fit with $\chi_{\text{red}}^2 = 0.7$. Using this template we derive an IR ($8 - 1000\mu\text{m}$) luminosity of $3.8 \pm 0.5 \times 10^{12} L_\odot$ and a FIR ($42.5 - 122.5\mu\text{m}$) luminosity of $2.3 \pm 0.3 \times 10^{12} L_\odot$ (both corrected for gravitational amplification). For comparison, if we adopt instead a SMGs average template ($\chi_{\text{red}}^2 = 1.2$) we obtain $L_{\text{IR}} = 3.0 \pm 0.4 \times 10^{12} L_\odot$, which is in good agreement with the value derived using the Arp220 template. Using Kennicutt standard relation⁵¹ for a Chabrier initial mass function (IMF)⁵³, this IR luminosity corresponds to a star formation rate (SFR) of $380 \pm 50 M_\odot \text{ yr}^{-1}$, or to $570 \pm 70 M_\odot \text{ yr}^{-1}$ if the most recent relation⁵² is used. If we adopt instead the Kennicutt calibration⁵¹ for a Salpeter IMF⁵⁴, the SFR increases to $640 \pm 90 M_\odot \text{ yr}^{-1}$, still below the range probed by other SMGs at $z \gtrsim 5$.

We also use a modified blackbody function to fit our photometric measurements described by

$$S_\nu \propto \{1 - \exp[-(\nu/\nu_0)^\beta]\} B(\nu, T_d), \quad (1)$$

where S_ν is the flux density at frequency ν , ν_0 is the rest-frame frequency at which the emission becomes optically thick, T_d is the dust temperature, β is the emissivity index, and $B(\nu, T_d)$ is the Planck function at temperature T_d . To minimize the number of free parameters, the emissivity index is fixed (previous

observational works suggest $\beta = 1.5 - 2$; refs.⁵⁵⁻⁵⁷), as well as $\nu_0 = c/100 \mu\text{m}$ (refs.^{12,58}), where c is the speed of light. From the best fit ($\chi^2 \approx 1.1$) we derive $T_d = 49 \pm 3 \text{ K}$ for $\beta = 1.8$ and $T_d = 52 \pm 3 \text{ K}$ for $\beta = 1.5$. For these dust temperatures and at the redshift of our source the CMB effects⁵⁹ are not significant ($\Delta T \lesssim 1 \text{ K}$).

Assuming the dust is isothermal, the dust mass, M_d , is estimated from

$$M_d = \frac{S_\nu/(1+z)D_L^2}{(1+z)\kappa_\nu B(\nu, T_d)}, \quad (2)$$

where S_ν is the flux density at frequency ν , κ_ν is the dust mass absorption coefficient at ν , T_d is the dust temperature, and $B(\nu, T_d)$ is the Planck function at temperature T_d . The dust mass absorption follows the same power law as the optical depth, $\kappa \propto \nu^\beta$. Assuming normalization of $\kappa_d(850\mu\text{m}) = 0.07 \text{ m}^2 \text{ kg}^{-1}$ (ref.⁶⁰) and a dust temperature of $49 \pm 3 \text{ K}$, we estimate a dust mass of $M_d = 1.9 \pm 0.4 \times 10^8 M_\odot$ after correcting for the CMB effects⁵⁹ (although this correction is less than 5 per cent). These calculations do not include the uncertainties of the dust mass absorption coefficient, which could be at least a factor of 3 (ref.⁶¹). If we use instead a lower dust temperature of 35 K, the dust mass increases by a factor of ~ 2 .

We also fit the observed photometry with the MAGPHYS⁶² SED modelling code finding consistent results, within the error bars, with median values of $\text{SFR} = 360_{-70}^{+80} M_\odot \text{ yr}^{-1}$, $L_{\text{IR}} = 4.5 \pm 0.7 \times 10^{12} L_\odot$, $T_d = 40_{-2}^{+4} \text{ K}$, and $M_d = 4.2 \pm 0.7 \times 10^8 M_\odot$.

4 Spectral line properties

We calculate the line luminosity for each detected line following the standard relation⁶³ described by:

$$L'_{\text{CO}} = 3.25 \times 10^7 S_{\text{CO}} \Delta V \nu_{\text{obs}}^{-2} D_L^2 (1+z)^{-3}, \quad (3)$$

where L'_{CO} is the line luminosity in $\text{K km s}^{-1} \text{ pc}^2$, $S_{\text{CO}} \Delta V$ is the velocity-integrated line flux in Jy km s^{-1} , ν_{obs} is the observed central frequency of the line in GHz and D_L is the luminosity distance in Mpc . The integrated flux, $S_{\text{CO}} \Delta V$, is calculated as the integral of the best-fit Gaussian distri-

bution, and its associated uncertainty through Monte Carlo simulations taking into account the errors in the Gaussian parameters (i.e. peak flux density and line width). To estimate the line luminosity in L_\odot , we use $L = 3 \times 10^{-11} \nu_r^3 L'$, where ν_r is the rest frequency of the line⁶³. All properties are summarized in Table 1.

5 CO(1-0) line luminosity and molecular gas mass

The molecular gas mass, $M(\text{H}_2)$, can be derived using the CO luminosity to molecular gas mass conversion factor, α , following the relation

$$M(\text{H}_2) = \alpha L'_{\text{CO}(1-0)}. \quad (4)$$

For the $L'_{\text{CO}(1-0)}$ line luminosity we adopt the average value of $L'_{\text{CO}(1-0)} = 2.0 \pm 0.8 \times 10^{10} \text{ K km s}^{-1} \text{ pc}^{-2}$ extrapolated from our CO(6-5) and CO(5-4) transitions and correcting for gravitational amplification. The extrapolation was done using average brightness ratios found for lower-redshift SMGs²² ($L'_{\text{CO}(5-4)}/L'_{\text{CO}(1-0)} = 0.32 \pm 0.05$, $L'_{\text{CO}(6-5)}/L'_{\text{CO}(1-0)} = 0.21 \pm 0.04$), this sample includes galaxies with similar luminosities to G09 83808 and are in agreement with those found for local ULIRGs²⁴ (within the large scatter). On the other hand, if we use the relationship between the Rayleigh-Jeans specific luminosity and CO(1-0) luminosity⁶⁴, $L'_{\text{CO}(1-0)} [\text{K km s}^{-1} \text{ pc}^2] = 3.02 \times 10^{-21} L_\nu [\text{erg s}^{-1} \text{ Hz}^{-1}]$, we obtain a consistent line luminosity of $1 \pm 0.1 \times 10^{10} \text{ K km s}^{-1} \text{ pc}^{-2}$ (assuming a mass-weighted dust temperature of 35 K, which is different from the luminosity-weighted dust temperature determined from the SED fitting⁶⁴). Using the former value and $\alpha = 0.8 M_\odot (\text{K km s}^{-1} \text{ pc}^{-2})^{-1}$, which is appropriate for starburst galaxies²³ (although some studies suggest larger values⁶⁵), we derive a molecular mass of $M(\text{H}_2) = 1.6 \pm 0.6 \times 10^{10} M_\odot$.

6 Dynamical mass and gas mass fraction

Dynamical mass has been derived using the ‘isotropic virial estimator’, which has been shown to

be appropriate for lower-redshift SMGs⁶⁶:

$$M_{\text{dyn}}[M_{\odot}] = 2.8 \times 10^5 \Delta\nu_{\text{FWHM}}^2 [\text{km s}^{-1}] R_{1/2} [\text{kpc}], \quad (5)$$

where $\Delta\nu_{\text{FWHM}}$ is the integrated line FWHM, which has been assumed to be 400 km/s (as the average between the CO and [CII] lines), and $R_{1/2}$ is the half-light radius of ~ 0.6 kpc (derived from the lesing model of the continuum emission). This results in a dynamical mass of $M_{\text{dyn}} = 2.6 \times 10^{10} M_{\odot}$. Using this estimation we calculate a gas mass fraction of $f_{\text{gas}} = M_{\text{H}_2}/M_{\text{dyn}} \approx 60\%$. This constrain the CO luminosity to molecular gas mass conversion factor to $\alpha \lesssim 1.4 M_{\odot} (\text{K km s}^{-1} \text{pc}^{-2})^{-1}$, otherwise the molecular gas mass would exceed the dynamical mass.

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