



# A Tale of Three: Magnetic Fields along the Orion Integral-shaped Filament as Revealed by the JCMT BISTRO Survey

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## Abstract

As part of the *B*-fields In Star-forming Region Observations survey, we present James Clerk Maxwell Telescope (JCMT) 850  $\mu\text{m}$  polarimetric observations toward the Orion integral-shaped filament (ISF) that covers three portions known as OMC-1, OMC-2, and OMC-3. The magnetic field threading the ISF seen in the JCMT POL-2 map appears as a tale of three: pinched for OMC-1, twisted for OMC-2, and nearly uniform for OMC-3. A multiscale analysis shows that the magnetic field structure in OMC-3 is very consistent at all the scales, whereas the field structure in OMC-2 shows no correlation across different scales. In OMC-1, the field retains its mean orientation from large to small scales but shows some deviations at small scales. Histograms of relative orientations between the magnetic field and filaments reveal a bimodal distribution for OMC-1, a relatively random distribution for OMC-2, and a distribution with a predominant peak at  $90^\circ$  for OMC-3. Furthermore, the magnetic fields in OMC-1 and OMC-3 both appear to be aligned perpendicular to the fibers, which are denser structures within the filament, but the field in OMC-2 is aligned along with the fibers. All these suggest that gravity, turbulence, and magnetic field are each playing a leading role in OMC-1, 2, and 3, respectively. While OMC-2 and 3 have almost the same gas mass, density, and nonthermal velocity dispersion, there are on average younger and fewer young stellar objects in OMC-3, providing evidence that a stronger magnetic field will induce slower and less efficient star formation in molecular clouds.

*Unified Astronomy Thesaurus concepts:* [Star formation \(1569\)](#); [Interstellar magnetic fields \(845\)](#); [Interstellar clouds \(834\)](#); [Polarimetry \(1278\)](#)

## 1. Introduction

During the star formation process, the dynamics and physical states of the molecular clouds are influenced by various physical mechanisms, especially self-gravity, turbulence, and magnetic field (*B*-field; C. F. McKee & E. C. Ostriker 2007). It has long been a subject of intense debate as to which force is playing a dominant role in regulating the cloud collapse and fragmentation (M.-M. Mac Low & R. S. Klessen 2004; T. C. Mouschovias et al. 2006; R. M. Crutcher 2012). Regarding the *B*-field, either the “strong-field models” that support a defining role played by the *B*-field (e.g., T. C. Mouschovias et al. 2006) or the “weak-field models” that pay more attention to turbulence (e.g., M.-M. Mac Low & R. S. Klessen 2004) cannot sufficiently explain all the observations toward star formation regions. The relative importance of turbulence and the *B*-field as well as their interactions with self-gravity in star formation remain to be explored in more case studies (H.-B. Li 2021). More reasonable scenarios may need to consider the essential roles of both processes, which have been explored in simulations (R. M. Crutcher 2012; P. Hennebelle & S. Inutsuka 2019).

Dense molecular filaments are important sites for star formation, with molecular gas accumulating and then fragmenting into star-forming cores due to gravitational instability (P. André et al. 2014; A. Hacar et al. 2023; J. E. Pineda et al. 2023). Observations have shown that *B*-fields appear to be perpendicular to high-density filaments, while they appear to be parallel to low-density elongated clouds or striations (e.g., N. L. J. Cox et al. 2016). Magnetic fields may also play a central role in shaping the fragmentation and physical states of filaments (e.g., Y.-W. Tang et al. 2019; D. Arzoumanian et al. 2021). More observations and dedicated studies are needed to reveal the relative importance of *B*-fields compared to other processes and to decipher how *B*-fields influence the gas dynamics during filament formation and fragmentation.

Situated at the head of the Orion A giant molecular cloud, the integral-shaped filament (ISF) is a well-known nearby star-forming filament (D. Johnstone & J. Bally 1999; J. Bally 2008) containing several portions, of which the more extensively studied are OMC-1, OMC-2, and OMC-3. Several studies present *B*-field results of the whole ISF (e.g., M. Houde et al. 2004; B. C. Matthews et al. 2009) or its portions OMC-1 (e.g., D. Ward-Thompson et al. 2017; D. T. Chuss et al. 2019;

H. Ajeddig et al. 2022), OMC-2/3 (F. Poidevin et al. 2010; P. S. Li et al. 2022; N. Zielinski & S. Wolf 2022), and OMC-4 (P. S. Li et al. 2022). With active massive star formation, the *B*-field in OMC-1 has been detected with a large-scale hourglass morphology associated with two molecular clumps, namely, Orion BN/KL and South (e.g., D. A. Schleuning 1998; K. Pattle et al. 2017; D. Ward-Thompson et al. 2017). The *B*-field orientations in OMC-2 exhibit more variations compared to the other portions of the ISF (e.g., F. Poidevin et al. 2010). As for OMC-3, observations have revealed a more ordered *B*-field (e.g., B. C. Matthews et al. 2001). Therefore, being the nearest filamentary molecular cloud (393 pc; J. E. Gröbbschedl et al. 2018) forming both massive and intermediate-to-low-mass stars, the OMC-1/2/3 region shows hints of varying *B*-field properties along the ISF, and a more comprehensive investigation is expected to provide new insights into the role of *B*-fields in filament dynamics and star formation. In this current work, as part of the *B*-fields In Star-forming Region Observations (BISTRO; D. Ward-Thompson et al. 2017; P. Bastien 2020), we use the James Clerk Maxwell Telescope (JCMT) to make submillimeter polarimetric observations of the ISF. The BISTRO team has previously observed the ISF (K. Pattle et al. 2017; D. Ward-Thompson et al. 2017). However, those observations were focused only on OMC-1. In this Letter, we have more than doubled the area studied to also include OMC-2 and 3. The aim is to set those earlier observations in the context of their environment and to understand the bigger picture of the role of magnetic fields in Orion A.

## 2. Observations

The observations of polarized dust emission (project ID: M17BL011, M20AL018) covering OMC-1, 2, and 3 in the Orion ISF were performed using POL-2 (P. Friberg et al. 2016) together with SCUBA-2 (W. S. Holland et al. 2013) on the JCMT. Some observations (project ID: M15BEC02) toward OMC-1 South were taken during the POL-2 commissioning stage. All the data were obtained using the POL-2 DAISY mode (P. Friberg et al. 2016).

The reduction of raw data involves three primary steps and uses two packages, SMURF and KAPPA (T. Jenness et al. 2013; M. J. Currie & D. S. Berry 2014), in the Starlink package (M. J. Currie et al. 2014). With an effective beam size of

14.''1 ( $\sim 0.027$  pc at 393 pc) at  $850 \mu\text{m}$  (J. T. Dempsey et al. 2013), we produced a synthesized map of Stokes parameters using a pixel size of 4''. We perform the absolute flux calibration with the flux conversion factor (FCF) estimated by adopting different recommended FCF values (S. Mairs et al. 2021) weighted with the observation time. In our work, FCFs were set to  $695 \text{ Jy beam}^{-1} \text{ pW}^{-1}$  for OMC-1 and  $668 \text{ Jy beam}^{-1} \text{ pW}^{-1}$  for OMC-2 and 3.

By using three models including background, source, and residual components, we smoothed the polarization maps to revise some obviously inaccurate measurements due to the uncertainty. With the calibrated Stokes parameters, the polarized intensities (PIs), polarization degrees ( $P$ ), and angles ( $\theta_P$ ) at different positions can be calculated using the following equations:

$$\text{PI} = \sqrt{Q^2 + U^2}, \quad P = \frac{\text{PI}}{I}, \quad \text{and} \quad \theta_P = 0.5 \tan^{-1} \left( \frac{U}{Q} \right). \quad (1)$$

Since both positive and negative  $Q$  and  $U$  values contribute to a positive PI value, a modified asymptotic estimator (S. Plaszczynski et al. 2014) is employed to debias the results to avoid the overestimation of PI.

Using these equations, the parameters and their uncertainties are determined to produce a catalog of polarization half-vectors. The polarization vectors with  $P/\sigma_P < 3$  or  $\sigma_P > 5\%$  are removed in our analysis, where  $\sigma_P$  is the uncertainty of polarization degree. All the selected polarization vectors, with their lengths proportional to the polarization degrees and plotted in an interval of  $8''$ , are shown in Figure 1(a).

### 3. Results

#### 3.1. Magnetic Field Morphology

Assuming aligned dust grains regulated by  $B$ -fields based on the radiation alignment theory (A. Lazarian 2007), the polarization angles of thermal dust emission enable one to infer the orientation of the  $B$ -field projected on the plane of the sky (PoS). Figures 1(b)–(d) show the statistics of  $B$ -field orientation distributions of the three clouds. It is clear that the  $B$ -fields in OMC-1 are mostly aligned along a northwest–southeast orientation with a position angle (PA) of about  $120^\circ$ , while the  $B$ -fields in OMC-3 are predominantly aligned along a northeast–southwest orientation with a PA of about  $45^\circ$ . On the other hand, the  $B$ -field orientations in OMC-2 have a broad distribution between  $50^\circ$  and  $130^\circ$  and another group between  $0^\circ$  and  $30^\circ$ , indicating a relatively more random distribution. In Figure 1(e), we present the half-vectors rotated by  $90^\circ$  in an interval of  $20''$  representing the corresponding  $B$ -field orientations across the filament.

*OMC-1.* Overall, the  $B$ -field appears to be perpendicular to the main axis of the cloud/filament. OMC-1 is associated with the Orion Nebula Cluster (ONC) and contains a high concentration of gas at a high temperature of  $>100$  K (e.g., D. Li et al. 2020). The maximum  $850 \mu\text{m}$  brightness of OMC-1 is approximately  $9 \times 10^5 \text{ mJy beam}^{-1}$  and is associated with the hot, high-mass star-forming clumps of Orion BN/KL. Moreover, as one approaches the location of Orion BN/KL, the hourglass pattern of the  $B$ -field becomes more prominent. This pinched morphology in the central cloud of OMC-1 indicates a strong interaction between gravity and the  $B$ -field, showcasing the effects of the  $B$ -field in high-mass star-forming regions. Figure 1 also shows  $B$ -field lines that are aligned parallel with the orientation of the cloud extension in the northeastern

subfilament of OMC-1, suggesting a gas accumulation process that is guided by the  $B$ -field surrounding the filament.

*OMC-2.* Our dust polarization map of OMC-2 is a marked improvement over previous observations, such as the SCUPOL results (F. Poidevin et al. 2010). As a site for intermediate-to-low-mass star formation, OMC-2 seems to have relatively chaotic  $B$ -field structures compared to the other two clouds. From Figure 1(e), the  $B$ -field lines seem to converge toward denser areas in the central parts of OMC-2, where gravitational contraction is likely taking place. In contrast, a subfilament to the west of the main filament is overall perpendicular to the  $B$ -field.

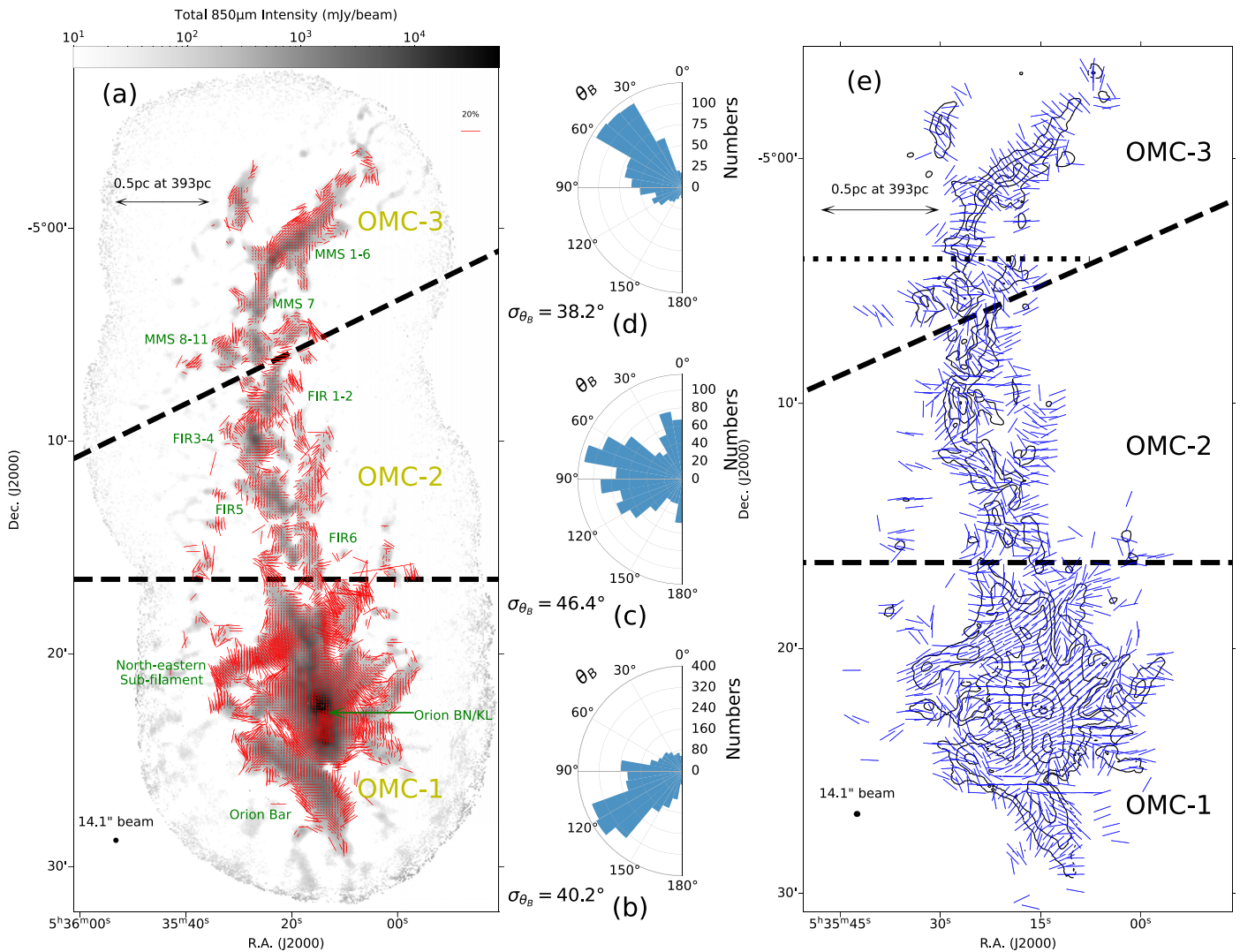
*OMC-3.* OMC-3 appears to be a filamentary cloud extending from southeast to northwest. The POL-2 observations reveal a nearly uniform  $B$ -field in the northern backbone of OMC-3. However, the OMC-3 South, which is suspected to be a ‘‘second filament’’ as noted by F. Poidevin et al. (2010), has disorderly  $B$ -field directions that are similar to the complicated structures of OMC-2. In the main body of OMC-3 north, the  $B$ -field orientations are orthogonal to the filament direction. In brief, OMC-3 exhibits very ordered to even uniform  $B$ -field structures.

#### 3.2. Multiscale View of B-field Geometries in Orion ISF

We utilized the 353 GHz polarization observations made with the High Frequency Instrument on Planck to infer the large-scale  $B$ -field (Planck Collaboration et al. 2015). The Stokes  $I$ ,  $Q$ , and  $U$  maps, which were corrected for the contamination from the cosmic microwave background and cosmic infrared background, were used to generate the large-scale polarization map at a resolution of  $5'$ . Figure 2(a) displays the large-scale  $B$ -field maps around the Orion A region. The  $B$ -field is roughly perpendicular to the ISF. Moreover, the field structure appears slightly pinched toward the filament.

In addition, we check the optical starlight polarization observations to further explore the large-scale  $B$ -field in relatively low-density regions (F. Poidevin et al. 2011). To limit our analysis to sources within the Orion cloud, we only consider starlight detections that fall within the region of our JCMT observations and have a distance of 360–500 pc based on Gaia parallax measurements (S. Rezaei et al. 2020; C. A. L. Bailer-Jones et al. 2021; Gaia Collaboration et al. 2021). The  $B$ -fields derived from 61 detections are shown in Figure 2(b). The majority of the  $B$ -field half-vectors have a west–east or northeast–southwest orientation, in general consistent with the Planck results. Given that the optical polarization data are presumably tracing the  $B$ -field threading the ISM around the ISF or that in the foreground toward the ISF, the  $B$ -field structure shows a small deviation compared to that seen in the Planck map.

The TADPOL survey (C. L. H. Hull et al. 2014) mapped the  $B$ -fields toward several selected sources in the ISF, including Orion KL in OMC-1, FIR 3 and FIR 4 in OMC-2, and MMS 5 and MMS 6 in OMC-3, at an angular resolution of  $2''.5$  ( $0.005$  pc) using the Combined Array for Research in Millimeter-wave Astronomy (CARMA). From Figures 2(c)–(e), for OMC-1, the orientation of the small-scale  $B$ -field revealed by CARMA largely follows that of the intermediate-scale  $B$ -field seen by the JCMT POL-2, though a small fraction of the CARMA  $B$ -field half-vectors are offset from parallel to even perpendicular to the JCMT  $B$ -field half-vectors; for OMC-2, the small-scale  $B$ -field is apparently decoupled from that on the intermediate scale, and there is no obvious correlation between the orientations of the  $B$ -fields on the two scales; for OMC-3, the small-scale and intermediate-scale  $B$ -fields both



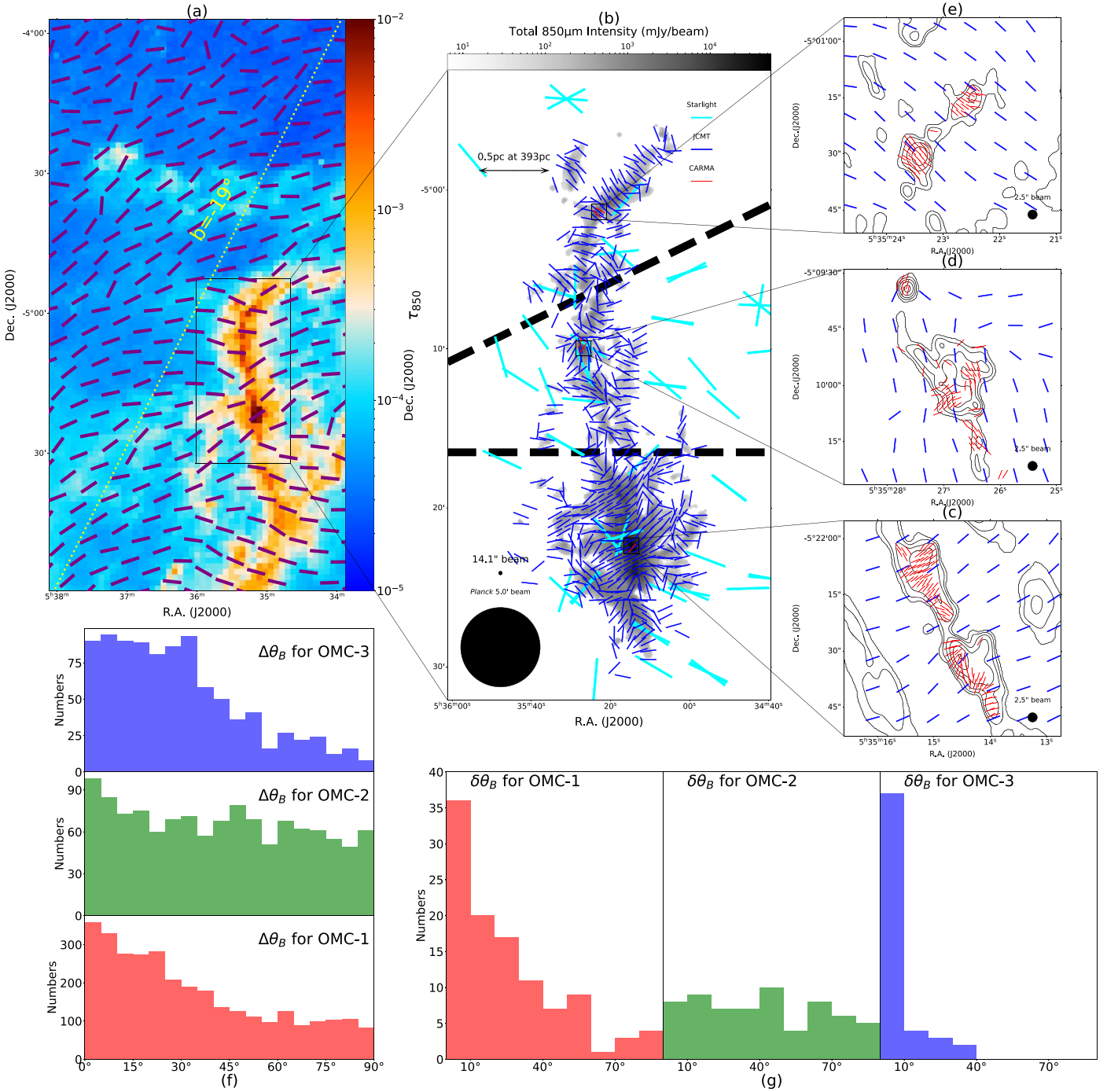
**Figure 1.** Panel (a): dust polarization observations of Orion A ISF made with the POL-2 on JCMT. The gray-scale image shows the  $850\ \mu\text{m}$  total intensity (Stokes  $I$ ). Red vectors are plotted in an interval of  $8''$ , showing the polarization angles with the length proportional to the polarization degree. Black dashed lines mark the divisions between three clouds, i.e., OMC-1 to the south, OMC-2 in the middle, and OMC-3 to the north. The locations of the 1.3 mm sources identified by R. Chini et al. (1997), including MMS 1–10 in OMC-3, FIR 1–6 in OMC-2, and the northeastern subfilament and Orion Bar in OMC-1, are marked on the image. Panels (b), (c), and (d) show the histograms of the PAs of the  $B$ -field orientations for OMC-1, OMC-2, and OMC-3, respectively. Panel (e): blue vectors with a uniform arbitrary length are plotted in an interval of  $20''$ , showing the magnetic field orientations, and are derived by rotating the polarization vectors by  $90^\circ$ . The  $850\ \mu\text{m}$  total intensity is shown in black contours at  $\log_{10}$  scale ( $\text{mJy beam}^{-1}$ ), which starts from 2.2 and continues in steps of 0.5. The black dotted line splits the OMC-3 cloud into the North and South parts.

appear to be uniform with almost the same orientation. There have been new Atacama Large Millimeter/submillimeter Array (ALMA) observations of dust polarization toward several sources in the ISF; however, these observations were made at subarcsecond resolutions, either probing  $B$ -field structures at too-small scales to be compared with the JCMT data (P. C. Cortes et al. 2021) or being dominated by self-scattering and thus unable to probe the  $B$ -field structure (S. Takahashi et al. 2019; Y. Liu et al. 2024).

More quantitatively, we compare the intermediate-scale  $B$ -field probed by JCMT POL-2 with the large-scale  $B$ -field probed by Planck by calculating the difference angle,  $\Delta\theta_B$ , between the orientations of the  $B$ -fields on the two scales. Since the Planck map covers an area much larger than what is covered by the JCMT map, we calculate  $\Delta\theta_B$  for each half-vector at  $8''$  intervals in the JCMT map; the  $B$ -field orientation at the corresponding position in the Planck map is derived by a weighted average of the  $B$ -field orientations at the nearest 4

pixels, where the pixel size of the Planck map is  $2'$  and the weighting is taken as the inverse of the square of the distance between the pixel center to the position of interest. In Figure 2(f), the histogram of  $\Delta\theta_B$  for OMC-1 is clearly peaking toward  $0^\circ$ , suggesting that the orientation of the intermediate-scale ( $0.03\ \text{pc}$ )  $B$ -field is predominantly parallel with that of the large-scale ( $0.6\ \text{pc}$ )  $B$ -field. For OMC-2,  $\Delta\theta_B$  appears to be widely distributed between  $0^\circ$  and  $90^\circ$ , with a very minor tendency of peaking at  $0^\circ$ , indicating that the  $B$ -field orientation on the intermediate scale has shown strong local variation and started to decouple from that on the large scale. For OMC-3,  $\Delta\theta_B$  are almost all below  $35^\circ$ , indicating that the intermediate-scale  $B$ -field is well aligned with that on the large scale.

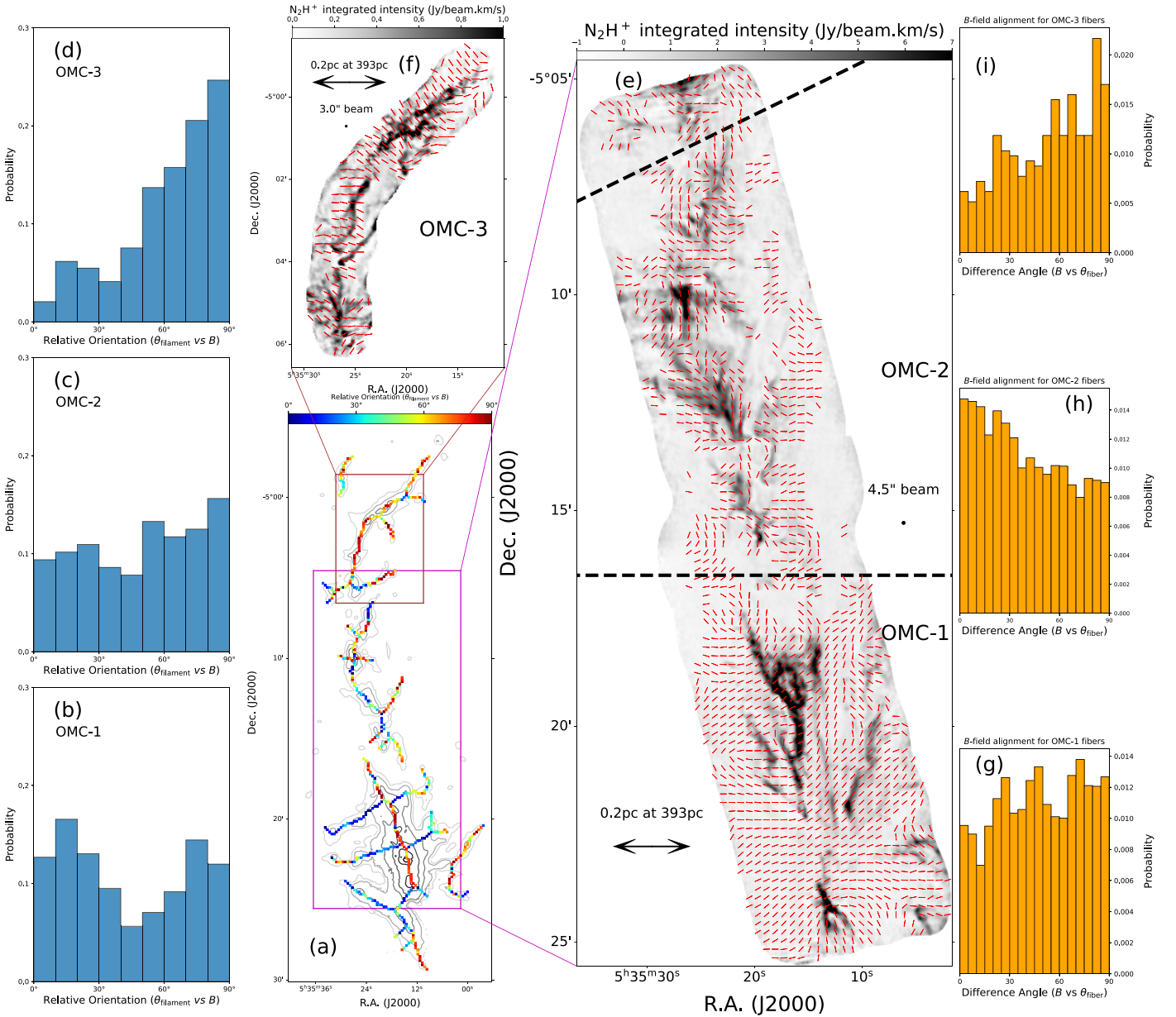
Similarly to  $\Delta\theta_B$ , we compute the difference angle,  $\delta\theta_B$ , between the orientations of the  $B$ -fields probed by JCMT and CARMA for each CARMA detection. Again, the  $B$ -field orientation at the corresponding position in the JCMT map is



**Figure 2.** Multiscale  $B$ -field orientations in the ISF. Panel (a): the background image displays the 850  $\mu\text{m}$  opacity map obtained from the Herschel and Planck data (M. Lombardi et al. 2014); purple segments indicate the large-scale  $B$ -field orientations inferred from the Planck 353 GHz data. A yellow dotted line marks the galactic latitude  $b = 19^\circ$ . Panel (b):  $B$ -field orientations derived from starlight, JCMT POL-2, and CARMA observations. The background image shows the JCMT 850  $\mu\text{m}$  total intensity map; blue vectors plotted at an interval of  $32''$  denote the  $B$ -field orientations observed by JCMT POL-2, cyan vectors represent the  $B$ -field orientations revealed by starlight polarization observations (F. Poidevin et al. 2011), and red vectors show the averaged  $B$ -field orientations obtained by the CARMA TADPOL survey (C. L. H. Hull et al. 2014). In panels (c), (d), and (e), red vectors indicate the  $B$ -field orientations obtained by the CARMA TADPOL survey in selected dense cores located in OMC-1, OMC-2, and OMC-3, respectively; blue vectors show the  $B$ -field orientations derived with the JCMT POL-2 observations; and the CARMA observations of the total dust emission at 1.3 mm are shown in black contours at  $\log_{10}$  scale ( $\text{mJy beam}^{-1}$ ), which starts from  $-2.0$  and continues in steps of 0.5 in panel (c), starts from  $-2.0$  and continues in steps of 0.2 in panel (d), and starts from  $-1.0$  and continues in steps of 0.3 in panel (e). Panel (f): histograms of the difference angles between the  $B$ -field orientations unveiled by Planck and that by the JCMT POL-2 for OMC-1, OMC-2, and OMC-3. Panel (g): histograms of the difference angles between the  $B$ -field orientations obtained by the JCMT POL-2 and by the CARMA TADPOL survey for OMC-1, OMC-2, and OMC-3.

derived by a weighted average of the  $B$ -field orientations at the nearest 4 pixels, and the weighting is taken as the inverse of the square of the distance between the pixel center to the position of interest. In Figure 2(g), the distribution of  $\delta\theta_B$  for OMC-1 shows a clear peak at  $0^\circ$ – $10^\circ$  and gradually declines

toward  $90^\circ$ ; for OMC-2,  $\delta\theta_B$  has a nearly flat distribution, again indicating that the  $B$ -fields on the two scales are apparently decoupled; for OMC-3, the distribution of  $\delta\theta_B$  depicts that the  $B$ -field orientations almost do not change across the two scales.



**Figure 3.** Panel (a): colored points show locations along the skeleton of each filament where the filament orientation is compared with the  $B$ -field orientation, and the color scale denotes the derived difference angle between the two orientations, as indicated by a color bar on the top; the  $850\ \mu\text{m}$  total intensity is shown in contours with levels at a  $\log_{10}$  scale ( $\text{mJy beam}^{-1}$ ), starting from 2.2 and continuing at steps of 0.5. Panels (b), (c), and (d): histograms of the difference angles between the filament skeleton and  $B$ -field orientations for OMC-1, OMC-2, and OMC-3. Panel (e): red vectors plotted at an interval of  $12''$  show the  $B$ -field orientations derived by the JCMT POL-2, and the gray-scale image shows the  $\text{N}_2\text{H}^+$  (1–0) velocity-integrated emission in OMC-1 and OMC-2 (A. Hacar et al. 2018). Panel (f): same as panel (e) but for OMC-3, and the  $\text{N}_2\text{H}^+$  data are taken from C. Zhang et al. (2020). Panels (g), (h), and (i) show the histograms of the difference angles between the  $\text{N}_2\text{H}^+$  fibers and POL-2  $B$ -field orientations in OMC-1, OMC-2, and OMC-3, respectively.

### 3.3. Relations between $B$ -field and Filamentary Structures

To investigate how the  $B$ -field orientation is aligned with the filamentary structures in the ISF, we employed the *filfinder* algorithm (E. W. Koch & E. W. Rosolowsky 2015) to extract filament skeletons. Figure 3(a) shows the derived skeletons along the main filament, the branches connected to the main filament, and some minor structures detached from the main filament. To quantify the filament orientations, we utilize the principal component analysis method on 10 adjacent pixels of the skeletons to determine the PA of the filaments at each position. We then compute the difference angles between the filament and  $B$ -field orientations. In Figure 3(a), the color scale of the skeletons visualizes the spatial distribution of the

difference angles. Figures 3((b)–(d)) show the histograms of the difference angles for OMC-1, OMC-2, and OMC-3, respectively. Along the filamentary cloud, three drastically different distributions for the relative orientation between the  $B$ -fields and filaments are seen: a bimodal distribution for OMC-1, nearly flat distribution for OMC-2, and a distribution with a predominant single peak at  $90^\circ$  for OMC-3. From the skeleton color scale representing the difference angles (Figure 3(a)), we can see that the bimodal distribution in OMC-1 is due to a combined effect that along the main filament, the  $B$ -field orientation is perpendicular to the filament axis, while along the relatively low-density branches, the  $B$ -field orientation is parallel to the branch axis; on the other hand,

for the distribution of relative orientation in OMC-3, a tail toward  $0^\circ$  is mostly attributed to OMC-3 South.

Molecular filaments may have complex internal structures, such as intertwined filamentary bundles or fibers. We identified the fibers with the filfinder algorithm from  $\text{N}_2\text{H}^+$  maps (see Appendix A). Figure 3(e) shows a comparison between the  $B$ -field orientations derived from our POL-2 observations and the fiber structures revealed by the combined ALMA and IRAM 30 m  $\text{N}_2\text{H}^+$  (1–0) observations of OMC-1 and 2 (A. Hacar et al. 2018). Such a comparison for OMC-3 is shown in Figure 3(f), where the ALMA  $\text{N}_2\text{H}^+$  (1–0) data were taken from C. Zhang et al. (2020). We calculate the difference angles between the fiber and  $B$ -field orientations, as shown in Figures 3(g)–(i). In OMC-1, the fibers tend to be perpendicular to the  $B$ -field; this is not difficult to understand considering the bimodal distribution for the relative orientation between the  $B$ -field and filaments (Figure 3(b)), and here the fibers traced by the  $\text{N}_2\text{H}^+$  emission represent the high-density part of the filaments. In OMC-3, the fibers are clearly perpendicular to the  $B$ -field, consistent with the distribution of relative orientation between the  $B$ -field and filaments. Interestingly, the fibers in OMC-2 appear to be predominantly parallel to the  $B$ -field, in contrast to the random distribution of relative orientation between the  $B$ -field and filaments.

## 4. Discussion and Summary

### 4.1. A Tale of Three: Gravitational, Turbulent, and Magnetic Interpretations for OMC-1, 2, and 3, Respectively

We have presented JCMT POL-2 dust polarization observations of a remarkable molecular filament containing OMC-1, OMC-2, and OMC-3 in the Orion ISF. Combining the POL-2 data with the Planck and CARMA polarization observations, we clearly see how the  $B$ -fields vary from the large ( $\sim 0.6$  pc) to intermediate ( $\sim 0.03$  pc) and small ( $\sim 0.005$ ) scales: for OMC-1, the  $B$ -field retains its mean orientation on all the scales, with some local variations on intermediate-to-small scales; for OMC-2, the  $B$ -fields on different scales are apparently decoupled, showing relatively disordered morphology on the intermediate and small scales; and for OMC-3, the  $B$ -field shows a uniform morphology, and the orientation does not change all the way from the large to intermediate and small scales. A natural and straightforward interpretation of such  $B$ -field morphologies, in particular their variation across different scales, is that the  $B$ -field in OMC-1 is channeling the gas accretion from the ambient medium to the filament, but as the mass continues to grow, forming massive dense cores within the filament, gravity overcomes the magnetic force, pulling the  $B$ -field into an hourglass shape (see also K. Pattle et al. 2017; D. Ward-Thompson et al. 2017). The  $B$ -field in OMC-2 appears highly twisted on intermediate and small scales, suggesting that turbulence is dominating over the  $B$ -field; the  $B$ -field in OMC-3, especially OMC-3 North, has a nearly uniform morphology from large to small scales, indicating that the  $B$ -field is strong enough to dominate the gas dynamics (e.g., E. C. Ostriker et al. 2001). Below, we test this simple interpretation by comparing the orientations between the  $B$ -fields and the dense gas structures.

Filamentary clouds naturally define an axis to be compared to the  $B$ -field, and such a comparison for ISF again reveals a trio: bimodal for OMC-1, random for OMC-2, and perpendicular for OMC-3 (Figures 3(b)–(d)). It immediately renders strong support to the above ternary interpretation. The

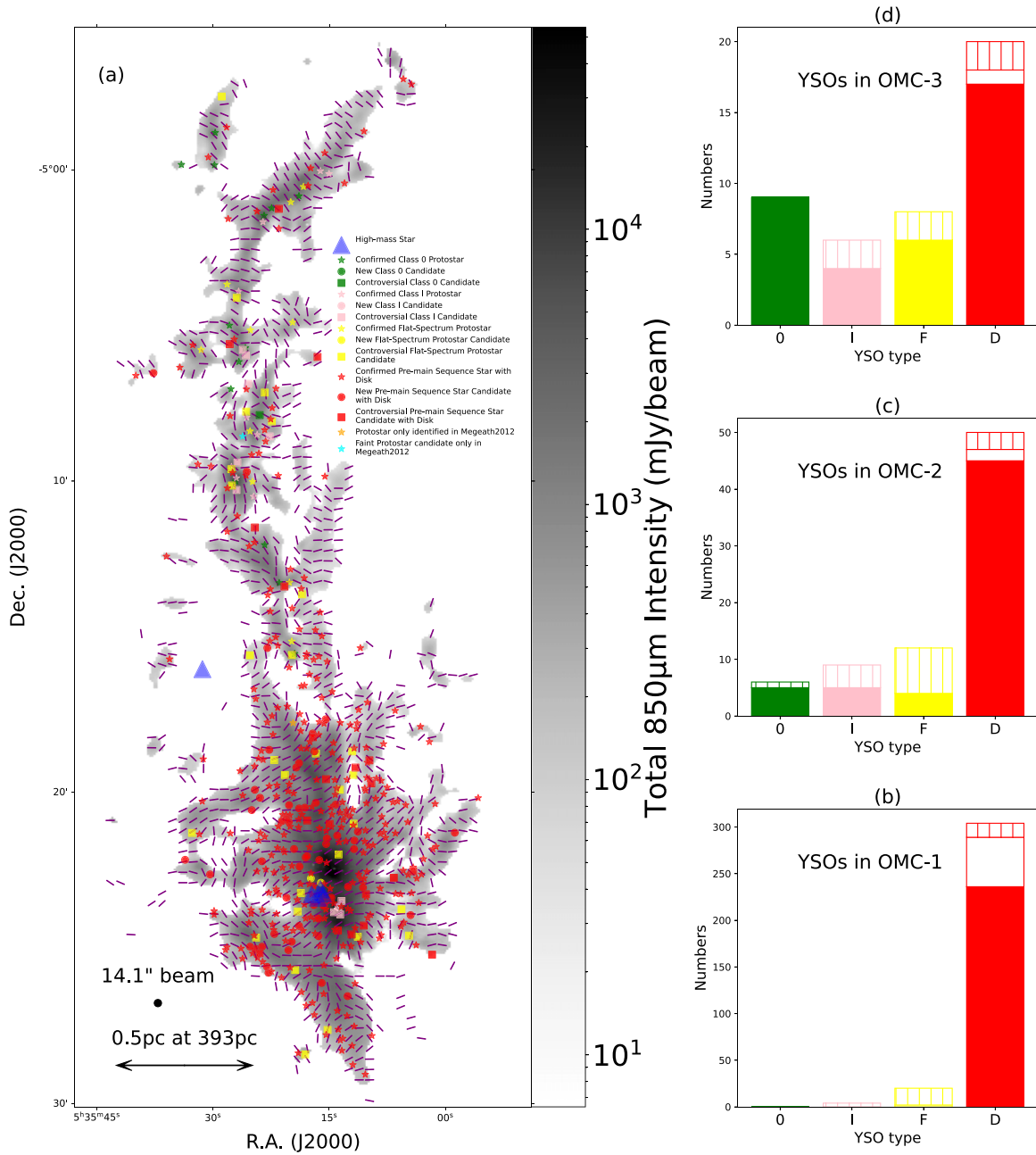
bimodal distribution for the relative orientation between the  $B$ -field and filaments in OMC-1 is clearly correlated to the gas density, with the high-density filament skeleton perpendicular to the  $B$ -field and low-density skeletons parallel to the  $B$ -field, consistent with the scenario that the  $B$ -field is channeling gas flows toward the high-density filament (e.g., D. Ward-Thompson et al. 2017; T. G. S. Pillai et al. 2020; P. Girichidis 2021). Such a correlation is strengthened by looking into the filament internal structures, i.e., the  $\text{N}_2\text{H}^+$  fibers: as the high-density part of the filament, the fibers are preferentially perpendicular to the  $B$ -field (Figure 3(g)). In OMC-2, the random distribution is apparently a consequence of the disordered nature of the  $B$ -field structure. Very interestingly, the fibers in OMC-2 are largely parallel to the  $B$ -field (Figure 3(h)), showing a pattern that is consistent with the results of simulations of super-Alfvénic turbulence (see, e.g., Figures 2 and 3 in P. Padoan et al. 2001), suggesting that turbulence is dynamically more important than the  $B$ -field in OMC-2. For OMC-3, the  $B$ -field is simply perpendicular to both the filament (Figure 3(d)) and fibers (Figure 3(i)), indicating that the  $B$ -field is strong enough to counteract gravity and turbulence.

### 4.2. The $B$ -field Strength Estimates

To further quantify the impact of the  $B$ -field on the dynamical evolution of the filament, it is desirable to estimate the  $B$ -field strength. However, deriving the  $B$ -field strength with the David–Chandrasekhar–Fermi (DCF) method (L. Davis 1951; S. Chandrasekhar & E. Fermi 1953) or its variants is subject to large uncertainty and in some cases is not applicable (J. Liu et al. 2021, 2022a, 2022b; C.-Y. Chen et al. 2022). First, the method requires calculating the polarization angle dispersion due to turbulent disturbance, or decomposing the  $B$ -field into turbulent and ordered components and calculating their ratio. This step is not always feasible, especially when the  $B$ -field structure is complicated. Second, under the assumption of energy equipartition between turbulence and the turbulent  $B$ -field, and adopting a gas density and turbulent velocity dispersion obtained from other observations, the PoS  $B$ -field strength can be derived. It should be noted that the energy equipartition assumption may not be valid when the  $B$ -field is weak. The estimates of the gas density and turbulent velocity dispersion often suffer large uncertainties. Nevertheless, the method has been widely used. Several such estimates for the sources in the ISF exist in the literature, and the results vary a lot, ranging from 0.3 to 6.6 mG for OMC-1 and from 0.13 to 0.64 mG for OMC-3 (B. C. Matthews et al. 2005; J. P. Vallée & J. D. Fiege 2007; R. H. Hildebrand et al. 2009; M. Houde et al. 2009; F. Poidevin et al. 2013; K. Pattle et al. 2017; D. T. Chuss et al. 2019; J. A. Guerra et al. 2021; J. Hwang et al. 2021; P. S. Li et al. 2022; N. Zielinski & S. Wolf 2022). Here we try with the best effort to estimate the PoS  $B$ -field strengths in the three regions with the new data, obtaining 0.45, 0.25, and 0.37 mG for OMC-1, 2, and 3, respectively (see Appendix B).

Given the aforementioned cautions and uncertainties, we only make a comparative rather than a more detailed quantitative analysis based on the derived  $B$ -field strengths. It is worth mentioning that OMC-1 has a width about 2 times greater than OMC-2 and 3 under the same column density threshold, resulting in a volume density in this region slightly lower than that in the latter two. But about half of the gas mass





**Figure 4.** Panel (a): distribution of detected YSOs overlaid on the  $B$ -field orientation maps. All the YSO candidates and high-mass stars are taken from the literature (S. T. Megeath et al. 2012; E. Furlan et al. 2016; J. E. Großschedl et al. 2019). Purple vectors at a  $20''$  interval indicate the  $B$ -field orientations observed by the JCMT POL-2. The class 0, class I, flat-spectrum sources, and pre-main-sequence stars with disks are denoted in green, pink, yellow, and red colors, respectively. For each YSO type, the confirmed ones that are consistent in different literature are represented by star symbols; newly discovered candidates by J. E. Großschedl et al. (2019) are indicated with filled circles; controversial candidates, showing inconsistencies in different literature, are marked with filled squares. Panels (b), (c), and (d): histograms of the four YSO types in OMC-1, OMC-2, and OMC-3, respectively; class 0, class I, flat-spectrum sources, and pre-main-sequence stars with disks are labeled as “0,” “I,” “F,” and “D,” respectively, with the same colors as in panel (a); filled histograms represent confirmed YSOs, open histograms depict newly discovered YSOs, and vertical gridded histograms indicate controversial YSOs.

in OMC-1 is attributed to the central high-density part with a width of  $\sim 0.06$  pc, and within that area, the average volume density reaches  $2.5 \times 10^6 \text{ cm}^{-3}$ . Therefore, though the  $B$ -field in OMC-1 is stronger, considering a much greater mass and central density, it is completely plausible that gravity is overwhelmingly more important in this region. OMC-2 and 3 have almost the same mass and nonthermal velocity dispersion (see Appendix B), while the  $B$ -field in OMC-3 is stronger than that in OMC-2; the relative  $B$ -field strength of the two regions is at least compatible with the interpretation that turbulence in OMC-2 and the  $B$ -field in OMC-3 are taking a leading role.

### 4.3. The Impact of the Magnetic Field on Star Formation

Given the markedly different  $B$ -field properties across the three regions in the ISF, it is of great interest to examine how the star formation activity is affected. We collect a catalog of young stellar objects (YSOs), which are classified into class 0, class I, flat-spectrum, and disk-bearing pre-main-sequence stars, based on the works of S. T. Megeath et al. (2012), E. Furlan et al. (2016), and J. E. Großschedl et al. (2019). Figure 4 shows all the YSOs in OMC-1, 2, and 3 and the statistics of each type in each of the three regions. The star

formation activity in OMC-1 is far more vigorous and complicated than that in OMC-2 and 3. OMC-1 is the only region of the three forming high-mass stars, containing several well-known high-mass protostellar objects. It is located behind the luminous Trapezium cluster, which is the central part of the ONC. The collected YSOs in this region are completely dominated by the disk sources and are heavily contaminated by the foreground ONC sources (C. J. Lada et al. 2000; J. Otter et al. 2021). Here we focus on the comparison between OMC-2 and OMC-3. From Figure 4, OMC-2 has a higher fraction of disk sources (50/78) than OMC-3 (20/43), indicating a younger age of the cluster in OMC-3. The total number of YSOs in OMC-2 is higher than that in OMC-3. Note that the mass, mean density, and nonthermal velocity dispersion in the two regions are almost the same, and the only appreciable difference lies in the  $B$ -field geometries and the relative orientation between the  $B$ -fields and filaments/fibers. Therefore, the differing YSO populations in the two regions are mostly likely due to the  $B$ -field effect, providing compelling evidence that a dynamically more important  $B$ -field leads to slower (or delayed) and less efficient star formation in molecular clouds.

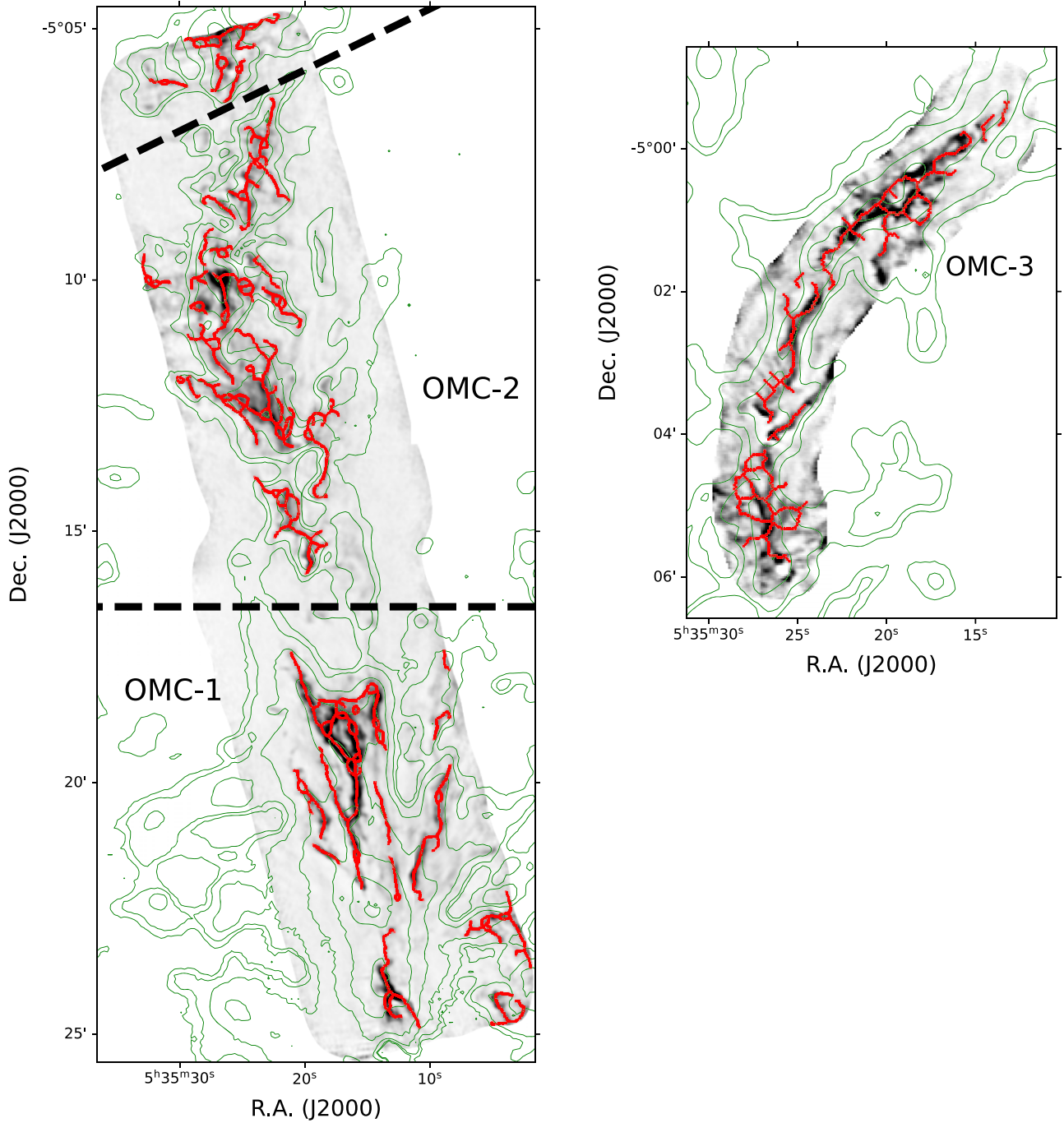
To summarize, concerning which mechanism is shaping the dynamics of molecular clouds on  $\sim 0.01$ – $1$  pc scales, each of the three clouds (OMC-1, 2, and 3 in the Orion ISF) seems to be telling a different story based on our JCMT POL-2 observations along with the Planck and CARMA data. Therefore, it is probably an oversimplified interpretation to claim that either the magnetic field or turbulence is universally more important in molecular cloud evolution and star formation. By comparing the YSO populations in OMC-2 and 3, we find evidence that a strong  $B$ -field could make star formation relatively slower and less efficient. Y. Zhang et al. (2019) carried out MHD simulations of sub-Alfvénic molecular clouds, focusing on the  $B$ -field orientation variation across various scales. They found that on small ( $< 0.1$  pc) scales, the cores are super-Alfvénic as a consequence of turbulent energy concentration induced by gravity, and thus the  $B$ -field on small scales exhibits a wide range of deviation in orientation from that on large scales. If one takes an average  $B$ -field orientation for each of the dense cores in the CARMA maps (Figures 2(c)–(e)) and compares to the  $B$ -field revealed by Planck, the offset distribution could be to some extent consistent with the work of Y. Zhang et al. (2019). However, a detailed comparison shows that the cross-scale correlation in  $B$ -field orientation (Figures 2(f) and (g)) is distinctly different from region to region, certainly not random in OMC-1 and 3. The observed relation between the  $B$ -field and filament/fiber orientations and the star formation activity variation further suggests a tale-of-three interpretation of the three regions regarding the interplay between gravity,  $B$ -fields, and turbulence.

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## Appendix A Identified Fiber Structures

The fibers within the filament are extracted using the filfinder algorithm from the  $N_2H^+$  (1–0) velocity-integrated emission maps. Figure 5 shows a comparison between the derived fibers, the  $N_2H^+$  emission, and the total 850  $\mu\text{m}$  emission.



**Figure 5.** Gray-scale images show the velocity-integrated  $\text{N}_2\text{H}^+$  (1–0) emissions, overlaid with the extracted fibers shown in red lines and the total  $850\ \mu\text{m}$  intensity shown in green contours. The left panel shows the OMC-1 and 2 region, with the  $\text{N}_2\text{H}^+$  data taken from A. Hacar et al. (2018), and the right panel shows OMC-3, with the  $\text{N}_2\text{H}^+$  data taken from C. Zhang et al. (2020).

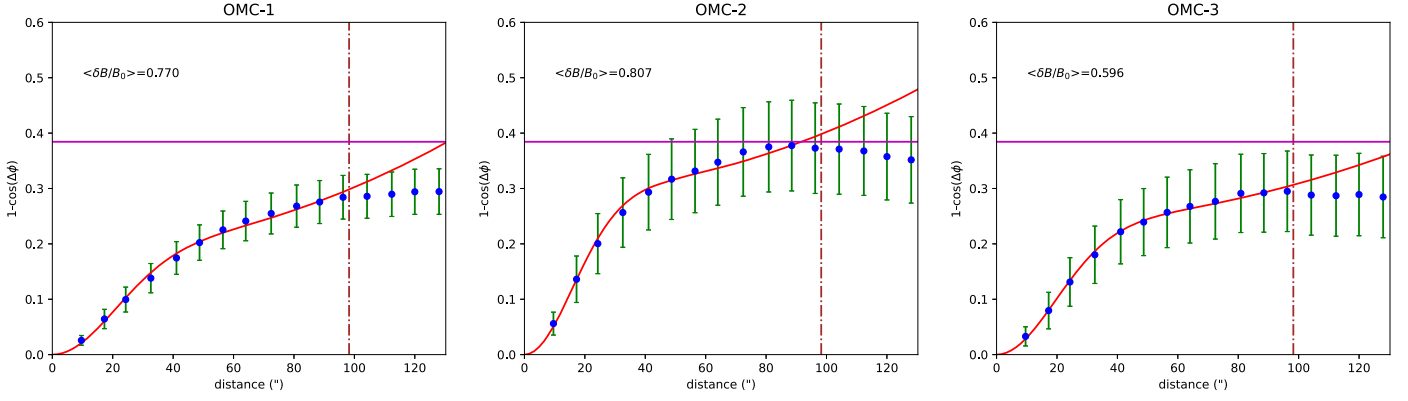
### Appendix B Details for $B$ -field Strength Calculation

In the DCF assumption, the PoS  $B$ -field strength of the molecular cloud is estimated by interpreting the observed deviation of polarization angles from a mean polarization angle distribution as a result of Alfvén waves induced by turbulent perturbations, i.e.,

$$B_0 = \sigma_v \sqrt{\mu_0 \rho} \left( \frac{\delta B}{B_0} \right)^{-1}, \quad (\text{B1})$$

where  $\sigma_v$  represents the turbulence-induced velocity dispersion that could approximately equate to the nonthermal velocity dispersion and  $\rho$  denotes the gas mass density.  $\delta B/B_0$  denotes the turbulent-to-ordered magnetic field ratio.

To obtain the mass density, we modeled the three star-forming clouds within the Orion A ISF as cylindrical filaments. We use the column density map at  $\sim 8''$  resolution produced by F. Schuller et al. (2021) to measure the mass of the three regions, obtaining  $\sim 660$ ,  $\sim 250$ , and  $\sim 260 M_\odot$  for OMC-1, 2, and 3. The dimensions of the three regions are measured to be approximately  $0.93\ \text{pc} \times 0.23\ \text{pc}$  for OMC-1,  $1.0\ \text{pc} \times 0.1\ \text{pc}$



**Figure 6.** Fitting results of the ACF for OMC-1, 2, and 3 from left to right. For each panel, blue filled circles with error bars denote data points derived from the polarization observations, the best-fitting result is shown by a red solid line, a horizontal magenta line marks the value expected for a random field ( $52^\circ$ ; F. Poidevin et al. 2010), and a vertical brown dashed–dotted line marks the right boundary of the points to be fitted.

for OMC-2, and  $1.0 \text{ pc} \times 0.1 \text{ pc}$  for OMC-3. Assuming a cylinder geometry lying in the PoS, the volume densities are found to be  $\sim 2.4 \times 10^5$ ,  $\sim 4.5 \times 10^5$ , and  $\sim 4.7 \times 10^5 \text{ cm}^{-3}$  for OMC-1, 2, and 3, respectively.

To estimate the velocity dispersion in the ISF, we utilized the  $\text{NH}_3$  (1, 1) observation data from the Green Bank Ammonia Survey (R. K. Friesen et al. 2017) with a resolution of  $36''$ . To extract the nonthermal velocity dispersion, we subtracted the thermal components of the observed velocity dispersion with the temperature map provided by F. Schuller et al. (2021). Our analysis revealed that the mean nonthermal velocity dispersion in OMC-1, OMC-2, and OMC-3 is  $0.90 \text{ km s}^{-1}$ ,  $0.38 \text{ km s}^{-1}$ , and  $0.41 \text{ km s}^{-1}$ , respectively.

The turbulent-to-ordered magnetic field ratio  $\delta B/B_0$  is determined by the dispersion of polarization angles. However, quantifying the turbulent  $B$ -field components could have bias due to the effects of nonturbulent field structure in dense clouds. So the angular dispersion function method has been developed to reduce the bias. Moreover, by considering the effect of signal integration along the line of sight and within the beam in the analysis, M. Houde et al. (2009) proposed the autocorrelation function (ACF) form to precisely derive the turbulent-to-ordered magnetic field ratios. The angular dispersion function could be expressed as

$$1 - \langle \cos[\Delta\Phi(l)] \rangle \simeq \frac{1}{N} \frac{\langle \delta B^2 \rangle}{\langle B_0^2 \rangle} \times [1 - e^{-l^2/2(\delta^2+2W^2)}] + a_2 l^2, \quad (\text{B2})$$

where  $N$  is the number of turbulent cells probed by the telescope beam,  $\Delta\Phi(l)$  represents the PA differences of two vectors at a distance  $l$ ,  $a_2$  signifies the slope of the second-order term in the Taylor expansion,

$$N = \frac{(\delta^2 + 2W^2)\Delta'}{\sqrt{2\pi}\delta^3}, \quad (\text{B3})$$

$W$  denotes the beam radius ( $6''.0$  for JCMT  $850 \mu\text{m}$  observations),  $\Delta'$  depicts the cloud depth, and  $\delta$  stands for the turbulent correlation length.

Setting cloud depths to  $0.23 \text{ pc}$ ,  $0.1 \text{ pc}$ , and  $0.1 \text{ pc}$  for OMC-1, OMC-2, and OMC-3, respectively, we derived the ACF of the three clouds with the JCMT POL-2 polarization vectors (2952 vectors in OMC-1, 1118 vectors in OMC-2, 890 vectors

in OMC-3). Equation (B2) is valid when  $l$  is not too big compared to a few times of  $W$  (M. Houde et al. 2009). In addition, we have a polarization map with a finite size, and thus the number of polarization detections on which the ACF could be derived at high intensities decreases as  $l$  increases, leading to degrading statistics for the data points on large  $l$ . We therefore limit our fitting to the data points with  $l < 100''$ . In Figure 6, the fitting results revealed that OMC-3 has the smallest  $\delta B/B_0$  value of 0.596, and OMC-2 has the largest  $\delta B/B_0$  value of 0.807, while OMC-1 has a  $\delta B/B_0$  value of 0.770. We also obtain  $\delta = 4.39$ ,  $3.15$ , and  $4.08 \text{ mpc}$  for OMC-1, 2, and 3, respectively, with the fitting. We note that  $\delta$  cannot be resolved with a telescope beam of  $27 \text{ mpc}$  ( $14''$  at a distance of  $393 \text{ pc}$ ). Such an issue occurs in other works applying the ACF fitting to dust polarization data (e.g., M. Houde et al. 2009; K. Qiu et al. 2013). Thus, the inferred  $\delta$  is more like a numerical artifact from the fitting, and the turbulence correlation scale is still to be explored. We finally estimated the strength of the PoS component of the  $B$ -field for OMC-1, OMC-2, and OMC-3 as  $0.45 \text{ mG}$ ,  $0.25 \text{ mG}$ , and  $0.37 \text{ mG}$ , respectively.

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