

LETTER TO THE EDITOR

## HerMES: SPIRE galaxy number counts at 250, 350, and 500 $\mu\text{m}$ <sup>\*</sup>

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### ABSTRACT

Emission at far-infrared wavelengths makes up a significant fraction of the total light detected from galaxies over the age of Universe. *Herschel* provides an opportunity for studying galaxies at the peak wavelength of their emission. Our aim is to provide a benchmark for models of galaxy population evolution and to test pre-existing models of galaxies. With the *Herschel* Multi-tiered Extra-galactic survey, HerMES, we have observed a number of fields of different areas and sensitivity using the SPIRE instrument on *Herschel*. We have determined the number counts of galaxies down to  $\sim 20$  mJy. Our constraints from directly counting galaxies are consistent with, though more precise than, estimates from the BLAST fluctuation analysis. We have found a steep rise in the Euclidean normalised counts  $< 100$  mJy. We have directly resolved  $\sim 15\%$  of the infrared extra-galactic background at the wavelength near where it peaks.

**Key words.** galaxies: evolution – submillimeter: galaxies – surveys

## 1. Introduction

The statistical properties of galaxy populations are important probes for understanding the evolution of galaxies. The most basic statistic of galaxy populations is the number counts i.e. the number density of galaxies as a function of flux. The first strong evidence for cosmological evolution came through studying number counts of radio galaxies (e.g. Longair 1966).

The number counts at far-infrared and sub-mm wavelengths are well known to exhibit strong evolution, e.g. from IRAS (Oliver et al. 1992, and references therein), ISO (Oliver et al. 2002; Héraudeau et al. 2004, and references therein), *Spitzer* (Shupe et al. 2008; Frayer et al. 2009, and references therein) and ground-based sub-mm surveys (Maloney et al. 2005; Coppin et al. 2006; Khan et al. 2007; Greve et al. 2008; Weiß et al. 2009; Scott et al. 2010, and references therein).

These results are underlined by the discovery of a significant extragalactic infrared background (Puget et al. 1996; Fixsen et al. 1998; Lagache et al. 1999). The background measures the flux weighted integral of the number counts over all redshifts

plus any diffuse cosmological component. This indicates that as much energy is received from galaxies after being reprocessed through dust as is received directly. It is only very recently, using BLAST, that count models have been probed using fluctuation techniques (Patanchon et al. 2009) or directly (Béthermin et al. 2010) at the wavelength where the background peaks.

Far-infrared and sub-mm counts and background measurements have been modelled phenomenologically with strongly evolving populations (see Sect. 4). Physical models (e.g. so-called semi-analytic models) struggle to explain these counts and solutions include altering the initial mass function (e.g. Baugh et al. 2005) or exploiting AGN/Supernovae feedback (e.g. Granato et al. 2004).

A primary goal of *Herschel* (Pilbratt et al. 2010) is to explore the evolution of obscured galaxies. *Herschel* opens up a huge region of new parameter space of surveys in area, depth and wavelength.

The *Herschel* Multi-tiered Extragalactic Survey (HerMES<sup>1</sup> Oliver et al. 2010, in prep.) is the largest project being undertaken by *Herschel* and consists of a survey of many well-studied extra-galactic fields (totalling  $\sim 70$  deg<sup>2</sup>) at various depths. This

<sup>\*</sup> *Herschel* is an ESA space observatory with science instruments provided by European-led Principal Investigator consortia and with important participation from NASA.

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**Table 1.** HerMES SPIRE SDP observations.

Name	Size	RA /°	Dec /°	Roll /°	Mode	Scan rate	Repeats	$t_{\text{AOR}}$ /hr	$\langle N_{\text{samp}} \rangle$	$S_{50\%}^{250\ \mu\text{m}}$ /mJy	$S_{50\%}^{350\ \mu\text{m}}$ /mJy	$S_{50\%}^{500\ \mu\text{m}}$ /mJy
A2218	9' × 9'	248.98	66.22	217	Lrg. Map	30"/s	100	9.2	1622	13.8	16.0	15.1
FLS	155' × 135'	258.97	59.39	185	Parallel	20"/s	1	16.8	30	17.5	18.9	21.4
Lockman-North	35' × 35'	161.50	59.02	91	Lrg. Map	30"/s	7	3.9	117	13.7	16.5	16.0
Lockman-SWIRE	218' × 218'	162.00	58.11	92	Lrg. Map	60"/s	2	13.4	16	25.7	27.5	33.4
GOODS-N	30' × 30'	189.23	62.24	132	Lrg. Map	30"/s	30	13.5	501	12.0	13.7	12.8

**Notes.** PACS observations will be discussed in Aussel et al. (in prep.). Size is approximate extent of region with typical coverage. Roll angle is measured East of North. Repeats is total number of pairs of scans in both *A* and *B* directions.  $t_{\text{AOR}}$  is total time in execution of the observations.  $\langle N_{\text{samp}} \rangle$  is the mean number of bolometer samples per pixel in the same typical-coverage region of the 250  $\mu\text{m}$  map (6" × 6" pixels).  $S_{50\%}^{250\ \mu\text{m}}$  is the flux density at which 50% of sources injected into the 250  $\mu\text{m}$  map are faithfully recovered.

letter is the first number count analysis from the HerMES science demonstration phase (SDP) SPIRE data. Even these preliminary results will be able to eliminate some existing models and provide a benchmark on which future models can be tested.

## 2. SPIRE data

### 2.1. SDP observations

The observations described here were carried out on the *Herschel* Space Observatory (Pilbratt et al. 2010) using the Spectral and Photometric Imaging REceiver (SPIRE). The SPIRE instrument, its in-orbit performance, and its scientific capabilities are described by Griffin et al. (2010), and the SPIRE astronomical calibration methods and accuracy are outlined in Swinyard et al. (2010). They were undertaken as part of the HerMES programme during the SDP between 12-Sep.-2009 and 25-Oct-2009 under the proposal identification SDP\_soliver\_3. The fields and observations are summarised in Table 1.

### 2.2. SPIRE catalogue data processing

For this paper it is sufficient to note that the same source detection method is applied to the simulations as to the real data so we sketch the details only briefly. The single-band SPIRE catalogues have been extracted from the maps using a version of the SUSSEXtractor method (Savage & Oliver 2007) as implemented in HIPE (Ott 2010). The processing of the SPIRE data is summarised here, with details of the approach given by Smith et al. (in prep.) Calibrated timelines were created using HIPE development version 2.0.905, with a fix applied to the astrometry (included in more recent versions of the pipeline), with newer calibration files (beam-steering mirror calibration version 2, flux conversion version 2.3 and temperature drift correction version 2.3.2) and with a median and linear slope subtracted from each timeline. The default HIPE naïve map-maker was then used to create maps, which were given a zero mean. The shallow fields (Lockman-SWIRE and FLS) were smoothed with a point-source optimised filter (see Smith et al. in prep. for details). Peaks in the map were identified and the flux was estimated based on an assumed (Gaussian) profile for a point source, through a weighted sum of the map pixels close to the centre of the source. This filtering means we underestimate the flux of extended sources. The SPIRE CATalogue (SCAT) processing is assessed by injecting synthetic sources on a grid into the real maps. We then run the SCAT source extraction pipeline on these maps and claim success if the closest detection to the injected source is within a search radius of FWHM of the beam and has a flux within a factor of two of the injected flux (see Smith et al. in prep. for more details). The resulting 50% completeness estimated in this way is tabulated in Table 1 but is not used to assess the counts.

## 3. Method

The *Herschel* beam is broad compared with the number density of sources, i.e. the maps are confused. Nguyen et al. (2010) measure a variance in nominal map pixels due to confused sources finding  $\sigma_{\text{conf}}$  5.8, 6.3 and 6.8 mJy/beam. This confusion means care has to be taken in the estimation of number counts. Our technique follows the standard approach for sub-mm surveys, correcting for flux boosting and incompleteness.

We determined the false detection rate by applying the source extraction on maps obtained from the difference between two independent observations of the same field. These maps are expected to have zero mean, no sources, but similar noise properties to mean maps. We thus estimate that the reliability for the samples in this paper is better than 97% for all fields and bands.

A source we measure to have flux,  $S_m$ , and noise,  $\sigma_m$ , is more likely to be a dimmer source on top of a positive noise fluctuation than the converse; this is known as flux boosting. We follow the Bayesian method of Crawford et al. (2009) for estimating fluxes of individual sources (“de-boosting”). We estimate the posterior probability distribution of the true flux of the source ( $S_i$ ) that contributed the most flux to a given detection. Note that it is similar to the now-standard flux de-boosting method (e.g. Coppin et al. 2006) but with an additional exponential suppression term at low intrinsic flux. We derive counts by randomly sampling the posterior distribution ten thousand times. The flux de-boosting procedure has some dependency on the choice of prior number count model and so these samples are drawn from distributions produced for the full range of models discussed in Sect. 4. These samples provide a direct estimate of the confidence region of our counts. For faint sources, the posterior probability function rises beyond the sampled flux range ( $S_m/5 < S_i < 5 S_m$ ) so that the deboosted flux is highly uncertain. In those cases we flag the deboosted flux as “bad” and the derived number counts at the flux level where deboosted fluxes are flagged “bad” are unreliable. We exclude count bins in which the fraction of “bad” sources is  $> 20\%$ . We also estimate the uncertainty in this by looking at the variation in derived counts from the range of models. Errors are included in the plots. The flux de-boosting procedure assumes no clustering. Clustering will affect this and will be addressed in a later paper.

We estimate the incompleteness in the whole process by running full simulations. We have constructed input maps various from various number count models (Pearson & Khan 2009; Lagache et al. 2004; Patanchon et al. 2009) and models from Xu et al. 2003; Lacey et al. 2010). These input maps are then processed by the SPIRE photometer simulator (SPS, Sibthorpe et al. 2009) for observational programmes exactly the same as the real data. The timeline output of the SPS, map-making and source extraction are then processed in the same way as the real data, including the flux-deboosting. The ratio of input to output

**Table 2.** Table of HerMES SDP number counts at 250, 350 and 500  $\mu\text{m}$ .

	Bin 1	Bin 2	Bin 3	Bin 4	Bin 5	Bin 6
$S_{\min}$	20	29	51	69	111	289
$S_{\max}$	29	51	69	111	289	511
$S_{\text{euc}}$	23.8	37.5	58.9	85.9	166.2	374.1
250 $\mu\text{m}$						
$\frac{dN}{dS}$	$2.0 \times 10^8$	$6.4 \times 10^7$	$1.2 \times 10^7$	$3.1 \times 10^6$	$2.1 \times 10^5$	$1.7 \times 10^4$
err <sub>1</sub>	38	22	28	35	23	6
err <sub>2</sub>	10	6	7	6	4	19
err <sub>3</sub>	4	3	7	4	7	23
err <sub>4</sub>	11	19	30	28	30	8
350 $\mu\text{m}$						
$\frac{dN}{dS}$	$1.1 \times 10^8$	$3.5 \times 10^7$	$5.3 \times 10^6$	$1.1 \times 10^6$	$6.2 \times 10^4$	$4.7 \times 10^3$
err <sub>1</sub>	49	34	44	56	25	129
err <sub>2</sub>	18	14	17	6	5	12
err <sub>3</sub>	7	4	10	7	15	43
err <sub>4</sub>	23	13	33	8	12	64
500 $\mu\text{m}$						
$\frac{dN}{dS}$	$3.6 \times 10^7$	$1.1 \times 10^7$	$1.6 \times 10^6$	$2.3 \times 10^5$	$1.3 \times 10^4$	$1.3 \times 10^3$
err <sub>1</sub>	83	50	62	56	45	0
err <sub>2</sub>	31	18	25	15	18	7
err <sub>3</sub>	10	6	18	14	33	50
err <sub>4</sub>	5	17	48	27	20	0

**Notes.** Bin limits and Euclidian weighted central fluxes are given in mJy. Counts are in  $/\text{sr}^{-1}\text{Jy}^{-1}$ . Errors 1–4 are fractional errors in percentages arising from: flux-deboosting, completeness corrections, Poisson statistics and field-field variations respectively.

counts gives us the completeness with the standard deviation between models providing an estimate of an error in that estimate.

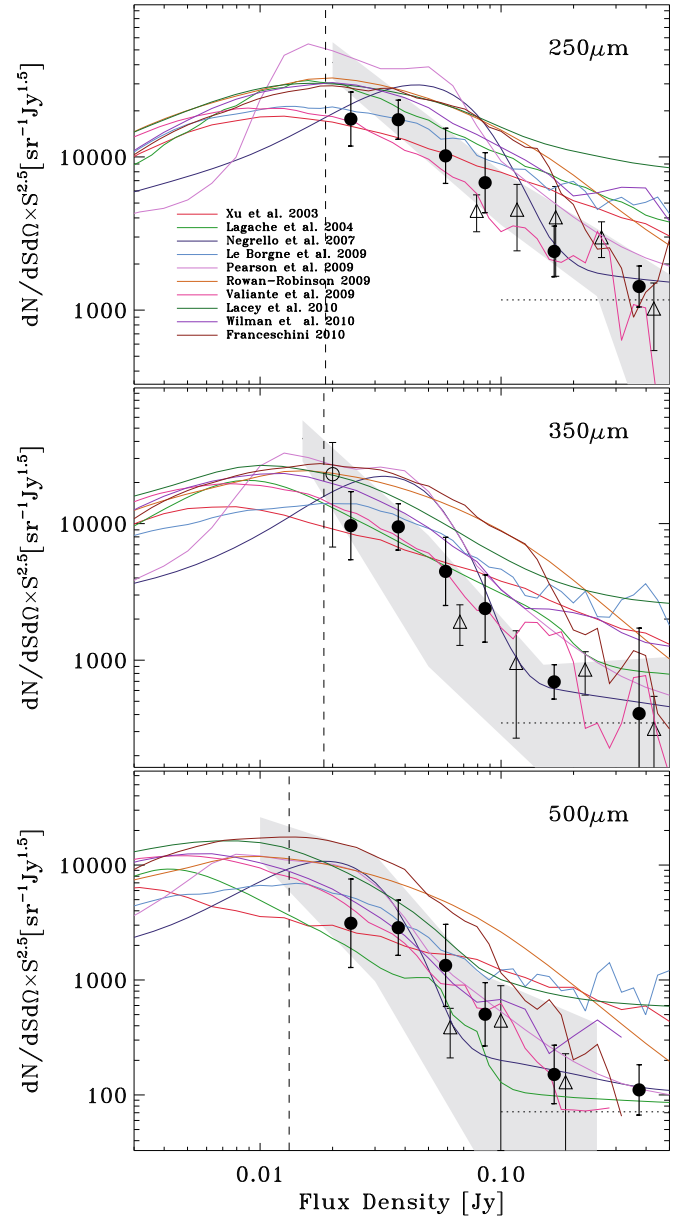
#### 4. Results

The results are presented in Table 2 and in Fig. 1 with Euclidian normalization. There are several sources of uncertainties in the number counts: Poisson noise from the raw counts; “sampling variance” due to additional fluctuations from real large-scale structure; additional Poisson noise from the sampling of the posterior flux distribution and systematic errors from the corrections and assumptions about priors and the effect of clustering on the de-boosting. We measure the standard deviation of counts between fields which includes Poisson errors and some of the other systematic errors. The errors plotted are the field to field variations (or the Poisson errors if larger) with the errors from flux-boosting and completeness corrections added in quadrature.

We see approximately flat counts for  $S > 100$  mJy and then a steep rise. There is flattening to about 20 mJy. We find very good agreement with the number counts estimated from a  $P(D)$  fluctuation analysis of the BLAST maps (Patanchon et al. 2009).

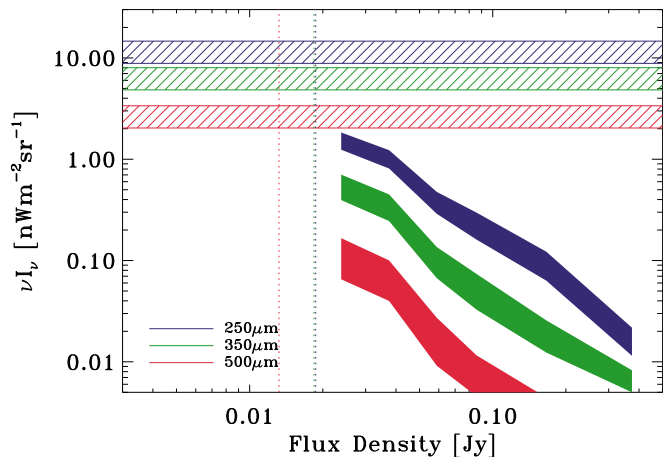
We have also estimated, but do not show, the integral counts. The flux density at which the integral source counts reach 1 source per 40 beams (with beams defined as  $3.87 \times 10^{-5}$ ,  $7.28 \times 10^{-5}$ ,  $1.48 \times 10^{-4}$   $\text{deg}^2$ ) is  $18.7 \pm 1.2$ ,  $18.4 \pm 1.1$  and  $13.2 \pm 1.0$  mJy at 250, 350 and 500  $\mu\text{m}$  respectively (N.B. these fluxes are slightly below our secure estimation of counts). Likewise the number density at 100 mJy is  $12.8 \pm 3.5$ ,  $3.7 \pm 0.4$  and  $0.8 \pm 0.1$   $\text{deg}^{-2}$ . These last measurements alone will be sufficient to rule out many models.

Since the first IRAS results, many empirical models have been developed to predict and interpret the numbers and luminosities of IR galaxies as a function of redshift. Empirical models are based on a similar philosophy. The spectral energy distributions of different galaxy populations are fixed and the mid-IR, far-IR and submm data are used to constrain the luminosity function evolution. Current limits come from the mid-IR,



**Fig. 1.** Number counts obtained from HerMES source catalogues. Filled circles are the mean number counts averaged over the following fields. GOODS-N & Lockman-North (faintest five bins only) and FLS & Lockman-SWIRE (brightest six bins only) with flux-deboosting, completeness corrections and field-field error bars. Model fit to fluctuations of BLAST maps (omitting upper-limits, Patanchon et al. 2009) – shaded region; BLAST resolved counts (B  thermin et al. 2010) – open triangles; Khan et al. (2007) data point – open circle; asymptote from modelling of IRAS data (Serjeant & Harrison 2005) – dotted line. Models are discussed in the text. Dashed line indicates the flux at which the integrated number density is  $(40 \text{ beams})^{-1}$ .

far-IR and submm number counts, redshift distributions, luminosity functions, and cosmic IR background. Models all agree on the general trends, with a very strong evolution of the bright-end ( $> 10^{11} L_{\odot}$ ) of the luminosity function and they yield approximately the same comoving number density of infrared luminous galaxies as a function of redshift. We compare the number counts with eight models, one pre-*Spitzer* (Xu et al. 2003), two based on the ISO, SCUBA and *Spitzer* first results (Lagache et al. 2004; Negrello et al. 2007) and 5 being more constrained by deep *Spitzer*, SCUBA, AzTEC, and recent BLAST observations



**Fig. 2.** The integrated background light at 250, 350, 500  $\mu\text{m}$  from the HerMES counts determined in Fig. 1. Dotted lines are the flux at which the integrated density is  $(40 \text{ beams})^{-1}$ . The hatched regions are measurements of the COBE background (Lagache et al. 1999).

(Le Borgne et al. 2009; Pearson & Khan 2009; Rowan-Robinson 2009; Valiante et al. 2009; Franceschini et al. 2009). The differences between the models are in several details, different assumptions leading sometimes to equally good fits to the current data. For example, Valiante et al. (2009) conclude that it is necessary to introduce both an evolution in the AGN contribution and an evolution in the luminosity-temperature relation, while Franceschini et al. (2009) reproduce the current data with only 4 galaxy populations and only one template for each population. We also compare with two semi-analytic models those of Lacey et al. (2010) and Wilman et al. (2010).

Comparison with SPIRE number counts shows that many models cannot fit the bright end ( $>100 \text{ mJy}$ ). Exceptions are the models of Negrello et al. (2007), Valiante et al. (2009), Franceschini et al. (2009) and Pearson & Khan (2009). Of these only Valiante et al. (2009) can fit the rise from  $(20 < S < 100) \text{ mJy}$ . The Valiante et al. (2009) model has “cooler” spectral energy distributions at higher redshift. However, increasing the number of higher redshift galaxies would have a similar effect on the counts so it would be premature to assume the spectral energy distributions need revision.

We have also calculated the contribution of the resolved sources to the background intensity as a function of flux (shown in Fig. 2). At the  $(40 \text{ beams})^{-1}$  depth we resolve  $1.73 \pm 0.33$ ,  $0.63 \pm 0.18$ ,  $0.15 \pm 0.07 \text{ nW m}^2$  or 15, 10, 6% of the nominal measured values at 250, 350 and 500  $\mu\text{m}$  (Lagache et al. 1999).

Future work will provide more detailed constraints at the fainter limits. This will include a  $P(D)$  analysis (Glenn et al. in prep.) and counts from catalogues extracted at known *Spitzer* source positions.

## 5. Conclusions

We present the first SPIRE number count analysis of resolved sources, conservatively within the limit of *Herschel* confusion. We have measured counts which resolve around 15% of the infrared background at 250  $\mu\text{m}$ . We see a very steep rise in the counts from 100 to 20 mJy in 250, 350 and 500  $\mu\text{m}$ . Few models

have quite such a steep rise. Many models fail at the bright counts  $>100 \text{ mJy}$ . This may suggest that models need a wider variety or evolution of the spectral energy distributions or changes in the redshift distributions. Future work is required to accurately constrain the fainter ends  $<20 \text{ mJy}$  where confusion is a serious challenge.

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