

# The mass, location, and heating of the dust in the Cassiopeia A supernova remnant

F. D. Priestley,<sup>1</sup>★ M. J. Barlow<sup>1</sup> and I. De Looze<sup>1,2</sup>

<sup>1</sup>*Department of Physics and Astronomy, University College London, Gower Street, London WC1E 6BT, UK*

<sup>2</sup>*Sterrenkundig Observatorium, Ghent University, Krijgslaan 281 – S9, 9000 Gent, Belgium*

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## ABSTRACT

We model the thermal dust emission from dust grains heated by synchrotron radiation and by particle collisions, under conditions appropriate for four different shocked and unshocked gas components of the Cassiopeia A (Cas A) supernova remnant (SNR). By fitting the resulting spectral energy distributions (SEDs) to the observed SNR dust fluxes, we determine the required mass of dust in each component. We find the observed SED can be reproduced by  $\sim 0.6 M_{\odot}$  of silicate grains, the majority of which is in the unshocked ejecta and heated by the synchrotron radiation field. Warmer dust, located in the X-ray emitting reverse shock and blastwave regions, contribute to the shorter wavelength infrared emission but make only a small fraction of the total dust mass. Carbon grains can at most make up  $\sim 25$  per cent of the total dust mass. Combined with estimates for the gas masses, we obtain dust-to-gas mass ratios for each component, which suggest that the condensation efficiency in the ejecta is high, and that dust in the shocked ejecta clumps is well protected from destruction by sputtering in the reverse shock.

**Key words:** supernovae: individual: Cassiopeia A – dust, extinction – ISM: supernova remnants.

## 1 INTRODUCTION

The detections of significant ( $\gtrsim 10^8 M_{\odot}$ ) masses of dust in high-redshift quasars (Bertoldi et al. 2003; Priddey et al. 2003), and dust-enriched galaxies at redshifts  $z > 7$  (Watson et al. 2015; Laporte et al. 2017; Hashimoto et al. 2018), require an explanation of how sufficient quantities of dust can be formed at such early epochs. Core-collapse supernovae (SNe) have been proposed as a potential source of this dust (Dunne et al. 2003; Gall, Hjorth & Andersen 2011), as their progenitors evolve rapidly compared to the age of the universe at these redshifts ( $\sim 1$  Gyr). Observations of supernova remnants (SNRs) have confirmed the presence of large quantities of dust formed in the ejecta, both through infrared (IR)/sub-millimetre (sub-mm) detection of dust emission (Barlow et al. 2010; Matsuura et al. 2011; Gomez et al. 2012; Matsuura et al. 2015; De Looze et al. 2017) and alteration of emission-line profiles due to dust extinction (Bevan & Barlow 2016; Bevan, Barlow & Milisavljevic 2017).

In order to explain the observed dust masses at high redshift, the average dust yield per SNe must exceed some minimum value, estimated as  $\sim 1 M_{\odot}$  by Dwek, Galliano & Jones (2007) and between 0.1 and  $1 M_{\odot}$  by Morgan & Edmunds (2003) and Michałowski et al. (2010), although Rowlands et al. (2014) found that higher yields may be necessary if dust destruction in the interstellar medium

(ISM) is taken into account. Dust masses observed in SNRs, such as the Crab Nebula ( $0.1\text{--}0.2 M_{\odot}$ ; Gomez et al. 2012) and SN 1987A ( $0.8 M_{\odot}$ ; Matsuura et al. 2015), approach or exceed this value, but the fraction that will survive passage through the SN reverse shock and into the ISM is uncertain (Bianchi & Schneider 2007; Nozawa et al. 2007; Nozawa et al. 2010; Biscaro & Cherchneff 2016; Bocchio et al. 2016; Micelotta, Dwek & Slavin 2016). In particular, large ( $a \gtrsim 0.1 \mu\text{m}$ ) dust grains are able to survive destruction by sputtering much more effectively than smaller grains (Silvia, Smith & Shull 2010).

Cassiopeia A (Cas A) is a Galactic SNR located 3.4 kpc away (Reed et al. 1995), with an age of approximately 330 yr (Fesen et al. 2006) and a radius of 1.7 pc (Reed et al. 1995). It provides a unique laboratory to test the efficiency of dust condensation in SN ejecta, and the subsequent destruction of dust by the reverse shock, and as such has been studied extensively in the past. IR/sub-mm observations have led to derived dust masses ranging from  $\sim 10^{-4} M_{\odot}$  of hot ( $T \sim 100$  K) dust (Arendt, Dwek & Moseley 1999; Douvion, Lagage & Pantin 2001) to  $2\text{--}4 M_{\odot}$  of cold dust emitting at sub-mm wavelengths (Dunne et al. 2003), although this higher mass has been attributed to foreground dust emission in the ISM (Krause et al. 2004). Analyses of integrated fluxes from *Spitzer* and *Herschel* observations (Rho et al. 2008; Barlow et al. 2010; Arendt et al. 2014) found  $\sim 0.01 M_{\odot}$  of hot dust, with  $\sim 0.1 M_{\odot}$  of cold, unshocked dust present in the central regions, in agreement with simulations of the dust formation and evolution in

\* E-mail: fdp@star.ucl.ac.uk

Cas A by Nozawa et al. (2010). Dunne et al. (2009) suggested the observed polarization of the sub-mm emission could be explained by  $\sim 1 M_{\odot}$  of cold dust, similar to the value of  $1.1 M_{\odot}$  given by Bevan et al. (2017) as the most likely mass based on the shape of emission-line profiles affected by extinction. De Looze et al. (2017) utilized spatially resolved *Herschel* and *Spitzer* observations of Cas A to fit the dust continuum emission, following the removal of line and synchrotron contamination, using a four-component model including ISM dust emission and three SNR dust temperature components. They found a large mass of unshocked cold dust in the centre of the SNR (up to  $0.6 M_{\odot}$ ), significantly above previous estimates based on the IR/sub-mm emission.

Previous modelling of the Cas A dust emission has been based on fitting the spectral energy distribution (SED) with some number of temperature components for a given dust composition (i.e. ‘hot’ and ‘cold’ dust). This assumes all dust grains radiate at the same temperature for each component, but grains of different sizes will in general have different equilibrium temperatures for the same heating source. Additionally, smaller grains can undergo large temperature fluctuations (e.g. Purcell 1976; Draine & Anderson 1985; Dwek 1986) and may not reach an equilibrium temperature at all. Temim & Dwek (2013) found that modelling the emission from a distribution of grain sizes in the Crab Nebula reduced the dust mass estimate by a factor of 2 compared to two-temperature fits by Gomez et al. (2012), demonstrating the importance of accounting for these effects. In this paper, we calculate the emission from a population of grains subjected to conditions appropriate for the various gas components in Cas A, and use it to constrain the mass, properties, and location of the newly formed dust.

## 2 PHYSICAL PROPERTIES OF THE CAS A SNR

In order to determine the dust emission from the remnant, several physical properties are required: the densities and temperatures of the electrons and nuclei, the dominant type of nucleus, and the radiation field strength and spectrum. Observations of Cas A reveal a complex structure, with material covering a wide range of densities and temperatures emitting at different wavelengths. The supernova explosion has driven a forward shock into the circumstellar material, thought to be from the stellar wind of the progenitor (Hwang & Laming 2009), while the ejecta from the supernova itself crosses the reverse shock as it expands (DeLaney et al. 2004). Both shocks are visible as X-ray emitting regions, with typical densities of  $n \sim 1\text{--}10 \text{ cm}^{-3}$  and temperatures  $T \gtrsim 10^7 \text{ K}$  (Willingale et al. 2003; Lazendic et al. 2006; Patnaude & Fesen 2014; Wang & Li 2016). The ejecta is mostly comprised of heavy elements, principally oxygen (Chevalier & Kirshner 1979; Willingale et al. 2003). As well as the X-ray emitting gas, the shocked ejecta also consists of denser clumps or knots, emitting in the optical and IR (Hurford & Fesen 1996; DeLaney et al. 2010; Patnaude & Fesen 2014) and associated with the dust emission (Arendt et al. 1999). Electron densities in the shocked clumps are  $n_e \sim 10^3\text{--}10^5 \text{ cm}^{-3}$  (Smith et al. 2009; DeLaney et al. 2010; Lee et al. 2017), while the gas temperatures are of order  $10^4 \text{ K}$  (Arendt et al. 1999; Docenko & Sunyaev 2010). The SNR also contains ejecta that has not yet encountered the reverse shock and is consequently much cooler. Smith et al. (2009) estimated a maximum electron density of  $n_e \lesssim 100 \text{ cm}^{-3}$  for the unshocked ejecta based on forbidden line ratios, while observations of radio absorption by DeLaney et al. (2014) and Arias et al. (2018) give  $n_e \sim 10 \text{ cm}^{-3}$  and  $T \sim 100 \text{ K}$ . Raymond et al. (2018) inferred a pre-shock temperature of  $\sim 100 \text{ K}$  from [Si I] IR emission lines.

Krause et al. (2008) determined that the Cas A SN was of type IIb from a spectrum of its light echo, meaning that the progenitor star must have lost most of its hydrogen envelope pre-explosion. Young et al. (2006) suggested a progenitor with main-sequence mass of  $15\text{--}25 M_{\odot}$  and a mass at explosion of  $4\text{--}6 M_{\odot}$ , based on a comparison of stellar evolution and explosion models with observed features of the SNR. Modelling of the X-ray spectra (Vink, Kaastra & Bleeker 1996; Willingale et al. 2003) gives ejecta masses in the  $2\text{--}4 M_{\odot}$  range, with the swept-up material in the forward shock contributing an additional  $8\text{--}10 M_{\odot}$ . Based on observed emission-line strengths from Smith et al. (2009), Docenko & Sunyaev (2010), and Milisavljevic & Fesen (2013) (see Appendix B), the optical and IR-emitting knots have a total gas mass of  $0.59 M_{\odot}$ . Arias et al. (2018) estimate an unshocked ejecta gas mass of  $\sim 3 M_{\odot}$ . They noted that this conflicts with some models of the emission from the shocked regions, which suggest most of the ejecta has already passed the reverse shock (Chevalier & Oishi 2003; Laming & Hwang 2003), and that their result is sensitive to both the assumed gas temperature and whether the gas is clumped. However, most of the dust ( $0.4\text{--}0.6 M_{\odot}$ ), whose mass was derived by De Looze et al. (2017), appears to be located inside the reverse shock. The gas temperature has since been measured to be  $\sim 100 \text{ K}$  (Raymond et al. 2018), the value used by Arias et al. (2018) to give the  $3 M_{\odot}$  mass estimate, although the (unknown) degree of clumping still allows for a potentially lower unshocked ejecta mass.

The forward and reverse shocks generate synchrotron radiation, at wavelengths ranging from radio to X-ray (e.g. Hwang et al. 2004; Wang & Li 2016). The low-frequency part of the synchrotron SED can be fit by a single power law (De Looze et al. 2017), but this overestimates the flux at the higher frequencies important for dust heating. We find that a power law with an exponential cut-off, of the form  $\nu F_{\nu} \propto \nu^{0.2} \exp(-(\nu/\nu_0)^{-0.4})$  with  $\nu_0 = 900 \text{ eV}$ , provides a good fit to the radio and X-ray data from Wang & Li (2016), similar to theoretical predictions from Zirakashvili & Aharonian (2007). The total synchrotron luminosity, assuming a distance of  $3.4 \text{ kpc}$  (Reed et al. 1995), is  $8 \times 10^{36} \text{ erg s}^{-1}$ .

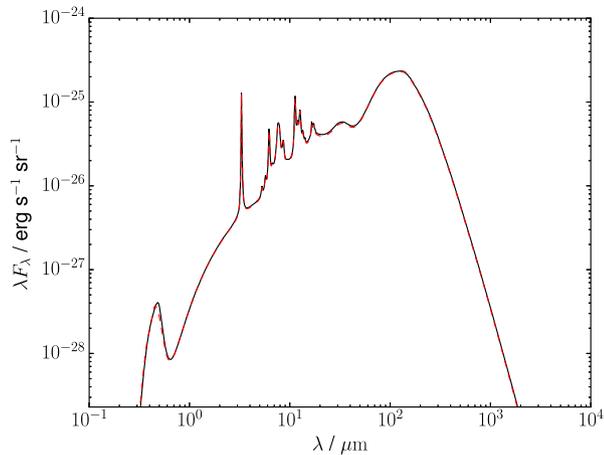
## 3 METHOD

### 3.1 DINAMO

We make use of a new dust emission code, DINAMO (DINAMO Is Not A MODified blackbody).<sup>1</sup> The code calculates the equilibrium temperature distributions and the thermal emission for a population of dust grains of arbitrary sizes and compositions, given the physical properties of the environment the grains are located in (radiation field, electron/ion density, and temperature). DINAMO includes heating from both radiation and particle collisions simultaneously, and treats cooling as a state-to-state process rather than using the continuous cooling approximation (Siebenmorgen, Kruegel & Mathis 1992). The code has been benchmarked against DustEM (Compiegne et al. 2011) using the test cases from Camps et al. (2015), and found to be in excellent agreement, as shown in Fig. 1.

The stochastic heating of dust grains is treated following the method of Guhathakurta & Draine (1989). For each grain size,  $N$  enthalpy bins are defined between a maximum and minimum enthalpy, with the probability of a grain being found in bin  $i$  at time  $t$  given by  $P_i(t)$ . The probability per unit time of a grain moving from

<sup>1</sup><https://fpriestley.github.io/dinamo>



**Figure 1.** Dust emission for the Mathis, Mezger & Panagia (1983) ISM radiation field benchmark case from Camps et al. (2015), as calculated by DINAMO (solid black line) and DustEM (dashed red line).

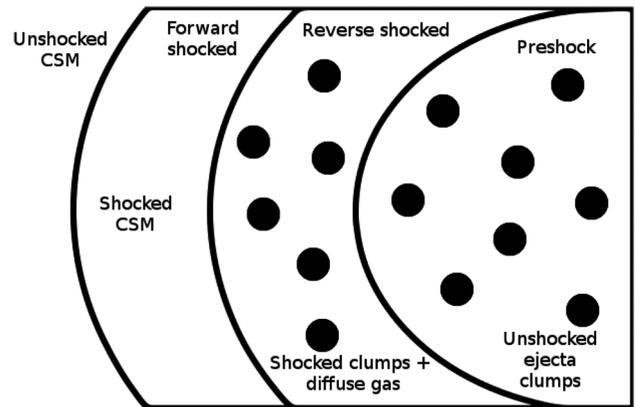
bin  $i$  to bin  $f$  is  $A_{fi}$ , with  $A_{ii} = -\sum_{f \neq i} A_{fi}$ , so that  $\frac{d}{dt} P_f = \sum_i A_{fi} P_i$ .

For the steady-state solution,  $\sum_i A_{fi} P_i = 0$  for all  $i$ , which with

the additional constraint that  $\sum_i P_i = 1$  gives a system of linear equations that can be solved to find the ‘equilibrium’ temperature distribution for a particular grain size and species. The details of the implementation of radiative and collisional heating, and cooling by thermal emission, are given in Appendix A.

### 3.2 Input parameters

Based on the above discussion, we model the SNR as consisting of four main components – the unshocked, clumped ejecta, the (reverse) shocked ejecta clumps responsible for the optical emission, the X-ray emitting diffuse reverse shocked ejecta and the X-ray emitting material swept up by the blast wave. We show a schematic diagram of the location of these components in Fig. 2. The adopted densities and temperatures of these components are listed in Table 1. For the unshocked ejecta, we take  $n_e = 100 \text{ cm}^{-3}$  and  $T = 100 \text{ K}$ , the upper limits from observations (Smith et al. 2009; Raymond et al. 2018) – we show later that lower values have no effect on the resulting emission, as the heating is dominated by the radiation field. We assume an ionization fraction of unity for convenience, whereas at these temperatures the gas is presumably mostly neutral. Arias et al. (2018) find  $n_e \sim 10 \text{ cm}^{-3}$  for the unshocked ejecta – for a nucleon density of  $n_i = 100 \text{ cm}^{-3}$ , this gives an ionization fraction of 0.1, not unreasonable for material irradiated by UV and X-ray photons and with an enhanced abundance of low ionization potential elements (e.g. Si, Fe). For the shocked clumps, we use  $n_e = 480 \text{ cm}^{-3}$ , used in deriving the total gas mass in this component, and  $T = 10^4 \text{ K}$  (Arendt et al. 1999; Docenko & Sunyaev 2010). From the ionization states of oxygen (Appendix B) in the shocked clumps, the ionization fraction must be approximately unity. As the values of  $n_e$  derived from observations range from  $10^2 - 10^5 \text{ cm}^{-3}$ , we also consider models with higher values of this parameter. For the two X-ray emitting components, we use the electron and ion temperatures and densities from Willingale et al. (2003). For the other two components, we assume that  $T_i = T_e$  and  $n_i = n_e$  – electrons are much more efficient at heating than ions for the same temperature due to their lower mass and correspondingly higher



**Figure 2.** Schematic diagram representing the assumed location of each component with respect to the forward and reverse shocks. The unshocked ejecta consists of clumped material towards the right of the diagram. Moving left, the ejecta encounters the reverse shock, compressing and heating the clumps, and ablating material that becomes the diffuse component. Further left is a contact discontinuity between the shocked ejecta and shocked circumstellar material (CSM) swept up by the blast wave of the forward shock, and the unshocked CSM that the forward shock is propagating into.

thermal velocities, so the exact value of  $n_i$  in these cases is unlikely to be significant. In the three ejecta components, we assume the ionic species is oxygen, using the heating efficiency from Dwek (1987), while in the forward shock we use hydrogen. Willingale et al. (2003) found mean ionic masses of  $15.6$  and  $1.33 m_H$  for the reverse and forward shocks, respectively, so oxygen (ejecta) and hydrogen (forward shock) are almost certainly the dominant constituents.

In addition to heating by the gas, dust in Cas A is also heated by the ambient radiation field. In Section 2, we determined the luminosity and SED of the synchrotron radiation emitted by the forward and reverse shocks, but the local intensity that enters the calculation of the dust heating depends on both the location of the dust, and the distribution of the emitting material. Based on an SNR radius of  $1.7 \text{ pc}$  (Reed et al. 1995), we assume all components are located  $1 \text{ pc}$  from a source with the luminosity of the whole remnant and explore the sensitivity of our results to changes in the intensity.

Given the oxygen-rich nature of the SNR, the dust in Cas A is expected to be primarily composed of silicates. Rho et al. (2008) and Arendt et al. (2014) found magnesium silicates of various compositions could reproduce a strong  $21 \mu\text{m}$  feature in the *Spitzer* dust emission spectra, although  $\text{Al}_2\text{O}_3$ , carbon grains, and other species were also suggested to be present. We use optical constants for magnesium and magnesium-iron silicates with varying elemental ratios (Jaeger et al. 1994; Dorschner et al. 1995; Jäger et al. 2003), which span the range  $0.2 - 500 \mu\text{m}$ . As these do not extend into the shorter-wavelength regions important for dust heating, we use the optical constants for astronomical silicates from Laor & Draine (1993) for  $0.001 - 0.2 \mu\text{m}$  for all silicate species, interpolating between the two data sets to avoid discontinuities, and we extrapolate the experimental data up to  $1000 \mu\text{m}$ . We also investigated carbon grains, using optical constants for the ACAR and BE samples from Zubko et al. (1996), extended to  $0.0003 \mu\text{m}$  with data from Uspenskii et al. (2006), as described by Owen & Barlow (2015). We assume mass densities of  $2.5$  and  $1.6 \text{ g cm}^{-3}$  for silicate and carbon grains, respectively, following De Looze et al. (2017), and sublimation temperatures of  $1500$  and  $2500 \text{ K}$ , although

**Table 1.** Adopted gas masses, ion and electron number densities and temperatures, and dominant ionic species for the four gas components.

Component	$M_{\text{gas}}/M_{\odot}$	$n_i/\text{cm}^{-3}$	$n_e/\text{cm}^{-3}$	$T_i/\text{K}$	$T_e/\text{K}$	Ion	Ref.
Pre-shock	3	100	100	100	100	O	(1), (2), (3)
Clumped	0.59	480	480	$10^4$	$10^4$	O	(4), (5)
Diffuse	1.68	7.8	61	$7.05 \times 10^8$	$5.22 \times 10^6$	O	(6)
Blastwave	8.32	14.3	16	$3.98 \times 10^8$	$3.79 \times 10^7$	H	(6)

References: (1) Smith et al. (2009); (2) Raymond et al. (2018); (3) Arias et al. (2018); (4) Docenko & Sunyaev (2010); (5) Appendix B; (6) Willingale et al. (2003).

**Table 2.** Dust species and their adopted densities  $\rho_g$ , sublimation temperatures  $T_{\text{sub}}$  and references for the optical constants.

Dust species	$\rho_g/\text{g cm}^{-3}$	$T_{\text{sub}}/\text{K}$	$n-k$
MgSiO <sub>3</sub>	2.5	1500	(1), (3)
Mg <sub>0.4</sub> Fe <sub>0.6</sub> SiO <sub>3</sub>	2.5	1500	(1), (3)
Mg <sub>0.7</sub> SiO <sub>2.7</sub>	2.5	1500	(2), (3)
Mg <sub>2.4</sub> SiO <sub>4.4</sub>	2.5	1500	(2), (3)
Am. carbon ACAR	1.6	2500	(4), (5)
Am. carbon BE	1.6	2500	(4), (5)

References: (1) Dorschner et al. (1995); (2) Jäger et al. (2003); (3) Laor & Draine (1993); (4) Zubko et al. (1996); (5) Uspenskii et al. (2006).

**Table 3.** Cas A SNR dust fluxes and uncertainties from De Looze et al. (2017), for an ISM radiation field strength  $G = 0.6G_0$ .

Waveband	$F_{\nu}/\text{Jy}$
IRAC 8 $\mu\text{m}$	$0.2 \pm 0.1$
WISE 12 $\mu\text{m}$	$3.4 \pm 0.3$
IRS 17 $\mu\text{m}$	$63.3 \pm 6.0$
WISE 22 $\mu\text{m}$	$202.0 \pm 19.3$
MIPS 24 $\mu\text{m}$	$153.4 \pm 15.0$
IRS 32 $\mu\text{m}$	$168.5 \pm 17.3$
PACS 70 $\mu\text{m}$	$149.5 \pm 20.1$
PACS 100 $\mu\text{m}$	$125.8 \pm 19.9$
PACS 160 $\mu\text{m}$	$69.9 \pm 12.0$
SPIRE 250 $\mu\text{m}$	$27.3 \pm 4.8$
SPIRE 350 $\mu\text{m}$	$10.9 \pm 1.9$
SPIRE 500 $\mu\text{m}$	$2.6 \pm 0.5$
SCUBA 850 $\mu\text{m}$	$0.4 \pm 0.1$

our results are not sensitive to the choice of this parameter. Dust properties used are summarized in Table 2.

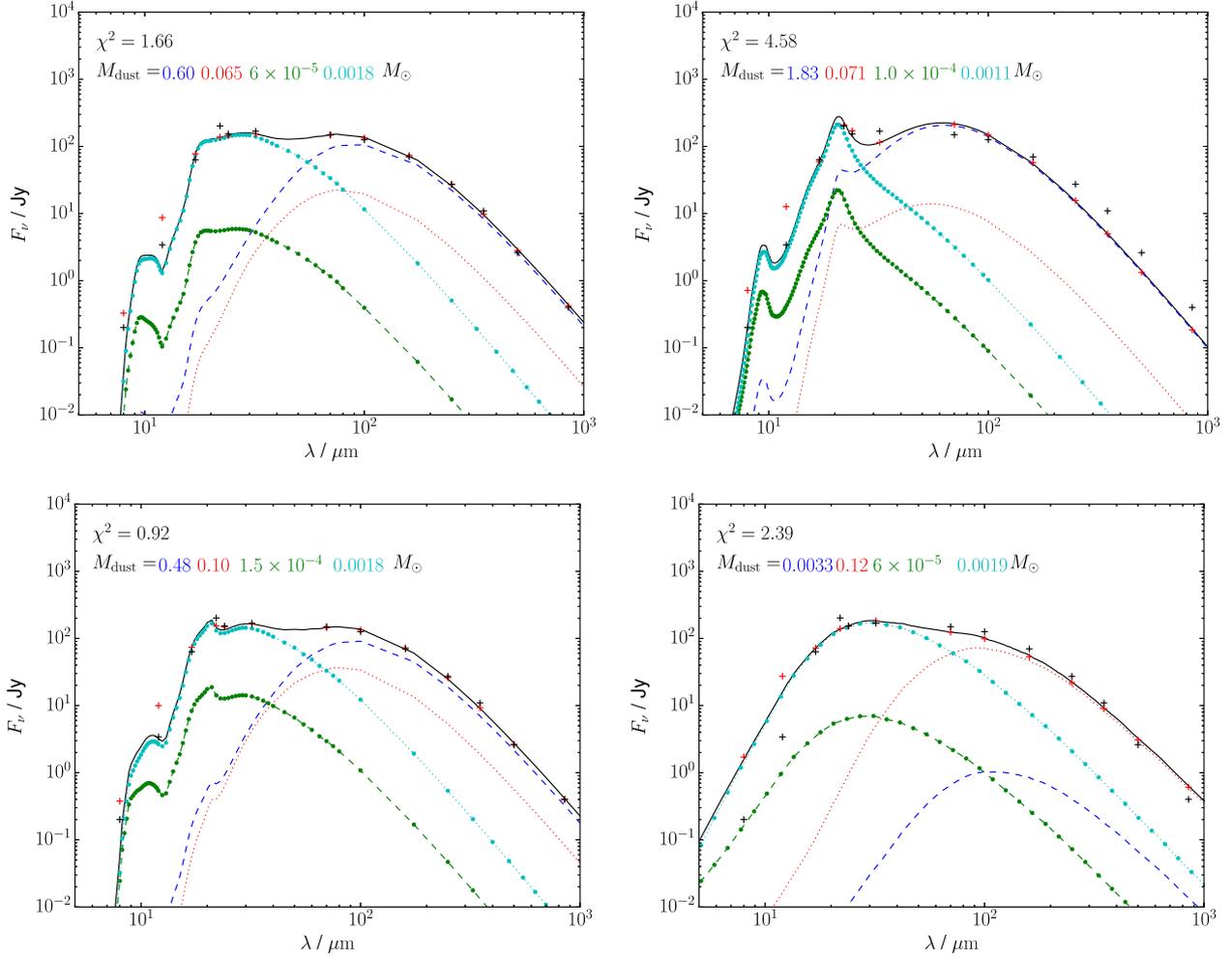
We initially assume an MRN size distribution (Mathis, Rumpl & Nordsieck 1977), with  $a_{\text{min}} = 0.005 \mu\text{m}$ ,  $a_{\text{max}} = 0.25 \mu\text{m}$  and a power-law size distribution with  $\frac{dn}{da} \propto a^{-3.5}$ , and calculate the SED per grain (averaged over size and temperature distributions) for each of the four components. As we do not consider dust self-absorption or other optically thick effects, the SEDs can be scaled to find the emission from an arbitrary number of grains, which can then be converted to a dust mass using the size distribution and grain density. Our observational data are the supernova dust fluxes reported by De Looze et al. (2017) for  $G = 0.6G_0$ , following the removal of line, synchrotron, and ISM foreground emission, as listed in Table 3. To fit the observed emission from Cas A, we run grids of models, with the number of dust grains in each component as the four free parameters, convolve the resulting SEDs with the filter profiles for each instrument and calculate the reduced  $\chi^2$  values for the models. We do not include the IRAC 8  $\mu\text{m}$  and WISE 12  $\mu\text{m}$  points in our  $\chi^2$  calculations, as at these wavelengths there is significant PAH emission from the ISM dust component

in the THEMIS-based models of De Looze et al. (2017) – as any real contribution from PAHs is highly uncertain, the reported SNR fluxes at these wavelengths cannot be used to constrain the dust emission.

## 4 RESULTS

Fig. 3 (upper left) shows our best-fitting model SED for MgSiO<sub>3</sub> grains, with optical data from Dorschner et al. (1995). The reduced  $\chi^2$  value, excluding the 8 and 12  $\mu\text{m}$  points, is 1.66. The model requires dust masses of 0.60, 0.065,  $6 \times 10^{-5}$  and  $0.0018 M_{\odot}$  for the pre-shock, clumped, diffuse and blastwave regions respectively, for a total dust mass of  $0.67 M_{\odot}$ , comparable to the value of  $0.5 \pm 0.1 M_{\odot}$  found by De Looze et al. (2017) for the same dust species. The De Looze et al. (2017) ‘hot’ dust component has a similar mass to the combined blastwave and diffuse components in our model, while their ‘warm’ component has  $\sim 10 \times$  less mass than our ‘clumped’ dust. The average (mass-weighted) dust temperatures of the four components are 29, 32, 97, and 94 K, respectively. The cold, warm, and hot dust temperatures found by De Looze et al. (2017) are 30, 79, and 100 K, respectively – the pre-shock and clumped components produce the ‘cold’ dust emission, while the diffuse and blastwave components reproduce the ‘hot’ dust. The emission from the De Looze et al. (2017) ‘warm’ dust in our models originates from a combination of the smallest grain in the two ‘cold’ components and the largest grains in the two ‘hot’ components. Our SED predicts more emission at 8 and 12  $\mu\text{m}$  than the reported SNR fluxes from De Looze et al. (2017), who did not include a high-temperature ( $\sim 500 \text{K}$ ) dust component capable of fitting the short-wavelength emission – we note that our predicted SED in this region appears similar to the ‘21  $\mu\text{m}$  peak’ IRS spectra of Cas A from Rho et al. (2008), reproducing the broad feature at  $\sim 10 \mu\text{m}$ .

Table 4 lists the best-fitting dust masses and reduced  $\chi^2$  values for each of the dust species listed in Table 2. Table 5 lists the dust-to-gas mass ratios for each component, assuming the gas masses listed in Table 1. The total dust masses for the four silicate species are similar, with the exception of Mg<sub>0.7</sub>SiO<sub>2.7</sub>, which required  $\sim 3 \times$  as much mass. This species was also found to require a significantly larger dust mass than other silicates by De Looze et al. (2017), although their value of  $21.4 M_{\odot}$  is much higher than ours, as they found a best-fitting temperature of 21 K, whereas even the largest dust grains are heated to  $\sim 30 \text{K}$  in our model, thus requiring less mass to produce the same flux. The Mg<sub>0.7</sub>SiO<sub>2.7</sub> SED is also a noticeably worse fit to the observations than the other models. Fig. 3 (upper right) shows its best-fitting SED – the slope at long wavelengths is inconsistent with that observed, and the dip in flux beyond the  $\sim 20 \mu\text{m}$  peak is too severe. Fig. 3 (lower left) shows the best-fitting SED for Mg<sub>0.4</sub>Fe<sub>0.6</sub>SiO<sub>3</sub>, the composition that gives the lowest  $\chi^2$  value of those we consider.



**Figure 3.**  $G = 0.6G_0$  Cas A dust fluxes from De Looze et al. (2017) (black crosses), and the best-fit total dust SEDs (black solid line) and model fluxes (red crosses) for grains with an MRN size distribution. The SEDs from each dust component, as defined in Table 1, are shown as blue dashed lines (pre-shock), red dotted lines (clumped), green circle-dashed lines (diffuse), and cyan circle-dotted lines (blastwave). The grain compositions are  $\text{MgSiO}_3$  (upper left),  $\text{Mg}_{0.7}\text{SiO}_{2.7}$  (upper right),  $\text{Mg}_{0.4}\text{Fe}_{0.6}\text{SiO}_3$  (lower left) and ACAR (lower right).

**Table 4.** Best-fitting model dust masses and (mass-weighted) average temperatures for each component; total dust masses and reduced  $\chi^2$  values for different grain species.

Species	$M_{\text{dust}}/M_{\odot} (< T > / \text{K})$				Total	$\chi^2$
	Pre-shock	Clumped	Diffuse	Blastwave		
$\text{MgSiO}_3$	0.60 (29)	0.065 (32)	$6 \times 10^{-5}$ (97)	0.0018 (94)	0.67	1.66
$\text{Mg}_{0.4}\text{Fe}_{0.6}\text{SiO}_3$	0.48 (29)	0.10 (33)	$1.5 \times 10^{-4}$ (97)	0.0018 (95)	0.58	0.92
$\text{Mg}_{0.7}\text{SiO}_{2.7}$	1.83 (40)	0.071 (45)	$1.0 \times 10^{-4}$ (110)	0.0011 (108)	1.90	4.58
$\text{Mg}_{2.4}\text{SiO}_{4.4}$	0.56 (29)	0.11 (33)	$6 \times 10^{-5}$ (98)	0.0018 (96)	0.67	1.32
Am. carbon ACAR	0.0033 (28)	0.12 (33)	$6 \times 10^{-5}$ (110)	0.0019 (105)	0.13	2.39
Am. carbon BE	0.002 (30)	0.19 (34)	0.0 (114)	0.0019 (108)	0.19	2.07

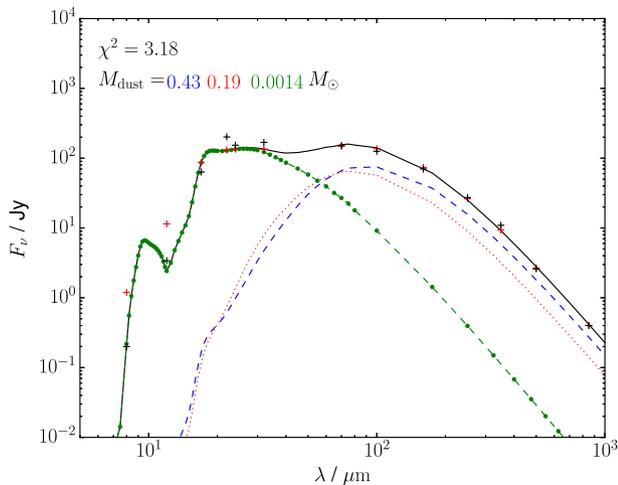
The two carbon species considered show similar behaviour, with lower dust masses than the silicates, almost all of which is contained in the clumped component, rather than in the pre-shock component as with the silicates. Given that the pre-shock gas mass is several times larger than the clumped component, it seems unlikely that the dust mass is distributed in the opposite fashion. Our dust masses are somewhat lower than the results for carbon grains from De Looze et al. (2017), again due to our models having higher grain temperatures than their ‘cold’ component, which contains most of the mass. Fig. 3 (lower right) shows the best-fitting SED for

the ACAR grains. While the model provides an acceptable fit to the  $> 20 \mu\text{m}$  data, the carbon dust SED at shorter wavelengths is drastically different to the IRS spectra from Rho et al. (2008), and the predicted flux at  $12 \mu\text{m}$  is higher than that observed, before the subtraction of line, synchrotron and ISM contributions. Carbon grains also fail to reproduce the  $21 \mu\text{m}$  feature.

The three grain species that produce good fits to the observed fluxes predict similar dust masses for each gas component, with the majority in the pre-shock region and a smaller amount in the clumps. The fitted dust masses in the two X-ray emitting components are

**Table 5.** Best-fitting model dust-to-gas mass ratios for each component for different grain species.

Species	Pre-shock	Clumped	Diffuse	Blastwave
MgSiO <sub>3</sub>	0.20	0.11	$3.6 \times 10^{-5}$	$2.2 \times 10^{-4}$
Mg <sub>0.4</sub> Fe <sub>0.6</sub> SiO <sub>3</sub>	0.16	0.17	$8.9 \times 10^{-5}$	$2.2 \times 10^{-4}$
Mg <sub>0.7</sub> SiO <sub>2.7</sub>	0.63	0.12	$6.0 \times 10^{-5}$	$1.3 \times 10^{-4}$
Mg <sub>2.4</sub> SiO <sub>4.4</sub>	0.19	0.19	$3.6 \times 10^{-5}$	$2.2 \times 10^{-4}$
Am. carbon ACAR	0.0011	0.20	$3.6 \times 10^{-5}$	$2.3 \times 10^{-4}$
Am. carbon BE	$6.7 \times 10^{-4}$	0.32	0.0	$2.3 \times 10^{-4}$


**Figure 4.**  $G = 0.6G_0$  Cas A dust fluxes from De Looze et al. (2017) (black crosses), and the best-fitting total dust SEDs (black solid line) and model fluxes (red crosses) for MgSiO<sub>3</sub> grains with an MRN size distribution, using only three ejecta components. The SEDs from each dust component as defined in Table 1 are shown as blue dashed lines (pre-shock), red dotted lines (clumped), and green circle-dashed lines (diffuse).

insignificant in terms of the total dust mass but are essential for reproducing the short wavelength data. The best-fitting models all predict most of this high-temperature dust to be in the blast wave component. However, given the similarity of the blast wave and diffuse dust SEDs, there is a large amount of degeneracy between them. Fig. 4 shows an SED fit for only the three ejecta components. While the  $\chi^2$  is worse than for models with all four components, the fit is still acceptable, with the diffuse dust mass comparable to the values found for the blast wave component in Table 4. The total dust mass is close to the four-component fit value, although the fraction of dust in the clumps is larger.

The clumped SED, although similar to the pre-shock component, has additional emission at wavelengths  $\lesssim 100 \mu\text{m}$ , which can be seen in Fig. 3, where the clumped and pre-shock SEDs are comparable at these wavelengths despite the pre-shock SED being dominant for  $\lambda > 100 \mu\text{m}$ . As the diffuse/blastwave components also emit strongly in this region, the amount of dust in the clumped component, and its contribution to the longer wavelength SED is constrained. The remainder of the  $> 100 \mu\text{m}$  emission must come from the pre-shock component, so the degeneracy is broken and the two components can be distinguished.

#### 4.1 Implications of model assumptions

While the temperatures and densities of the matter in the two X-ray emitting components have been derived from observations

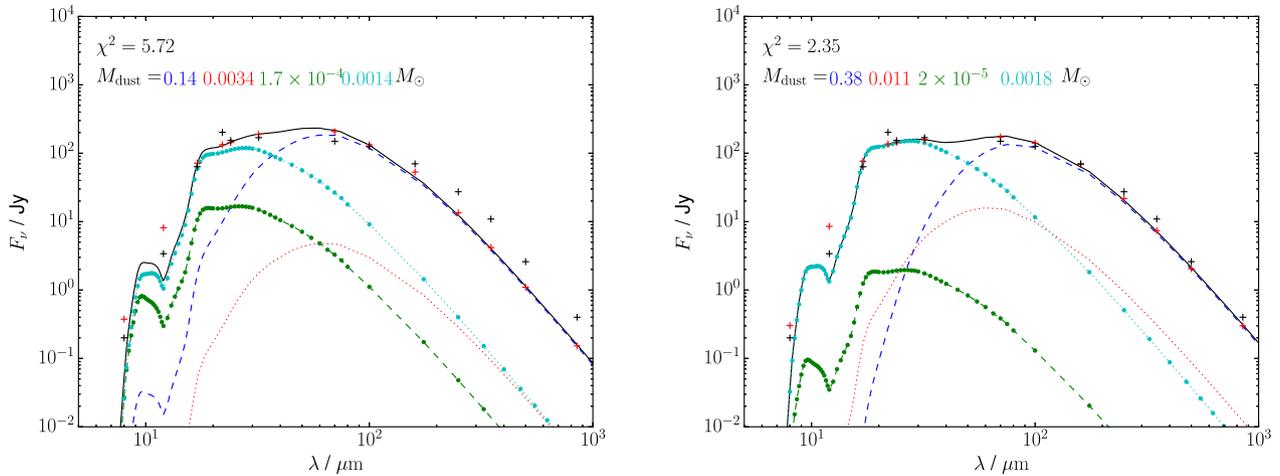
(Willingale et al. 2003), the two cooler components, which contain most of the dust mass, have more uncertain properties. The gas temperatures are fairly well constrained, while Smith et al. (2009) found an upper limit of  $n_e \lesssim 100 \text{ cm}^{-3}$  for the pre-shock ejecta, but the reported densities for the clumped material range up to  $10^5 \text{ cm}^{-3}$ , far higher than our model value of  $480 \text{ cm}^{-3}$ . Both dust components are primarily heated by the synchrotron radiation field, which we obtained by interpolating between radio and X-ray measurements and assuming a mean distance of 1 pc from the synchrotron radiation source in the blastwave – the actual radiation field could differ in both strength and spectral shape. We have also assumed an MRN size distribution for all dust components, which may not be justified given the expected growth of grains in the pre-shock component and the processing of grains in the reverse shock. In the following, we assess the impact of each of these parameters on the resulting dust masses. We use MgSiO<sub>3</sub> grains unless otherwise stated.

##### 4.1.1 Gas density

At the pre-shock ejecta temperature of 100 K (Raymond et al. 2018), the effects of collisional heating on the dust emission are negligible – even increasing the electron and ion densities to  $10^4 \text{ cm}^{-3}$ , we find no noticeable difference in the model SED. For the clumped component, a density of  $10^4 \text{ cm}^{-3}$  causes electron collisional heating to dominate the heating rate, and the dust emission becomes stronger and shifted towards shorter wavelengths. Repeating our fitting procedure, we find that the dust mass in the clumped component is reduced by a factor of a few, while the pre-shock component mass increases slightly so that the total dust mass is barely affected. The mass in the diffuse component is reduced to negligible levels in this case.

##### 4.1.2 Radiation field

To investigate the sensitivity of our results to the synchrotron radiation field, probably the least certain of our model inputs, we increased its strength by a factor of 10 for all components. Fig. 5 (left) shows the best-fitting SED for MgSiO<sub>3</sub> grains. The required dust masses in the pre-shock and clumped components are reduced to  $\sim 0.1 M_\odot$ , while the diffuse and blast wave components, which are heated by particle collisions, are less affected. The model is a poor fit, in particular to the long-wavelength data, where it predicts significantly less flux than observed. We find that this can be remedied by changing the size distribution in the unshocked component to grains only between radii of 0.1 and  $1.0 \mu\text{m}$ . This is shown in Fig. 5 (right) – the dust mass is still lower than models with the original choice of radiation field, but only by about a third, and the proportion in the pre-shock component is even higher. Increasing the radiation field by smaller factors also requires an increase in the



**Figure 5.**  $G = 0.6G_0$  Cas A dust fluxes from De Looze et al. (2017) (black crosses), and the best-fitting total dust SEDs (black solid line) and model fluxes (red crosses) for  $\text{MgSiO}_3$  grains with the synchrotron radiation field strength increased by a factor of 10. The SEDs from each dust component as defined in Table 1 are shown as blue dashed lines (pre-shock), red dotted lines (clumped), green circle-dashed lines (diffuse), and cyan circle-dotted lines (blastwave). In the left panel, all components have MRN size distributions; in the right panel, the pre-shock component values of  $a_{\min}/a_{\max}$  have been changed to 0.1/1.0  $\mu\text{m}$ .

number of large grains in the pre-shock component in order to reproduce the long-wavelength flux, with the total dust mass being mostly unchanged even for a  $5 \times$  increase in the radiation field.

#### 4.1.3 Grain size distribution

Altering the grain size power-law exponent to 2 results in an increased total dust mass of  $0.76 M_{\odot}$ , with  $0.084 M_{\odot}$  in the pre-shock component and  $0.67 M_{\odot}$  in the clumped component, although for these models these two components are almost entirely degenerate. The reduced  $\chi^2$  values is 1.24 – all models in the following discussion have  $\chi^2 \lesssim 2$ . In this case, the dust mass in the diffuse component ( $0.003 M_{\odot}$ ) dominates over the blastwave component ( $< 10^{-4} M_{\odot}$ ), although again with some degeneracy between the two. For a power-law exponent of 4, results are similar to the MRN case, although the total dust mass is slightly lower ( $0.54 M_{\odot}$ ). Changing  $a_{\min}$  and  $a_{\max}$  between 0.001–0.01 and 0.05–1.0  $\mu\text{m}$  respectively, acceptable fits can still be obtained, with the total dust mass staying within the range of 0.4– $8 M_{\odot}$  in all cases.

## 5 DISCUSSION

### 5.1 Comparison with previous results

Our total dust masses are comparable to those found by De Looze et al. (2017) – given that we use their best-fitting SNR dust fluxes as input, this is unsurprising. Although not directly equivalent, their ‘cold’ dust component corresponds to the flux produced by our pre-shock and clumped components, while their ‘warm’ and ‘hot’ dust corresponds to the two X-ray emitting components. The required dust masses are greater than several previous studies of the IR emission found for Cas A (Rho et al. 2008; Barlow et al. 2010; Arendt et al. 2014), generally  $\lesssim 0.1 M_{\odot}$  – as discussed by De Looze et al. (2017), this is due to the inclusion of better defined long wavelength data, allowing the cold dust contribution to be better determined. Modelling of Cas A’s integrated emission-line profiles by Bevan et al. (2017) found a dust mass of  $\sim 1 M_{\odot}$ , comparable to our values.

As well as the total dust mass, our preferred dust compositions are also similar to those found by De Looze et al. (2017), with various magnesium-containing silicates (with the exception of  $\text{Mg}_{0.7}\text{SiO}_{2.7}$ ) all proving acceptable. While De Looze et al. (2017) suggested some carbon dust could also be present, we find that even a small ( $\sim 10^{-3} M_{\odot}$ ) mass of carbon dust in the X-ray emitting gas predicts near-IR fluxes in conflict with observations, although larger carbon dust masses could be present in the cooler regions. Bevan et al. (2017) found that mixtures of carbon and silicate grains can explain line-profile asymmetries in Cas A, but increasing the proportion of silicates required increasingly large dust masses, up to unreasonably high values ( $6.5 M_{\odot}$  for 90 per cent silicates). Since the unshocked ejecta would be responsible for the majority of the line asymmetry, this is not necessarily in conflict with our results, although the lack of carbon in the shocked ejecta would require explanation if it makes up a significant fraction of the unshocked ejecta. Multiple dust species have been used by Rho et al. (2008) and Arendt et al. (2014) to fit the mid-IR spectrum, which we discuss further in the next section, but some form of magnesium silicate was required by both those studies.

Theoretical studies of the formation of dust in the specific case of Cas A have led to differing results. Nozawa et al. (2010) predicted a total dust mass of  $0.167 M_{\odot}$  with roughly equal quantities of carbon and magnesium silicates, and smaller contributions from other species. Bocchio et al. (2016) predicted an even higher dust mass of  $0.92 M_{\odot}$ , about half of which is in silicates and a smaller fraction in carbon, while Biscaro & Cherchneff (2016) predicted only  $\sim 10^{-2} M_{\odot}$  with the main dust component being  $\text{Al}_2\text{O}_3$ . Of these studies only that of Bocchio et al. (2016) is consistent with our results – the others predict dust masses significantly smaller than our required values and of different composition. While we did not investigate  $\text{Al}_2\text{O}_3$ , De Looze et al. (2017) found that the required mass of aluminium for this composition significantly exceeded that expected from calculations of the nucleosynthetic yields by Woosley & Weaver (1995).

### 5.2 Heating mechanisms

As mentioned previously, the pre-shock and clumped dust components are heated primarily by the synchrotron radiation field,

while the diffuse and blastwave components are heated by particle collisions. However, the additional collisional heating due to the higher temperature in the clumped component causes a subtle but important difference in the dust SED, compared to the pre-shock dust, which results in our best-fitting models having the majority of the dust mass in the pre-shock ejecta. This shows the importance of treating both processes simultaneously. Bocchio et al. (2013) have previously investigated the effects of the addition of electron collisional heating in the ISM, but calculations of dust heating in SNRs have generally considered only radiative (e.g. Temim & Dwek 2013; Owen & Barlow 2015) or collisional (Dwek 1987; Dwek et al. 2008) effects.

We also find that in the diffuse and blastwave components, the heating rate from collisions with nuclei is comparable to that from electrons ( $\sim 0.3$ – $0.5$  of the total heating rate), with the difference between  $T_i$  and  $T_e$  enough to counteract the higher mass of the ions (and subsequent reduction in velocities and collision rates). As this non-equilibrium between ions and electrons is expected to be a general feature of shocks in SNRs (Raymond 2018), it is not necessarily true that electrons will be the dominant source of dust heating even in situations where collisional heating can be expected to dominate.

### 5.3 The mid-IR spectrum

Rho et al. (2008) and Arendt et al. (2014) both used *Spitzer* IRS 5–30  $\mu\text{m}$  spectra of Cas A, in addition to longer-wavelength photometric data extending to 160  $\mu\text{m}$ , to investigate the dust composition in different regions of the SNR. Rho et al. (2008) categorized the spectra based on the strength of the 21  $\mu\text{m}$  peak compared to the continuum, whereas Arendt et al. (2014) defined various regions of the remnant characterized by particular dominant emission features (line or continuum), and treated the dust emission from these regions as separate populations. Rho et al. (2008) found 0.02–0.05  $M_\odot$  of dust to be responsible, with the dominant materials being carbon and simple iron and silicate species (e.g. FeO and SiO<sub>2</sub>), although also requiring some form of magnesium silicate and Al<sub>2</sub>O<sub>3</sub>. Arendt et al. (2014) found  $\sim 0.04 M_\odot$  of dust to be needed, which could mostly be fit by a combination of magnesium silicates and some additional featureless dust component (carbon, Al<sub>2</sub>O<sub>3</sub>, or Fe/FeS/Fe<sub>3</sub>O<sub>4</sub>), plus an additional featureless dust component of  $\lesssim 0.1 M_\odot$  associated with the [Si II] emission, which they could not constrain the composition of but associated with the unshocked ejecta.

Both studies found dust masses lower than our values by an order of magnitude – this is again due to the additional long wavelength fluxes we use, where the majority of the dust mass emits. Arendt et al. (2014) did include Herschel PACS fluxes to constrain the dust masses, in particular the [Si II] component, from which they determined an upper limit of 0.1  $M_\odot$  of unshocked dust, well below our values of  $\sim 0.5 M_\odot$ . However, in determining the ISM dust contribution at 160  $\mu\text{m}$ , they assumed a scaled version of the SPIRE 250  $\mu\text{m}$  map to represent the ISM dust emission, whereas De Looze et al. (2017) find a substantial SNR contribution to the 250  $\mu\text{m}$  flux. Their SNR 160  $\mu\text{m}$  fluxes, and the derived dust mass, are therefore lower than the values from De Looze et al. (2017) that took into account ISM and SNR dust emission simultaneously.

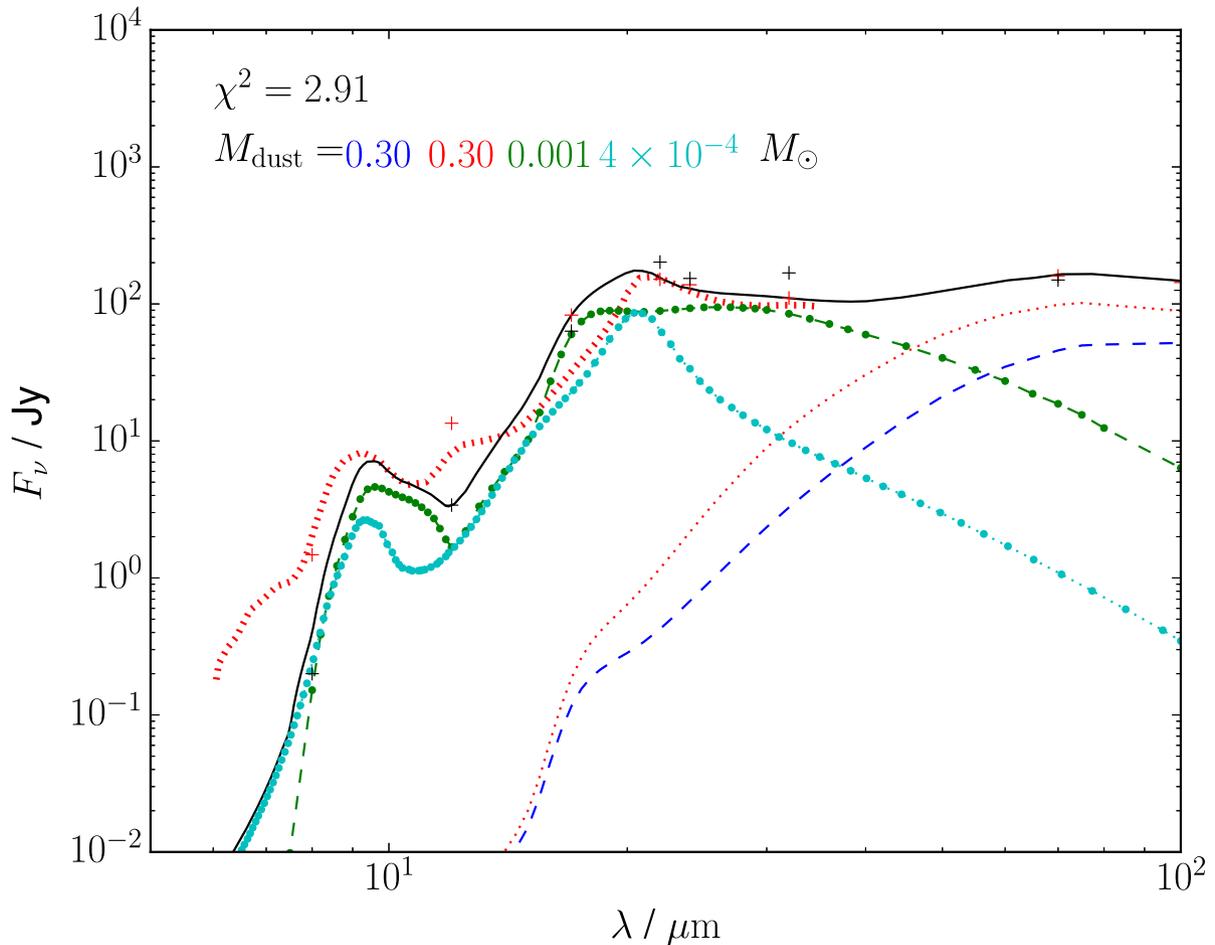
Although there are variations from position to position, the main features of the mid-IR spectra are two emission peaks of varying strength at  $\sim 10$  and 21  $\mu\text{m}$ , and an underlying continuum that rises to 21  $\mu\text{m}$  and stays roughly constant or falls to longer wavelengths. While Rho et al. (2008) and Arendt et al. (2014) used

various combinations of dust species and temperature components to fit this behaviour, our model SEDs approximately reproduce this naturally using single silicate species – the two peaks and the rising part of the continuum are produced by grains in the high temperature X-ray emitting parts of the ejecta, while dust in the cooler regions produces an increasing fraction of the flux towards longer wavelengths, preventing the decline found with single-temperature  $\sim 100$  K SEDs. Some IRS spectra [‘featureless’ in Rho et al. (2008) and [‘Ne II’] in Arendt et al. (2014)] required an additional dust component at  $\sim 10$  K to explain featureless emission, possibly carbon or Al<sub>2</sub>O<sub>3</sub>, but we find that the observations are consistent with magnesium silicate grains making up the majority of the dust present in the SNR.

Our best-fitting magnesium silicate models all underpredict the flux at 21  $\mu\text{m}$ , with the exception of Mg<sub>0.7</sub>SiO<sub>2.7</sub>, which otherwise provides a poor fit. The shape of our predicted SED in this region also clearly differs from many of the individual IRS spectra (the ‘21  $\mu\text{m}$  peak dust’ of Rho et al. 2008), suggesting that some additional component may be contributing at this wavelength. Arendt et al. (2014) attribute this component to Mg<sub>0.7</sub>SiO<sub>2.7</sub>, while Rho et al. (2008) adopt a combination of FeO and SiO<sub>2</sub>. By including an additional dust component from one of these species, assuming FeO and SiO<sub>2</sub> can be treated as silicates and using the diffuse X-ray emitting gas properties, we find that this 21  $\mu\text{m}$  excess can be reproduced with the addition of  $\sim 5 \times 10^{-4} M_\odot$  of dust. FeO and Mg<sub>0.7</sub>SiO<sub>2.7</sub> both produce similar SEDs, while SiO<sub>2</sub> produces an additional sharp peak at  $\sim 12 \mu\text{m}$ , similar to that seen in some of the IRS spectra (e.g. fig. 3 of Rho et al. 2008) and not produced by other silicate species. Fig. 6 shows the best-fitting dust SED for MgSiO<sub>3</sub> as before, except with the blastwave component replaced by Mg<sub>0.7</sub>SiO<sub>2.7</sub> using diffuse component properties, and the [Ar II] dust spectrum from Arendt et al. (2014; their fig. 3), scaled to the intensity of the total Cas A IRS spectrum at 21  $\mu\text{m}$ . We note that the model SED is a reasonable fit to the observed spectrum, in particular the feature at  $\sim 10 \mu\text{m}$ , despite only being fit to the flux points at  $> 17 \mu\text{m}$  and not to the spectral data at all.

### 5.4 Dust-to-gas mass ratios

For the three dust species considered that produce acceptable fits to the observed SED (MgSiO<sub>3</sub>, Mg<sub>0.4</sub>Fe<sub>0.6</sub>SiO<sub>3</sub> and Mg<sub>2.4</sub>SiO<sub>4.4</sub>), we find similar dust masses in each of the four components –  $\sim 0.5$ – $0.6 M_\odot$  in the pre-shock ejecta,  $\sim 0.1 M_\odot$  in the clumps, and  $10^{-4}/2 \times 10^{-3} M_\odot$  in the X-ray emitting reverse shock/blast wave gas. Taking the gas masses of these components as 3.0 (Arias et al. 2018), 0.59 (Appendix B), and 1.68 and 8.32  $M_\odot$  (both from Willingale et al. 2003), respectively, this gives dust-to-gas mass ratios of 0.20, 0.17,  $6.0 \times 10^{-5}$ , and  $2.4 \times 10^{-4}$  (Table 5), compared to a typical Galactic ISM value of  $\sim 0.0067$  (Draine 2011). The gas mass of 3  $M_\odot$  obtained by Arias et al. (2018) for the pre-shock component was noted by those authors to be higher than expected by many models of the emission from the shocked gas for the Cas A SNR, and could be lower if the unshocked ejecta gas has a lower temperature than assumed, or has a clumpy structure. Raymond et al. (2018) found a temperature of  $\sim 100$  K for this component, the same value used by Arias et al. (2018) to derive the 3  $M_\odot$  gas mass, but the level of clumping is still uncertain. The dust-to-gas ratio for this component may therefore be a lower limit (assuming our dust mass is accurate), while for the other components the gas masses are better defined. We note that assuming a lower total (pre-shock + clumped + diffuse) ejecta gas mass (e.g. 2  $M_\odot$ ; Laming & Hwang 2003) would result in extremely high ( $\gtrsim 1$ ) dust-to-gas ratios in



**Figure 6.**  $G = 0.6G_0$  Cas A dust fluxes from De Looze et al. (2017; black crosses), and the best-fitting total dust SEDs (black solid line) and model fluxes (red crosses) for  $\text{MgSiO}_3$  and  $\text{Mg}_{0.7}\text{SiO}_{2.7}$  grains with an MRN size distribution. The SEDs from dust components as defined in Table 1 are shown as blue dashed lines (pre-shock), red dotted lines (clumped), green circle-dashed lines (diffuse  $\text{MgSiO}_3$ ), and cyan circle-dotted lines (diffuse  $\text{Mg}_{0.7}\text{SiO}_{2.7}$ ). The ‘[Ar II] region’ dust spectrum from Arendt et al. (2014) is shown as a stippled red curve.

either the pre-shock or clumped components, given the relatively well-constrained diffuse mass and the need for  $\sim 0.5 M_\odot$  of dust in one of the other two components to reproduce the far-IR flux.

The ejecta dust-to-gas ratios are  $\sim 20 \times$  higher than in the ISM, implying that a significant fraction of the metals are condensed into dust grains (0.17, assuming all the ejecta mass is condensible material). Raymond et al. (2018) estimated a similar efficiency of  $\sim 0.1$  from the gas and dust masses from Arias et al. (2018) and De Looze et al. (2017), respectively, combined with their measurement of the pre-shock gas temperature. Owen & Barlow (2015) found lower ratios (0.026–0.038) for the Crab Nebula, which has no reverse shock processing the ejecta; however, the Crab Nebula contains significant quantities of hydrogen and helium, unlike Cas A, so the fraction of condensible material locked up in dust grains is higher than the value inferred from this ratio. Nozawa et al. (2010) found a condensation efficiency of 0.13 in their model of dust formation in Cas A, although they predicted a lower dust mass than our models require, and also predicted that the main dust component is carbon, which we rule out. The lower-than-ISM dust-to-gas mass ratios in the two diffuse shocked components are consistent with significant dust destruction by both the forward and reverse shocks. Even if the blastwave dust mass is attributed instead to the reverse shock (which is possible, due to the degeneracy between the two

components), the dust-to-gas ratio would be 0.0013, implying that  $< 1$  per cent of the original dust mass has survived, assuming the original ratio is that of the pre-shock component. We note that dust condensation efficiencies are expected to be lower in the outer ejecta than the centre, which may suggest a lower initial dust-to-gas ratio for the material which has been processed by the reverse shock at present, and a correspondingly higher survival fraction. The surviving dust mass in the Cas A model of Bocchio et al. (2016) is  $\sim 1$  per cent of the initial mass.

Conversely, we find that the dust-to-gas ratio in the clumps is very similar to that in the pre-shock ejecta, suggesting that dust within the clumps has been protected from destruction by the reverse shock. Biscaro & Cherchneff (2016) predicted a surviving dust mass fraction of 6–11 per cent for clumps in the Cas A SNR, while Micelotta et al. (2016) predicted a surviving fraction of 12 per cent for silicate grains, lower than our implied value even for conservative estimates of the unshocked ejecta gas mass – for a pre-shock gas mass of  $1 M_\odot$  the surviving fraction is 28 per cent, rising to 84 per cent for the  $3 M_\odot$  value given by Arias et al. (2018). Dust masses from an increased radiation field model (Fig. 5) suggest a surviving fraction of 15 per cent, closer to the theoretical estimates, although this would require all the dust to be located at an average distance of  $\sim 0.3$  pc from the synchrotron radiation

source. Biscaro & Cherchneff (2016) and Micelotta et al. (2016) include sputtering of dust grains expelled from the clumps into the inter-clump medium – accounting only for sputtering within the clumps before injection into the inter-clump medium, the value from Biscaro & Cherchneff (2016) is 28–58 per cent, which is consistent with our results if the unshocked ejecta gas mass was  $\sim 1\text{--}2 M_{\odot}$ . The dust production efficiency of SNRs, and their overall contribution to the dust budget in the ISM, can therefore be strongly affected by the degree of clumping and the detailed, multidimensional hydrodynamical evolution of the ejecta as it passes through the reverse shock.

## 6 CONCLUSIONS

We have modelled the emission from dust grains subjected to the physical conditions present in the Cas A SNR, accounting for radiative and collisional heating mechanisms using a new dust emission code DINAMO, and used the SNR dust fluxes describing the IR SED from 17 to  $850 \mu\text{m}$  from De Looze et al. (2017) to constrain both the mass of dust present, and its distribution between the various components of the remnant (unshocked ejecta, shocked ejecta clumps, diffuse shocked ejecta, and material swept up by the blast wave). We find dust masses of  $\sim 0.6\text{--}0.7 M_{\odot}$  depending on the silicate composition, with the majority being located in the unshocked ejecta ( $\sim 90$  per cent) and post-shock clumps ( $\sim 10$  per cent), and only a small fraction present in the hot X-ray emitting gas. The dust-to-gas ratio in the shocked clumpy ejecta is  $\sim 0.17$ , similar to that in the unshocked region ( $\sim 0.2$ ), while in the diffuse components it is significantly lower ( $\lesssim 0.001$ ). This is consistent with dust grains being efficiently sputtered at high temperatures, whereas in the ejecta clumps that have passed through the reverse shock, the dust is more resilient to destruction. Magnesium silicate grains, with possible iron inclusions, are found to reproduce almost all of the observed Cas A dust spectrum, with a relatively minor amount of another species ( $\text{FeO}$ ,  $\text{SiO}_2$ , or  $\text{Mg}_{0.7}\text{SiO}_{2.7}$ ) required to reproduce the  $21 \mu\text{m}$  emission peak. While carbon grains may be present, they cannot make up a large fraction of the dust mass in the X-ray emitting gas ( $< 25$  per cent) without predicting NIR fluxes in excess of those observed. If the mass fractions do not vary significantly between the shocked and unshocked ejecta, carbon dust can be ruled out as a major constituent of the ejecta dust. The unshocked and clumped ejecta dust, making up the majority of the mass, is heated mostly by the remnant's synchrotron radiation field, while the diffuse and blastwave dust, which dominates the total SED luminosity, is heated by collisions with electrons and nuclei. The total dust mass in Cas A is consistent with CCSNe being significant contributors to the dust in high-redshift galaxies, particularly if much of it is present in clumps that survive the passage of the reverse shock without disruption.

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## APPENDIX A: DUST HEATING AND COOLING RATES

### A1 Radiative heating

A dust grain in enthalpy bin  $i$ , with enthalpy  $H_i$ , absorbing a photon of energy  $h\nu$  will increase its enthalpy to  $H_i + h\nu$ , which may move it into a higher enthalpy bin. The transition rate between bins  $i$  and  $f$  due to radiative heating is given by

$$A_{fi} = 4\pi^2 a^2 Q_{\lambda_{fi}} J_{\lambda_{fi}} \frac{hc \Delta H_f}{(H_f - H_i)^3} \quad (\text{A1})$$

for  $f > i$ , where  $a$  is the grain radius,  $\lambda_{fi} = \frac{hc}{H_f - H_i}$  is the wavelength of a photon with energy corresponding to the difference between enthalpy bins,  $Q_{\lambda_{fi}}$  and  $J_{\lambda_{fi}}$  are the absorption efficiency and radiation field strength at wavelength  $\lambda_{fi}$  and  $\Delta H_f$  is the width of enthalpy bin  $f$  (Camps et al. 2015). Photons energetic enough to heat the grain beyond the highest enthalpy bin are included in the rate to bin  $N$ , giving an additional term

$$A'_{Ni} = 4\pi^2 a^2 \int_0^{\lambda_{\max}} \lambda Q_{\lambda} J_{\lambda} d\lambda / hc, \quad (\text{A2})$$

where  $\lambda_{\max} = \frac{hc}{H_{\max} - H_i}$  is the wavelength of the least energetic photon capable of heating a grain beyond the maximum temperature. Photons which are not energetic enough to heat a grain out of the enthalpy bin contribute to a continuous heating rate

$$\frac{dH_{\text{heat}}}{dt} = 4\pi^2 a^2 \int_{\lambda_0}^{\infty} Q_{\lambda} J_{\lambda} d\lambda, \quad (\text{A3})$$

where  $\lambda_0 = \frac{hc}{H_f - H_i}$ . If the continuous heating rate is greater than the equivalent cooling rate, then

$$A_{fi}^{\text{cont}} = \frac{1}{\Delta H_i} \frac{dH_{\text{net}}}{dt} \quad (\text{A4})$$

for  $f = i + 1$ , where  $\Delta H_i$  is the width of bin  $i$  and  $H_{\text{net}}$  is the net heating rate  $\frac{dH_{\text{heat}}}{dt} - \frac{dH_{\text{cool}}}{dt}$ .

### A2 Radiative cooling

Dust grains of temperature  $T$  emit radiation at wavelength  $\lambda$  with intensity  $Q_{\lambda} B(\lambda, T)$ , where  $B(\lambda, T)$  is the Planck function, causing them to lose energy. The transition rates are similar to those for absorption of a photon, with

$$A_{fi} = 4\pi^2 a^2 Q_{\lambda_{fi}} B(\lambda, T_i) \frac{hc \Delta H_f}{(H_i - H_f)^3} \quad (\text{A5})$$

for  $f < i$ , where  $T_i$  is the temperature of a grain in enthalpy bin  $i$ , and

$$A'_{ii} = 4\pi^2 a^2 \int_0^{\lambda_{\min}} \lambda Q_{\lambda} B(\lambda, T_i) d\lambda / hc, \quad (\text{A6})$$

where  $\lambda_{\min} = \frac{hc}{H_i - H_{\min}}$ . The continuous cooling rate is given by

$$\frac{dH_{\text{cool}}}{dt} = 4\pi^2 a^2 \int_{\lambda_0}^{\infty} Q_{\lambda} B(\lambda, T_i) d\lambda, \quad (\text{A7})$$

where  $\lambda_0 = \frac{hc}{H_i - H_f}$ , and for  $\frac{dH_{\text{cool}}}{dt} > \frac{dH_{\text{heat}}}{dt}$ ,

$$A_{fi}^{\text{cont}} = -\frac{1}{\Delta H_i} \frac{dH_{\text{net}}}{dt} \quad (\text{A8})$$

for  $f = i - 1$ .

### A3 Collisional heating

As with photons, a collision between a dust grain and a particle (either an electron or an atom/ion) can result in a transfer of energy to the dust grain. Unlike with photons, a colliding particle does not necessarily transfer all its energy to the grain, and the amount of heating depends on the particle energy as well as the dust properties. The transition rate between enthalpy bins due to particle heating is given by

$$A_{fi} = \pi a^2 n \int f(E) v(E) \delta(\Delta E) dE, \quad (\text{A9})$$

where  $n$  is the number density of particles,  $f(E)$  is the probability distribution of particle energies,  $v(E)$  is the velocity of a particle with energy  $E$ ,  $\Delta E$  is the energy transferred to the dust grain, and  $\delta(\Delta E)$  is a function such that

$$\delta(\Delta E) = \begin{cases} 0 & |\Delta E - (H_f - H_i)| < \Delta H_f / 2, \\ 1 & \text{otherwise,} \end{cases} \quad (\text{A10})$$

where  $\Delta H_f$  is the width of enthalpy bin  $f$ . The additional heating rate to enthalpies higher than  $H_N$  is given by

$$A'_{Ni} = \pi a^2 n \int f(E) v(E) \delta'(\Delta E) dE, \quad (\text{A11})$$

where

$$\delta'(\Delta E) = \begin{cases} 0 & \Delta E < H_{\max} - H_i \\ 1 & \text{otherwise} \end{cases} \quad (\text{A12})$$

and the continuous heating rate is given by

$$\frac{dH_{\text{heat}}}{dt} = \pi a^2 n \int f(E) v(E) \delta''(\Delta E) \Delta E dE, \quad (\text{A13})$$

where

$$\delta''(\Delta E) = \begin{cases} 1 & \Delta E < H_f - H_i \\ 0 & \text{otherwise} \end{cases} \quad (\text{A14})$$

for  $f = i + 1$ . For electrons, the transferred energy  $\Delta E$  is determined as a function of  $E$  using the method described by Dwek & Smith (1996). For a dust grain of stopping thickness  $R_0 = 4a\rho/3$  where  $\rho$  is the density, if the electron range  $R_1(E) \leq R_0$  then  $\Delta E = E$ . For  $R_1(E) > R_0$ ,  $\Delta E = E - E'$  where  $R(E') = R_1 - R_0$ . A function for  $R(E)$  based on fits to experimental data is given in Dwek & Smith (1996). For atoms and ions, Dwek (1987) gives the transferred energy as

$$\Delta E = \begin{cases} E & E \leq E' \\ E' & E > E' \end{cases}, \quad (\text{A15})$$

where  $E'$  is listed for various nuclei as a function of grain radius  $a$  in Dwek (1987).

## APPENDIX B: THE GAS MASS FOR THE CLUMPED COMPONENT

The ejecta in Cas A is primarily composed of oxygen (Chevalier & Kirshner 1979; Willingale et al. 2003). For neutral and doubly ionized oxygen, integrated line fluxes were estimated by scaling up the mean [O I] 63  $\mu\text{m}$  and [O III] 52  $\mu\text{m}$  fluxes (per LWS

aperture) from Docenko & Sunyaev (2010) by a factor of 10.5 to account for the ratio of remnant to aperture area, giving  $1.3 \times 10^{-10} \text{ erg cm}^{-2} \text{ s}^{-1}$  and  $5.4 \times 10^{-10} \text{ erg cm}^{-2} \text{ s}^{-1}$ , respectively. For triply ionized oxygen, the integrated [O IV] 25.9  $\mu\text{m}$  flux from Smith et al. (2009) is  $4.95 \times 10^{-10} \text{ erg cm}^{-2} \text{ s}^{-1}$ . Singly ionized oxygen, which has no IR transitions, has an [O II] 7325  $\text{\AA}$  line flux of  $5.93 \times 10^{-13} \text{ erg cm}^{-2} \text{ s}^{-1}$  in the integrated Cas A spectrum of Milisavljevic & Fesen (2013). The flux in the nearby [Ar III] 7136  $\text{\AA}$  line is  $8.13 \times 10^{-14} \text{ erg cm}^{-2} \text{ s}^{-1}$ , while the [Ar III] 8.99  $\mu\text{m}$  integrated flux from Smith et al. (2009) is  $1.24 \times 10^{-10} \text{ erg cm}^{-2} \text{ s}^{-1}$ . The predicted ratio of  $F(8.99 \mu\text{m})/F(7136 \text{\AA})$  from EQUIB (Howarth & Adams 1981), using an electron density of  $480 \text{ cm}^{-3}$  [the mean value from the [O III] 52  $\mu\text{m}/88 \mu\text{m}$  line ratios from Docenko & Sunyaev (2010)] and a temperature of  $10^4 \text{ K}$ , is 0.898 – the observed 7136  $\text{\AA}$  flux is therefore 1698 times weaker than expected, corresponding to 8.07 mag of extinction. Using the Galactic extinction law from Howarth (1983) with  $R = 3.1$ , the expected extinction at 7325  $\text{\AA}$  is 7.77 mag, and the dereddened integrated [O II] 7325  $\text{\AA}$  flux is therefore  $7.64 \times 10^{-10} \text{ erg cm}^{-2} \text{ s}^{-1}$ . With the same values for  $n_e$  and  $T$ , and a distance to the SNR of 3.4 kpc (Reed et al. 1995), we use EQUIB to determine the number of oxygen nuclei at each stage of ionization and the corresponding masses, given in Table B1. The total oxygen gas mass is found to be  $0.49 M_{\odot}$ . Docenko & Sunyaev (2010) find that  $N(\text{other})/N(\text{O}) = 0.10$ , where  $N(\text{other})$  is the number of all non-oxygen nuclei, with sulphur comprising half of  $N(\text{other})$ . Assuming a mass per ‘other’ nuclei of  $32m_{\text{H}}$ , the ratio  $M(\text{other})/M(\text{O})$  is 0.20, and the total gas mass of the optical/IR-emitting clumps is  $0.59 M_{\odot}$ .

**Table B1.** Deduced number of oxygen nuclei in various ionization states, and corresponding masses, in the optical/IR-emitting gas component in the Cas A SNR.

Ion stage	O <sup>0</sup>	O <sup>+</sup>	O <sup>2+</sup>	O <sup>3+</sup>	Total
Number	$5.29 \times 10^{54}$	$2.91 \times 10^{55}$	$1.74 \times 10^{54}$	$2.01 \times 10^{53}$	$3.63 \times 10^{55}$
Mass/ $M_{\odot}$	0.07	0.39	0.02	0.003	0.49

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