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Spectral Measures for G_2 II: finite subgroups

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Abstract

Joint spectral measures associated to the rank two Lie group G_2 , including the representation graphs for the irreducible representations of G_2 and its maximal torus, nimrep graphs associated to the G_2 modular invariants have been studied. In this paper we study the joint spectral measures for the McKay graphs (or representation graphs) of finite subgroups of G_2 . Using character theoretic methods we classify all non-conjugate embeddings of each subgroup into the fundamental representation of G_2 and present their McKay graphs, some of which are new.

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1 Introduction

Two fundamental algebraic structures in conformal quantum field theories are modular tensor categories and conformal nets of factors. The representation theory of the conformal

nets of factors in the algebraic formulation of a conformal field theory, is a modular tensor category describing the Verlinde ring of primary fields. The full conformal field theory is described by extensions of the conformal net, studied through subfactors. The study of modular tensor categories and subfactors and their classification leads to the quest to develop new invariants associated to these categories. Among the most basic invariant for a semisimple tensor category is the underlying fusion or Grothendieck ring, whilst for a subfactor basic invariants which are not complete are the Jones index, the principal graph and its dual [13, 7]. At a more sophisticated level for a subfactor is the standard invariant (also known as the λ -lattice [29], or planar algebra [20]). A modular tensor category is described by its braided structure from which one can derive modular data (in particular its S and T matrices). In the cases of module categories for modular tensor categories (or equivalently braided subfactors in the conformal net picture), invariants include its nimrep (Non-negative Integer Matrix REPresentation of the fusion rules of the underlying modular tensor category) and the cell system built on this nimrep. Another invariant, considered in the work presented in this paper, is the spectral measure associated to the nimrep [1].

The representation ring of a finite group has a basis indexed by simple representations with multiplication given by tensor product. Each representation acts by left multiplication and it is natural to ask what the eigenvalues (with multiplicity) and eigenvectors are for this linear operator. For compact groups a Hilbert space structure is needed on the representation ring and instead of eigenvalues with multiplicity one needs to consider the spectral measure.

Spectral measures associated to the compact Lie groups $SU(2)$ and $SU(3)$ and their maximal tori, nimrep graphs associated to the $SU(2)$ and $SU(3)$ modular invariants, and the McKay graphs (or representation graphs) for finite subgroups of $SU(2)$ and $SU(3)$ were studied in [1, 8, 9]. Spectral measures associated to the compact Lie group $Sp(2)$ are studied in [11], whilst spectral measures associated to other compact rank two Lie groups and their maximal tori are studied in [12].

The real rank two Lie group $G_2 = \text{Aut}(\mathbb{O})$ is the automorphism group of the octonions \mathbb{O} . It is a simply connected and compact Lie group of dimension 14 and is the smallest of the exceptional Lie groups. It is isomorphic to the subgroup of $SO(7)$ that fixes any particular vector in its 8-dimensional real spinor representation. Spectral measures for the representation graphs for the irreducible representations of G_2 and its maximal torus \mathbb{T}^2 , and for nimrep graphs associated to the G_2 modular invariants, were studied in [10]. The spectral measure for G_2 itself appears in [32] as the reduced Haar measure, and is useful for the study of how gluons transform under G_2 at finite temperature [35]. These measures are also useful in other applications, e.g. the computation of glueball masses in Hamiltonian lattice gauge theories or in studies of the strong coupling limit of these theories [23].

In this paper we consider the McKay graphs for finite subgroups of G_2 and study their spectral measures. In the context of finite subgroups of a compact group, the spectral measure simply encodes the character table of the group. However in the context of quantum subgroups discussed above, the spectral measure provides the appropriate generalisation of the character table. We trust that framing this work in terms of spectral measures will make the case of quantum groups more accessible for those who are more familiar with the classical group case.

The spectral measure of a self-adjoint operator a in a unital C^* -algebra A with state φ is a compactly supported probability measure ν_a on the spectrum $\sigma(a) \subset \mathbb{R}$ of a , uniquely determined by its moments

$$\varphi(a^m) = \int_{\sigma(a)} x^m d\nu_a(x),$$

for all non-negative integers m . We need to have a concept to encode the corresponding notions for commuting operators. This is familiar in classical and even free probability [15] as joint laws.

The self-adjoint operators we consider here are the adjacency matrices of the McKay graphs for finite subgroups of G_2 . The characters of the rank two Lie group G_2 are functions on \mathbb{T}^2 , and it is convenient to write the spectral measures for these operators as measures ε over the torus \mathbb{T}^2 . However, \mathbb{T}^2 has dimension one greater than $\sigma(a) \subset \mathbb{R}$, so that there is an infinite family of pullback measures ε over \mathbb{T}^2 for any spectral measure ν_a . The details of the relation between the measures ε and ν_a are given in Section 2.1.

In order to remove this ambiguity, we also consider joint spectral measures, that is, measures over the joint spectrum $\sigma(a, b) \subset \sigma(a) \times \sigma(b) \subset \mathbb{R}^2$ of commuting self-adjoint operators a and b . The abelian C^* -algebra B generated by a , b and the identity 1 is isomorphic to $C(X)$, where X is the spectrum of B . Then the joint spectrum is defined as $\sigma(a, b) = \{(a(x), b(x)) \mid x \in X\}$. In fact, one can identify the spectrum X with its image $\sigma(a, b)$ in \mathbb{R}^2 , since the map $x \mapsto (a(x), b(x))$ is continuous and injective, and hence a homomorphism since X is compact [31]. In the case where the operators a , b act on a finite-dimensional Hilbert space, this is the set of all pairs of real numbers (λ_a, λ_b) for which there exists a vector ϕ , $\|\phi\| = 1$, such that $a\phi = \lambda_a\phi$, $b\phi = \lambda_b\phi$. Then there exists a compactly supported probability measure $\tilde{\nu}_{a,b}$ on $\sigma(a, b)$, which is uniquely determined by its cross moments

$$\varphi(a^m b^n) = \int_{\sigma(a,b)} x^m y^n d\tilde{\nu}_{a,b}(x, y), \tag{1}$$

for all non-negative integers m, n . Such joint spectral measures specialize to the spectral measures ν_a (respectively ν_b) by integrating over all y for which $(\lambda_a, y) \in \sigma(a, b)$ (respectively all x for which $(x, \lambda_b) \in \sigma(a, b)$). As discussed in Section 2.1, such a measure uniquely defines a measure over \mathbb{T}^2 invariant under an action of the Weyl group of G_2 . In this paper we determine the joint spectral measure for each finite subgroup Γ of G_2 for all non-conjugate embeddings of Γ into the fundamental representations of G_2 . In the process, we correct a number of errors that have previously appeared in the literature.

A number of recent papers have considered G_2 modular tensor categories. Kuperberg [25] gives a skein-theoretic description of quantum G_2 spider categories, which for generic q agree with the fusion category of representations of the quantum group [28]. In the case of q a root of unity, Ostrik and Snyder have shown that the spider category agrees with the category of tilting modules for G_2 except for a few small roots of unity q [30]. As a consequence, the semisimplified versions of these categories obtained by quotienting by the ideal of negligible morphisms, again agree with the fusion category of representations of the quantum group. These categories are in fact C^* -tensor categories and have been shown to have property (T) [19].

The paper is organised as follows. In Section 2 we present some preliminary material, including a discussion on the relation between spectral measures over certain different

domains in Section 2.1 and a summary of relevant results for G_2 and its nimreps from [10].

In Section 3 we discuss the finite subgroups of G_2 , including their embeddings into the fundamental representations of G_2 . We also give a general discussion of their spectral measures. In Section 4 we discuss each finite subgroup of G_2 individually, including determining all non-conjugate embeddings into the fundamental representations of G_2 . For each such embedding we construct their McKay graphs, some of which have appeared before in [17] whilst the rest are new, and we determine their joint spectral measures.

2 Preliminaries

2.1 Spectral measures over different domains

The Weyl group of G_2 is the dihedral group D_{12} of order 12. As a subgroup of $GL(2, \mathbb{Z})$, D_{12} is generated by matrices T_2, T_6 , of orders 2, 6 respectively, given by

$$T_2 = \begin{pmatrix} 0 & -1 \\ -1 & 0 \end{pmatrix}, \quad T_6 = \begin{pmatrix} 0 & 1 \\ -1 & 1 \end{pmatrix}, \quad (2)$$

where the action of D_{12} on \mathbb{T}^2 is given by $T(\omega_1, \omega_2) = (\omega_1^{a_{11}}\omega_2^{a_{12}}, \omega_1^{a_{21}}\omega_2^{a_{22}})$, for $T = (a_{ij}) \in D_{12}$. This action leaves $\chi_\mu(\omega_1, \omega_2)$ invariant, for any $\mu \in P_{++} = \{(\mu_1, \mu_2) \in \mathbb{N}^2 \mid \mu_1 \geq \mu_2\}$, the interior of the Weyl alcove for G_2 . Any D_{12} -invariant measure ε_μ on \mathbb{T}^2 yields a pushforward probability measure ν_μ on $I_\mu = \chi_\mu(\mathbb{T}^2) \subset \mathbb{R}$ by

$$\int_{I_\mu} \psi(x) d\nu_\mu(x) = \int_{\mathbb{T}^2} \psi(\chi_\mu(\omega_1, \omega_2)) d\varepsilon_\mu(\omega_1, \omega_2), \quad (3)$$

for any continuous function $\psi : I_\mu \rightarrow \mathbb{C}$, where $d\varepsilon_\mu(\omega_1, \omega_2) = d\varepsilon_\mu(g(\omega_1, \omega_2))$ for all $g \in D_{12}$. There is a loss of dimension here, in the sense that the integral on the right hand side is over the two-dimensional torus \mathbb{T}^2 , whereas on the right hand side it is over the interval I_μ . Thus there is an infinite family of pullback measures ε_μ over \mathbb{T}^2 for any measure ν_μ on I_μ , that is, any ε_μ such that $\varepsilon_\mu(I_\mu^{-1}[x]) = \nu_\mu(x)$ for all $x \in I_\mu$ will yield the probability measure ν_μ on I_μ as a pushforward measure by (3). As in [10], we instead work with an intermediate probability measure $\tilde{\nu}_{\lambda, \mu}$ which lives over the joint spectrum $\mathfrak{D}_{\lambda, \mu} \subset I_\lambda \times I_\mu \subset \mathbb{R}^2$, for $\lambda, \mu \in P_+$, where there is no loss of dimension.

A fundamental domain C of \mathbb{T}^2 under the action of the dihedral group D_{12} is illustrated in Figure 2, where the axes are labelled by the parameters θ_1, θ_2 in $(e^{2\pi i\theta_1}, e^{2\pi i\theta_2}) \in \mathbb{T}^2$, which is a quotient of the fundamental domain of \mathbb{T}^2/S_3 illustrated in Figure 1 (see [8]) by the \mathbb{Z}_2 -action given by the matrix -1 . Note that in Figure 2, the lines $\theta_1 = 0$ and $\theta_2 = 0$ are also boundaries of copies of the fundamental domain C under the action of D_{12} , whereas in Figure 1 they are not boundaries of copies of the fundamental domain under the action of S_3 . The torus \mathbb{T}^2 contains 12 copies of C , so that

$$\int_{\mathbb{T}^2} \phi(\omega_1, \omega_2) d\varepsilon(\omega_1, \omega_2) = 12 \int_C \phi(\omega_1, \omega_2) d\varepsilon(\omega_1, \omega_2), \quad (4)$$

for any D_{12} -invariant function $\phi : \mathbb{T}^2 \rightarrow \mathbb{C}$ and D_{12} -invariant measure ε over \mathbb{T}^2 . The only fixed point of \mathbb{T}^2 under the action of D_{12} is the point $(1, 1)$.

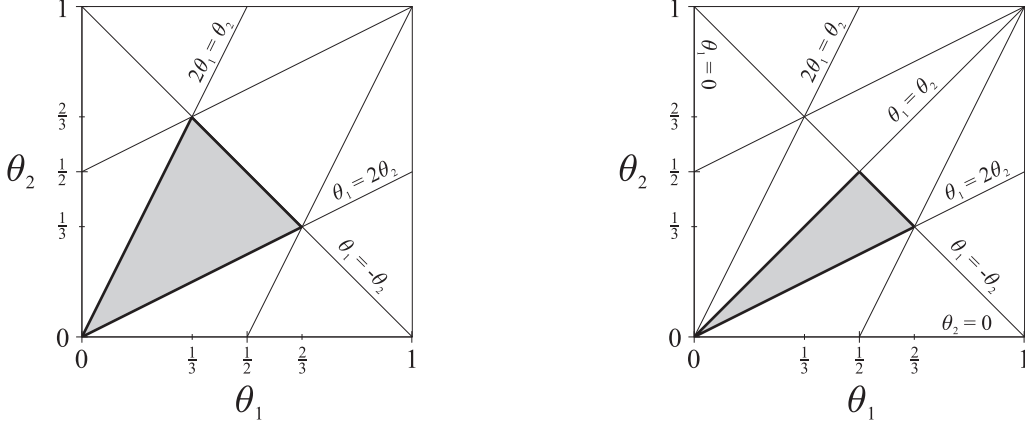


Figure 1: A fundamental domain of \mathbb{T}^2/S_3 . Figure 2: A fundamental domain C of \mathbb{T}^2/D_{12} .

Let $x_\lambda = \chi_\lambda(\omega_1, \omega_2)$ and let $\Psi_{\lambda, \mu}$ be the map $(\omega_1, \omega_2) \mapsto (x_\lambda, x_\mu)$. We denote by $\mathfrak{D}_{\lambda, \mu}$ the image of $\Psi_{\lambda, \mu}(C) (= \Psi_{\lambda, \mu}(\mathbb{T}^2))$ in \mathbb{R}^2 . Note that we can identify $\mathfrak{D}_{\lambda, \mu}$ with $\mathfrak{D}_{\mu, \lambda}$ by reflecting about the line $x_\lambda = x_\mu$. Then the joint spectral measure $\tilde{\nu}_{\lambda, \mu}$ is the measure on $\mathfrak{D}_{\lambda, \mu}$ uniquely determined by its cross-moments as in (1). Then there is a unique D_{12} -invariant pullback measure ε on \mathbb{T}^2 such that

$$\int_{\mathfrak{D}_{\lambda, \mu}} \psi(x_\lambda, x_\mu) d\tilde{\nu}_{\lambda, \mu}(x_\lambda, x_\mu) = \int_{\mathbb{T}^2} \psi(\chi_\lambda(\omega_1, \omega_2), \chi_\mu(\omega_1, \omega_2)) d\varepsilon_{\lambda, \mu}(\omega_1, \omega_2), \quad (5)$$

for any continuous function $\psi : \mathfrak{D}_{\lambda, \mu} \rightarrow \mathbb{C}$.

Any probability measure on $\mathfrak{D}_{\lambda, \mu}$ yields a probability measure on the interval I_λ , given by the pushforward $(p_\lambda)_*(\tilde{\nu}_{\lambda, \mu})$ of the joint spectral measure $\tilde{\nu}_{\lambda, \mu}$ under the orthogonal projection p_λ onto the spectrum $\sigma(\lambda)$ (see [10] for more details). Since the spectral measure ν_λ over I_λ is also uniquely determined by its (one-dimensional) moments $\tilde{\zeta}_m = \int_{I_\lambda} x_\lambda^m d\nu_\lambda(x_\lambda)$ for all $m \in \mathbb{N}$, one could alternatively consider the moments in (1) with $n = 0$ to determine the measure ν_λ over I_λ .

Let

$$C_k^W = \{(e^{2\pi i q_1/3(k+4)}, e^{2\pi i q_2/3(k+4)}) \in \mathbb{T}^2 \mid q_1, q_2 = 0, 1, \dots, 3k+11; q_1 + q_2 \equiv 0 \pmod{3}\} \quad (6)$$

which is the support (over \mathbb{T}^2) of the spectral measure of the nimrep graph $\mathcal{A}_k(G_2)$ associated to the trivial G_2 modular invariant at level k . The following G_2 -invariant measures will be useful later, c.f. [9].

Definition 2.1. Let $\omega = e^{2\pi i/3}$, $\tau = e^{2\pi i/n}$. We define the following measures on \mathbb{T}^2 :

- (1) $d_m \times d_n$, where d_k is the uniform measure on the k^{th} roots of unity, for $k \in \mathbb{N}$.
- (2) $d^{(n)}$, the uniform measure on C_n^W for $n \in \mathbb{N}$.
- (3) $d^{((n))}$, the uniform measure on the S_3 -orbit of the points (τ, τ) , $(\bar{\omega}\bar{\tau}, \omega)$, $(\omega, \bar{\omega}\bar{\tau})$, for $n \in \mathbb{Q}$, $n \geq 2$.
- (4) $d^{(n,k)}$, the uniform measure on the S_3 -orbit of the points $(\tau e^{2\pi i k}, \tau)$, $(\tau, \tau e^{2\pi i k})$, $(\bar{\omega}\bar{\tau}, \omega e^{2\pi i k})$, $(\omega e^{2\pi i k}, \bar{\omega}\bar{\tau})$, $(\bar{\omega}\bar{\tau} e^{-2\pi i k}, \omega e^{-2\pi i k})$, $(\omega e^{-2\pi i k}, \bar{\omega}\bar{\tau} e^{-2\pi i k})$, for $n, k \in \mathbb{Q}$, $n > 2$, $0 \leq k \leq 1/n$.

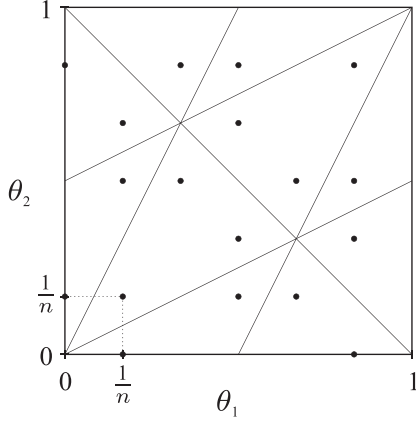


Figure 3: $\text{Supp}(d^{(n)})$

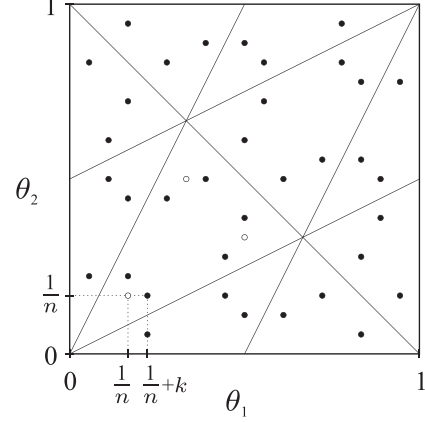


Figure 4: $\text{Supp}(d^{(n,k)})$

The sets $\text{Supp}(d^{(n)})$, $\text{Supp}(d^{(n,k)})$ are illustrated in Figures 3, 4 respectively, where $\text{Supp}(d\mu)$ denotes the set of points $(\theta_1, \theta_2) \in [0, 1]^2$ such that $(e^{2\pi i\theta_1}, e^{2\pi i\theta_2})$ is in the support of the measure $d\mu$. The white circles in Figure 4 denote the points given by the measure $d^{(n)}$. The cardinality $|\text{Supp}(d_m \times d_n)|$ of $\text{Supp}(d_m \times d_n)$ is mn , whilst $|\text{Supp}(d^{(n)})| = |D_n| = 3n^2$ was shown in [8, Section 7.1]. For $n > 2$ and $0 < k < 1/n$, $|\text{Supp}(d^{(n)})| = 18$, whilst $|\text{Supp}(d^{(n,k)})| = 36$. The cardinalities of the other sets are $|\text{Supp}(d^{(n,0)})| = |\text{Supp}(d^{(n,1/n)})| = 18$ for $n > 2$, and $|\text{Supp}(d^{(2)})| = 9$. Some relations between these measures are given in [9, Section 2].

2.2 Spectral measures for G_2

Here we review the results, determined in [10], for the spectral measures for G_2 . Let ρ_1, ρ_2 denote the fundamental representations of G_2 of dimensions 7, 14 respectively. The restrictions of the characters χ_{ρ_j} of the fundamental representations of G_2 to \mathbb{T}^2 yield maps from the torus to the interval $I_j = \chi_{\rho_j}(\mathbb{T}^2) \subset \mathbb{R}$:

$$\begin{aligned} \chi_{\rho_1}(\omega_1, \omega_2) &= 1 + \omega_1 + \omega_1^{-1} + \omega_2 + \omega_2^{-1} + \omega_1\omega_2^{-1} + \omega_1^{-1}\omega_2 \\ &= 1 + 2\cos(2\pi\theta_1) + 2\cos(2\pi\theta_2) + 2\cos(2\pi(\theta_1 - \theta_2)), \\ \chi_{\rho_2}(\omega_1, \omega_2) &= \chi_{\rho_1}(\omega_1, \omega_2) + 1 + \omega_1\omega_2 + \omega_1^{-1}\omega_2^{-1} + \omega_1^2\omega_2^{-1} + \omega_1^{-2}\omega_2 + \omega_1\omega_2^{-2} + \omega_1^{-1}\omega_2^2 \\ &= \chi_{\rho_1}(\omega_1, \omega_2) + 1 + 2\cos(2\pi(\theta_1 + \theta_2)) + 2\cos(2\pi(2\theta_1 - \theta_2)) + 2\cos(2\pi(\theta_1 - 2\theta_2)), \end{aligned}$$

where $\omega_j = e^{2\pi i\theta_j} \in \mathbb{T}$ for $\theta_j \in [0, 1]$, $j = 1, 2$.

Let

$$x := \chi_{\rho_1}(\omega_1, \omega_2), \quad y := \chi_{\rho_2}(\omega_1, \omega_2), \quad (7)$$

and denote by Ψ the map $\Psi_{(1,0),(1,1)} : (\omega_1, \omega_2) \mapsto (x, y)$. The image $\mathfrak{D} = \Psi(C)$ of the fundamental domain C of \mathbb{T}^2/D_{12} (illustrated in Figure 2) under Ψ is illustrated in Figure 5, where the boundaries of C given by $\theta_1 = 2\theta_2$, $\theta_1 = -\theta_2$, $\theta_1 = \theta_2$ yield the curves c_1, c_2, c_3 respectively. These curves are given by [32]

$$\begin{aligned} c_1 : & \quad y = -5(x+1) + 2(x+2)^{3/2}, & x \in [-2, 7], \\ c_2 : & \quad y = -5(x+1) - 2(x+2)^{3/2}, & x \in [-2, -1], \\ c_3 : & \quad 4y = x^2 + 2x - 7, & x \in [-1, 7]. \end{aligned}$$

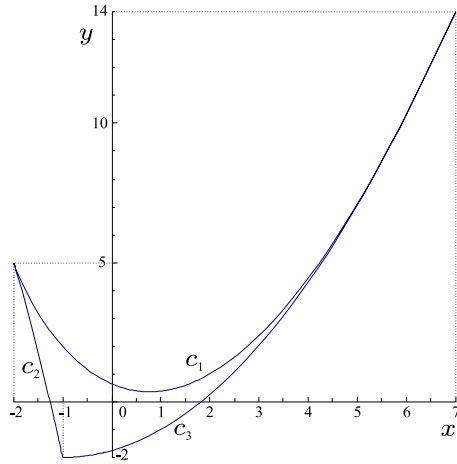


Figure 5: The domain $\mathfrak{D} = \Psi(C)$.

The fixed point $(1, 1)$ of \mathbb{T}^2 under the action of D_{12} maps to 7, 14 in the intervals I_1, I_2 respectively.

Under the change of variables (7) the Jacobian is given by [10]

$$J = 8\pi^2(\cos(2\pi(2\theta_1 + \theta_2)) + \cos(2\pi(\theta_1 - 3\theta_2)) + \cos(2\pi(3\theta_1 - 2\theta_2)) - \cos(2\pi(\theta_1 + 2\theta_2)) - \cos(2\pi(3\theta_1 - \theta_2)) - \cos(2\pi(2\theta_1 - 3\theta_2))). \quad (8)$$

The Jacobian is real and vanishes in \mathbb{T}^2 only on the boundaries of the images of the fundamental domain C under D_{12} . Again, J^2 can be written in terms of the D_{12} -invariant elements x, y as $J^2 = (4x^3 - x^2 - 2x - 10xy - y^2 - 10y + 7)(x^2 + 2x - 7 - 4y)$ (see also [32]), which is non-negative since J is real. We write J in terms of x and y as

$$|J| = 4\pi^2 \sqrt{(4x^3 - x^2 - 2x - 10xy - y^2 - 10y + 7)(x^2 + 2x - 7 - 4y)}. \quad (9)$$

The joint spectral measure over \mathfrak{D} is then $|J(x, y)|dx dy/16\pi^4$, which is the reduced Haar measure on G_2 [32, Sect. 6.3].

2.3 Spectral measures for nimrep graphs associated to G_2 modular invariants

Suppose G is the nimrep associated to a G_2 braided subfactor at some finite level k with vertex set G_0 . We define a state φ on $\ell^2(G_0)$ by $\varphi(\cdot) = \langle \cdot e_*, e_* \rangle$, where e_* is the basis vector in $\ell^2(G_0)$ corresponding to the distinguished vertex $*$ with lowest Perron-Frobenius weight.

If we consider the nimrep graphs G_λ, G_μ , which have joint spectrum $\mathfrak{D}_{\lambda, \mu}$, then the m, n^{th} cross moment $\varsigma_{m, n} = \varphi(G_\lambda^m G_\mu^n) = \int_{\mathfrak{D}_{\lambda, \mu}} x^m y^n d\tilde{\nu}(x, y)$, where $x = x_\lambda, y = x_\mu$, is given by $\langle G_\lambda^m G_\mu^n e_*, e_* \rangle$. Let $\beta_\lambda^{(\nu)} = \chi_\lambda(t_\nu)$ be the eigenvalues of G_λ , where $t_\nu = (\exp(\xi(\nu_1 + 1)), \exp(-3\xi(\nu_2 + 1)))$ for $\xi = 6\pi i/(k+4)$, with corresponding eigenvectors $\psi^{(\nu)}$ (note that the eigenvectors of G_λ are the same for all λ , since the G_λ are a family of commuting matrices). Each eigenvalue $\beta_\lambda^{(\mu)}$ is also given by a ratio of the modular S -matrix for G_2 at

level k , $\beta_\lambda^{(\mu)} = S_{\lambda_\mu}/S_{0\mu}$, where $\mu \in \text{Exp}(G) \subset P_+^k = \{(\lambda_1, \lambda_2) \mid \lambda_1, \lambda_2 \geq 0; \lambda_1 + 2\lambda_2 \leq k\}$ are given by the modular invariant Z . Then $G_\lambda^m G_\mu^n = \mathcal{U} \Lambda_\lambda^m \Lambda_\mu^n \mathcal{U}^*$, where Λ_λ is the diagonal matrix with the eigenvalues $\beta_\lambda^{(\nu)}$ on the diagonal, and \mathcal{U} is the matrix whose columns are given by the eigenvectors $\psi^{(\nu)}$, so that

$$\varsigma_{m,n} = \langle \mathcal{U} \Lambda_\lambda^m \Lambda_\mu^n \mathcal{U}^* e_*, e_* \rangle = \langle \Lambda_\lambda^m \Lambda_\mu^n \mathcal{U}^* e_*, \mathcal{U}^* e_* \rangle = \sum_{\nu} (\beta_\lambda^{(\nu)})^m (\beta_\mu^{(\nu)})^n |\psi_*^{(\nu)}|^2, \quad (10)$$

where $\psi_*^{(\nu)} = \mathcal{U}^* e_*$ is the entry of the eigenvector $\psi^{(\nu)}$ corresponding to the distinguished vertex $*$. Then there is a D_{12} -invariant measure ε over \mathbb{T}^2 such that

$$\varsigma_{m,n} = \int_{\mathbb{T}^2} \chi_\lambda(\omega_1, \omega_2)^m \chi_\mu(\omega_1, \omega_2)^n d\varepsilon(\omega_1, \omega_2),$$

for all λ, μ .

Note from (10) that the measure ε is a discrete measure which has weight $|\psi_*^{(\nu)}|^2$ at the points $g(t_\nu) \in \mathbb{T}^2$ for $g \in D_{12}$, $\nu \in \text{Exp}(G)$, and zero everywhere else. Thus the measure ε does not depend on the choice of λ, μ , so that the spectral measure over \mathbb{T}^2 is the same for any pair (G_λ, G_μ) , even though the corresponding measures over $\mathfrak{D}_{\lambda,\mu} \subset \mathbb{R}^2$, and indeed the subsets $\mathfrak{D}_{\lambda,\mu}$ themselves, are different for each such pair. The same result holds for the spectral measure over \mathbb{T}^2 of a finite subgroup of G_2 .

3 The finite subgroups of G_2

The classification of finite subgroups of G_2 is due to [33, 3] (see also [14, 17]).

The reducible (i.e. block-diagonalizable) finite subgroups of G_2 are the finite discrete subgroups of $SU(2) \times SU(2)$ and $SU(3)$ [33]. These subgroups are thus well known, and the corresponding spectral measures can be obtained from [1, 8, 9].

The irreducible finite subgroups of G_2 , of which there are seven up to conjugacy in G_2 (or equivalently, up to conjugacy in $GL(V)$, where V is the natural 7-dimensional module for $O(7, \mathbb{C})$ [14, Corollary 1]), can be further classified into two types, primitive and imprimitive, where a linear group $\Gamma \subset GL(V)$ is imprimitive if there is a non-trivial decomposition $V = \bigoplus_i V_i$ such that Γ permutes the V_i . There are two imprimitive finite subgroups and five primitive ones. These finite subgroups are listed in Table 1, where type denotes whether an irreducible subgroup is primitive (P) or imprimitive (I).

The McKay graph \mathcal{G}_Γ^ρ is the the fusion graph of the irreducible representation ρ of Γ acting on the irreducible representations of Γ . This graph determines the Bratteli diagram for the tower of relative commutants of the subfactor $P^\Gamma \subset (M_n \otimes P)^\Gamma$, where n is the dimension the representation ρ and P is the type II₁ factor $\bigotimes_{n=1}^\infty M_n$ [34, §VI]. This graph is not however the principal graph of this subfactor as it is not bipartite. The principal graph is rather an unfolded version of the McKay graph \mathcal{G}_Γ^ρ , with adjacency matrix given by $\begin{pmatrix} 0 & \Delta_\Gamma^\rho \\ \Delta_\Gamma^\rho & 0 \end{pmatrix}$, where Δ_Γ^ρ is the adjacency matrix of the (folded) graph \mathcal{G}_Γ^ρ .

We will consider the (joint) spectral measure for the McKay graphs $\mathcal{G}_\Gamma^j := \mathcal{G}_\Gamma^{\rho_j}$ associated to a finite subgroup $\Gamma \subset G_2$, where ρ_j are the restrictions of the fundamental representations ρ_j of G_2 to Γ , $j = 1, 2$. We will consider all possible embeddings of the

Subgroup $\Gamma \subset G_2$	Type	$ \Gamma $
finite subgroups of $SU(2) \times SU(2), SU(3)$	-	-
$PSL(2; 7) \cong GL(3; 2) \cong \Sigma(168) \subset SU(3)$	I	168
$PSL(2; 7) \rtimes \mathbb{Z}_2^3$	I	1344
$PGL(2; 7)$	P	336
$PSL(2; 8)$	P	504
$PSL(2; 13)$	P	1092
$PU(3; 3) \cong G_2(2)'$	P	6048
$G_2(2)$	P	12096

Table 1: Finite subgroups of G_2 .

subgroup in G_2 . Any two such embeddings are conjugate in G_2 if and only if they afford the same character on the seven-dimensional representation ρ_1 [14, Corollary 1]. In some cases there is more than one non-conjugate embedding of the subgroup in G_2 . In these cases the restricted representation ϱ_1 is not necessarily irreducible. In all cases, even for irreducible ϱ_1 , the restriction ϱ_2 of the 14-dimensional representation is not necessarily irreducible.

We use the following methods to determine embeddings ϱ_1 of Γ in G_2 . First, take a seven-dimensional (not necessarily irreducible) representation γ_1 of Γ . The Kronecker square of ρ_1 decomposes into irreducible representations of G_2 as $\rho_1^2 = \text{id}_{G_2} + \rho_1 + \rho_2 + \lambda_{(2,0)}$, where $\lambda_{(2,0)}$ has dimension 27. The Kronecker square of ϱ_1 is obtained by restricting this decomposition to Γ , and we see that ϱ_1 appears in the decomposition of ϱ_1^2 into irreducible representations of Γ . Thus, if γ_1 is not contained in the decomposition of γ_1^2 into irreducible representations of Γ , we can eliminate γ_1 as a possible restriction ϱ_1 of ρ_1 . The decomposition of γ_1^2 into irreducible representations can be obtained using the character table for Γ , by decomposing the character $\chi_{\gamma_1^2} = \chi_{\gamma_1}^2$ of γ_1^2 into the characters of the irreducible representations λ of Γ : $\chi_{\gamma_1^2} = \sum_{\lambda} a_{\lambda} \chi_{\lambda}$, where $a_{\lambda} = \langle \gamma_1^2, \lambda \rangle / |\Gamma| = \sum_{g \in \Gamma} \chi_{\gamma_1^2}(g) \chi_{\lambda}(g) / |\Gamma|$.

We next consider the eigenvalues of the representation matrices of Γ . If the elements in a conjugacy class C_n of Γ have order n , then $(C_n)^n = Z(\Gamma)$, where $Z(\Gamma)$ is the center of Γ . If the center is trivial, which is the case for all subgroups of G_2 , then the eigenvalues ξ of the matrices representing these elements must satisfy $\xi^n = 1$. Since $\chi_{\lambda}(\Gamma_j)$ is the sum of the eigenvalues ξ , it is usually possible to write down the complete set of eigenvalues from the information provided by the character table of Γ and the fact that the eigenvalues must be powers of n^{th} roots of unity. Where there is some ambiguity, we can pin down the correct choice for the set of eigenvalues from the following considerations. Suppose there is ambiguity regarding the eigenvalues of the conjugacy class C_{mn} whose elements have order mn , $m, n \in \mathbb{N}$. If there is only one conjugacy class C_n whose elements have order n , then for $g \in C_{mn}$, $g^m \in C_n$, and since g, g^m commute, their corresponding representation matrices can be simultaneously diagonalised, and thus the eigenvalues of C_n must be m^{th} powers of those for C_{mn} . Suppose now that there is more than one conjugacy class $C_n^{(j)}$ whose elements have order n . Since g^m are all conjugate for conjugate g , we see that there exists a j such that $g^m \in C_n^{(j)}$ for all $g \in C_{mn}$. It turns out in all the cases considered here that there is only one consistent choice of j such that the eigenvalues of $C_n^{(j)}$ are m^{th}

powers of those for C_{mn} for all (irreducible) representations.

As was shown in [8, §4], the eigenvalues of the representation matrices of Γ can be written in the form $\chi_{\varrho_j}(C) = \text{Tr}(\varrho_j(g))$, where g is any element of the conjugacy class C of Γ . Every element $g \in \Gamma$ is conjugate to an element d in the maximal torus of G_2 , i.e. $\varrho_j(h^{-1}gh) = \varrho_j(d) = (\varrho_j|_{\mathbb{T}^2})(t_1, t_2)$ for some $(t_1, t_2) \in \mathbb{T}^2$, for $j = 1, 2$, where $\varrho_j|_{\mathbb{T}^2}$ is given by [10]

$$(\rho_1|_{\mathbb{T}^2})(t_1, t_2) = \text{diag}(D(t_1), D(t_2^{-1}), D(t_1^{-1}t_2), 1), \quad (11)$$

$$\begin{aligned} & (\rho_2|_{\mathbb{T}^2})(t_1, t_2) \\ &= \text{diag}(D(t_1), D(t_2^{-1}), D(t_1^{-1}t_2), D(1), D(t_1t_2), D(t_1^2t_2^{-1}), D(t_1^{-1}t_2^2)), \end{aligned} \quad (12)$$

where $D(t_i) = \begin{pmatrix} \text{Re}(t_i) & -\text{Im}(t_i) \\ \text{Im}(t_i) & \text{Re}(t_i) \end{pmatrix}$ for $t_i \in \mathbb{T}$. Now $\text{Tr}(\varrho_j(g)) = \text{Tr}(\varrho_j(d)) = \Phi_j(t_1, t_2)$, thus the eigenvalues of $\varrho_j(g)$ are all of the form $\Phi_j(\omega_1, \omega_2)$ for $\omega_1, \omega_2 \in \mathbb{T}$, and hence its spectrum is contained in the interval I_j . As shown in [10, Sections 3,4] the spectrum of the fundamental representation ρ_j of G_2 , and its restriction to \mathbb{T}^2 , is the whole of the interval $I_j = \chi_{\rho_j}(\mathbb{T}^2)$, for $j = 1, 2$. Thus the support of the spectral measure μ_{Δ_j} of $\Delta_j = \Delta_{\mathcal{G}_\Gamma^j}$, the adjacency matrix of \mathcal{G}_Γ^j , is contained in I_j when Γ is G_2 or one of its finite subgroups. Then for $\Gamma \subset G_2$, the eigenvalues of every group element in ϱ_1 are necessarily of the form $\mathcal{E}_{t_1, t_2} := \{1, t_1, t_1^{-1}, t_2, t_2^{-1}, t_1t_2^{-1}, t_1^{-1}t_2\}$, where $t_i \in \mathbb{T}$. Thus, by [22, Proposition], if the eigenvalues of the group elements in the representation γ_1 have this form then γ_1 is a restriction ϱ_1 of the seven-dimensional fundamental representation ρ_1 of G_2 to Γ .

We now turn to consider the possible restrictions ϱ_2 of the fourteen-dimensional fundamental representation ρ_2 of G_2 to Γ , for fixed ϱ_1 . By dimension considerations one can determine the possible candidates for ϱ_2 from the Kronecker square of ϱ_1 . Let γ_2 be such a candidate, i.e. a fourteen-dimensional representation such that $\varrho_1^2 = \text{id}_\Gamma + \varrho_1 + \gamma_2 + \lambda$, where λ is (necessarily) some 27-dimensional representation of Γ . We make a choice of pair (t_1^C, t_2^C) from the set X_C of eigenvalues of group elements (from the conjugacy class C) in ϱ_1 such that $\mathcal{E}_{t_1^C, t_2^C} = X_C$. Note that the choice of (t_1^C, t_2^C) such that $\mathcal{E}_{t_1^C, t_2^C} = X_C$ is not unique. However, any other pair $(\tilde{t}_1^C, \tilde{t}_2^C)$ such that $\mathcal{E}_{\tilde{t}_1^C, \tilde{t}_2^C} = X_C$ will appear in the orbit of (t_1^C, t_2^C) under the action of the Weyl group D_{12} of G_2 , where the action of D_{12} on \mathbb{T}^2 is given in Section 2.1. Thus $\Phi_j(\tilde{t}_1^C, \tilde{t}_2^C) = \Phi_j(t_1^C, t_2^C)$ for $j = 1, 2$. Then one checks that $\Phi_2(t_1^C, t_2^C) = \chi_{\gamma_2}(C)$ for each conjugacy class C , in which case γ_2 is indeed a restriction ϱ_2 of the fourteen-dimensional fundamental representation ρ_2 of G_2 to Γ .

The following result present the classification of all non-conjugate embeddings of the finite subgroups of G_2 in G_2 . For the details of the embeddings for each individual subgroup see the corresponding subsections of Section 4.

Theorem 3.1. Up to conjugacy in G_2 the number of faithful embeddings of each finite subgroup of G_2 in the seven-dimensional fundamental representation of G_2 is:

- For $PU(3; 3)$ and $G_2(2)$ there is only one embedding;
- For $PSL(2; 7)$, $\mathbb{Z}_2^3 \cdot PSL(2; 7)$, $PGL(2; 7)$ and $PSL(2; 13)$ there are two non-conjugate embeddings;
- For $PSL(2; 8)$ there are three non-conjugate embeddings.

3.1 Spectral measures for finite subgroups of G_2

For finite groups, the analogue of the S -matrix of Section 2.3, which simultaneously diagonalizes the representations of Γ , is essentially the character table of Γ . This S -matrix, with rows, columns labelled by the irreducible characters, conjugacy classes respectively of Γ , is given by $S_{\lambda,C} = \sqrt{|C|}\chi_\lambda(C)/\sqrt{|\Gamma|}$, for representations λ and conjugacy classes C . Then the analogue of (10) for finite groups is that the m, n^{th} moment $\varsigma_{m,n}$ is given over \mathfrak{D} by [21] (c.f. [8, Section 4] for the case of finite subgroups of $SU(2)$)

$$\varsigma_{m,n} = \int_{\mathfrak{D}} x^m y^n d\nu(x, y) = \sum_C \frac{|C|}{|\Gamma|} \chi_{\varrho_1}(C)^m \chi_{\varrho_2}(C)^n. \quad (13)$$

There is an analogous statement to (13) for the joint spectral measure $\nu_{\lambda,\mu}$ over $\mathfrak{D}_{\lambda,\mu}$ for any irreducible representations λ, μ of Γ . The weight on the right hand side will again be $|C|/|\Gamma|$, since the same S -matrix simultaneously diagonalises all the representations of Γ . Since this weight does not depend on the representations λ, μ , we see that the D_{12} -invariant pullback measure ε over \mathbb{T}^2 will be the same for any joint spectral measure $\nu_{\lambda,\mu}$. This is analogous to the situation for nimrep graphs discussed in Section 2.3.

We wish to compute ‘inverse’ maps $\tilde{\Psi} : \mathfrak{D} \rightarrow \mathbb{T}^2$ such that $\Psi \circ \tilde{\Psi} = \text{id}$. The following equation can easily be checked by substituting in $x = \Phi_1(\omega_1, \omega_2)$, $y = \Phi_2(\omega_1, \omega_2)$:

$$(\omega_j + \omega_j^{-1})^3 + (1 - x)(\omega_j + \omega_j^{-1})^2 + (y - 2)(\omega_j + \omega_j^{-1}) + 2y - x^2 + 2x - 1 = 0,$$

where $j = 1, 2$. Solving this cubic in $\omega_j + \omega_j^{-1} = 2 \cos(\vartheta_j)$, we obtain for $l \in \{0, 1, 2\}$

$$\vartheta_j(l) = \cos^{-1} \left(\frac{1}{6} (x - 1 + 2^{-1/3} \epsilon_l P + 2^{1/3} \bar{\epsilon}_l (x^2 - 2x + 7 - 3y) P^{-1}) \right)$$

where $2^{1/3}$ takes a real value, $\epsilon_l = e^{2\pi i l/3}$ and $P = (2x^3 + 21x^2 - 30x + 7 - 45y - 9xy + \sqrt{(4x^3 - x^2 - 2x - 10xy - y^2 - 10y + 7)(x^2 + 2x - 7 - 4y)})^{1/3}$. We note that for the roots of a cubic equation it does not matter whether the square root in P is taken to be positive or negative. Then we set $\Psi_{l,l'}(x, y) = (e^{\vartheta_1(l)i}, e^{\vartheta_2(l')i})$, and we have that $\Psi(\Psi_{l,l'}(x, y)) = (x, y)$ for some $l, l' \in \{0, 1, 2\}$. The particular choice of pair l, l' such that the equality $\Psi \circ \Psi_{l,l'} = \text{id}$ is satisfied depends on x, y , but it is easy to check (eg. using Mathematica) whether a given choice satisfies this equality for any of the examples we consider. We present in Table 2 the values of the eigenvalues $(\chi_{\varrho_1}(C), \chi_{\varrho_2}(C)) = (x, y) \in \mathfrak{D}$ which will appear for the finite subgroups of G_2 , and (the orbits under D_{12} of) the corresponding points $(\theta_1, \theta_2) \in [0, 1]^2$ such that $\Psi(e^{2\pi i \theta_1}, e^{2\pi i \theta_2}) = (x, y)$.

Let $\Omega(\theta_1, \theta_2) := \Phi_1(e^{2\pi i \theta_1}, e^{2\pi i \theta_2})^m \Phi_2(e^{2\pi i \theta_1}, e^{2\pi i \theta_2})^n \in \mathfrak{D}$ and $\Omega^W(\theta_1, \theta_2)$ its orbit under $W = D_{12}$ given by $\Omega^W(\theta_1, \theta_2) := \sum_{g \in D_{12}} \Omega(g(\theta_1, \theta_2))/12$. We now determine the spectral measures associated with the pairs (θ_1, θ_2) that appear in Table 2.

$(x, y) \in \mathfrak{D}$	Orbit of $(\theta_1, \theta_2) \in [0, 1]^2$	$\frac{J^2(\theta_1, \theta_2)}{16\pi^4}$
(7, 14)	(0, 0)	0
(-2, 5)	$(\frac{1}{3}, \frac{2}{3}), (\frac{2}{3}, \frac{1}{3})$	0
(-1, -2)	$(0, \frac{1}{2}), (\frac{1}{2}, \frac{1}{2}), (\frac{1}{2}, 0)$	0
(1, -1)	$(0, \frac{1}{3}), (\frac{1}{3}, \frac{1}{3}), (\frac{1}{3}, 0), (0, \frac{2}{3}), (\frac{2}{3}, \frac{2}{3}), (\frac{2}{3}, 0)$	0
(3, 2)	$(0, \frac{1}{4}), (\frac{1}{4}, \frac{1}{4}), (\frac{1}{4}, 0), (0, \frac{3}{4}), (\frac{3}{4}, \frac{3}{4}), (\frac{3}{4}, 0)$	0
(-1, 2)	$(\frac{1}{4}, \frac{1}{2}), (\frac{1}{2}, \frac{1}{4}), (\frac{1}{4}, \frac{3}{4}), (\frac{3}{4}, \frac{1}{2}), (\frac{1}{2}, \frac{3}{4}), (\frac{3}{4}, \frac{1}{4})$	0
(2, 1)	$(\frac{1}{6}, \frac{1}{3}), (\frac{1}{3}, \frac{1}{6}), (\frac{1}{6}, \frac{5}{6}), (\frac{5}{6}, \frac{2}{3}), (\frac{2}{3}, \frac{5}{6}), (\frac{5}{6}, \frac{1}{6})$	0
(-1, 1)	$(\frac{1}{6}, \frac{1}{2}), (\frac{1}{2}, \frac{1}{3}), (\frac{1}{3}, \frac{5}{6}), (\frac{5}{6}, \frac{1}{2}), (\frac{1}{2}, \frac{2}{3}), (\frac{2}{3}, \frac{1}{6}),$ $(\frac{1}{2}, \frac{1}{6}), (\frac{1}{3}, \frac{1}{2}), (\frac{5}{6}, \frac{1}{3}), (\frac{1}{2}, \frac{5}{6}), (\frac{2}{3}, \frac{1}{2}), (\frac{1}{6}, \frac{2}{3})$	36
(0, 0)	$(\frac{1}{7}, \frac{3}{7}), (\frac{3}{7}, \frac{2}{7}), (\frac{2}{7}, \frac{6}{7}), (\frac{6}{7}, \frac{4}{7}), (\frac{4}{7}, \frac{5}{7}), (\frac{5}{7}, \frac{1}{7}),$ $(\frac{3}{7}, \frac{1}{7}), (\frac{2}{7}, \frac{3}{7}), (\frac{6}{7}, \frac{2}{7}), (\frac{4}{7}, \frac{6}{7}), (\frac{5}{7}, \frac{4}{7}), (\frac{1}{7}, \frac{5}{7})$	49
(1, 0)	$(\frac{1}{8}, \frac{3}{8}), (\frac{3}{8}, \frac{1}{4}), (\frac{1}{4}, \frac{7}{8}), (\frac{7}{8}, \frac{5}{8}), (\frac{5}{8}, \frac{3}{4}), (\frac{3}{4}, \frac{1}{8}),$ $(\frac{3}{8}, \frac{1}{8}), (\frac{1}{4}, \frac{3}{8}), (\frac{7}{8}, \frac{1}{4}), (\frac{5}{8}, \frac{7}{8}), (\frac{3}{4}, \frac{5}{8}), (\frac{1}{8}, \frac{3}{4})$	32
(-1, 0)	$(\frac{1}{8}, \frac{1}{2}), (\frac{1}{2}, \frac{3}{8}), (\frac{3}{8}, \frac{7}{8}), (\frac{7}{8}, \frac{1}{2}), (\frac{1}{2}, \frac{5}{8}), (\frac{5}{8}, \frac{1}{8}),$ $(\frac{1}{2}, \frac{1}{8}), (\frac{3}{8}, \frac{1}{2}), (\frac{7}{8}, \frac{3}{8}), (\frac{1}{2}, \frac{7}{8}), (\frac{5}{8}, \frac{1}{2}), (\frac{1}{8}, \frac{5}{8})$	32
$(-p, p + q + 1)$	$(\frac{1}{9}, \frac{4}{9}), (\frac{4}{9}, \frac{1}{3}), (\frac{1}{3}, \frac{8}{9}), (\frac{8}{9}, \frac{5}{9}), (\frac{5}{9}, \frac{2}{3}), (\frac{2}{3}, \frac{1}{9}),$ $(\frac{4}{9}, \frac{1}{9}), (\frac{1}{3}, \frac{4}{9}), (\frac{8}{9}, \frac{1}{3}), (\frac{5}{9}, \frac{8}{9}), (\frac{2}{3}, \frac{5}{9}), (\frac{1}{9}, \frac{2}{3})$	$9a_1$
$(-q, 1 - p)$	$(\frac{1}{9}, \frac{1}{3}), (\frac{1}{3}, \frac{2}{9}), (\frac{2}{9}, \frac{8}{9}), (\frac{8}{9}, \frac{2}{3}), (\frac{2}{3}, \frac{7}{9}), (\frac{7}{9}, \frac{1}{9}),$ $(\frac{1}{3}, \frac{1}{9}), (\frac{2}{9}, \frac{1}{3}), (\frac{8}{9}, \frac{2}{9}), (\frac{2}{3}, \frac{8}{9}), (\frac{7}{9}, \frac{2}{3}), (\frac{1}{9}, \frac{7}{9})$	$9a_2$
$(p + q, 1 - q)$	$(\frac{2}{9}, \frac{5}{9}), (\frac{5}{9}, \frac{1}{3}), (\frac{1}{3}, \frac{7}{9}), (\frac{7}{9}, \frac{4}{9}), (\frac{4}{9}, \frac{2}{3}), (\frac{2}{3}, \frac{2}{9}),$ $(\frac{5}{9}, \frac{2}{9}), (\frac{1}{3}, \frac{5}{9}), (\frac{7}{9}, \frac{1}{3}), (\frac{5}{9}, \frac{7}{9}), (\frac{2}{3}, \frac{4}{9}), (\frac{2}{9}, \frac{2}{3})$	$9a_3$
(0, -1)	$(\frac{1}{12}, \frac{5}{12}), (\frac{5}{12}, \frac{1}{3}), (\frac{1}{3}, \frac{11}{12}), (\frac{11}{12}, \frac{7}{12}), (\frac{7}{12}, \frac{2}{3}), (\frac{2}{3}, \frac{1}{12}),$ $(\frac{5}{12}, \frac{1}{12}), (\frac{1}{3}, \frac{5}{12}), (\frac{11}{12}, \frac{1}{3}), (\frac{7}{12}, \frac{11}{12}), (\frac{2}{3}, \frac{7}{12}), (\frac{1}{12}, \frac{2}{3})$	48
$(\frac{1+\sqrt{13}}{2}, 1)$	$(\frac{1}{13}, \frac{4}{13}), (\frac{4}{13}, \frac{3}{13}), (\frac{3}{13}, \frac{12}{13}), (\frac{12}{13}, \frac{9}{13}), (\frac{9}{13}, \frac{10}{13}), (\frac{10}{13}, \frac{1}{13}),$ $(\frac{4}{13}, \frac{1}{13}), (\frac{3}{13}, \frac{4}{13}), (\frac{12}{13}, \frac{3}{13}), (\frac{9}{13}, \frac{12}{13}), (\frac{10}{13}, \frac{9}{13}), (\frac{1}{13}, \frac{10}{13})$	13
$(\frac{1-\sqrt{13}}{2}, 1)$	$(\frac{2}{13}, \frac{7}{13}), (\frac{7}{13}, \frac{5}{13}), (\frac{5}{13}, \frac{11}{13}), (\frac{11}{13}, \frac{6}{13}), (\frac{6}{13}, \frac{8}{13}), (\frac{8}{13}, \frac{2}{13}),$ $(\frac{7}{13}, \frac{2}{13}), (\frac{5}{13}, \frac{7}{13}), (\frac{11}{13}, \frac{5}{13}), (\frac{6}{13}, \frac{11}{13}), (\frac{8}{13}, \frac{6}{13}), (\frac{2}{13}, \frac{8}{13})$	13

Table 2: $(x, y) \in \mathfrak{D}$ and (the orbits of) the corresponding points $(\theta_1, \theta_2) \in [0, 1]^2$. Here $p = 2 \cos(4\pi/9)$, $q = 2 \cos(8\pi/9)$, $a_1 = 3 + 2 \cos(\pi/9) + 2 \cos(2\pi/9)$, $a_2 = 3 - 2 \cos(2\pi/9) + 2 \sin(\pi/18)$, $a_3 = 3 - 2 \cos(\pi/9) - 2 \sin(\pi/18)$

We first note that

$$\Omega^W(0, 0) = \int_{\mathbb{T}^2} \Omega(\theta_1, \theta_2) d_1 \times d_1, \quad (14)$$

$$3\Omega^W(0, 1/2) = \int_{\mathbb{T}^2} \Omega(\theta_1, \theta_2) (4 d_2 \times d_2 - d_1 \times d_1), \quad (15)$$

$$6\Omega^W(0, 1/3) = \int_{\mathbb{T}^2} \Omega(\theta_1, \theta_2) (9 d_3 \times d_3 - 3 d^{(1)}), \quad (16)$$

$$2\Omega^W(1/3, 2/3) = \int_{\mathbb{T}^2} \Omega(\theta_1, \theta_2) (3d^{(1)} - d_1 \times d_1). \quad (17)$$

The orbits of $(0, 1/4)$, $(1/4, 1/2)$ are illustrated in Figure 6, represented by \bullet , $*$ respectively. Both orbits lie on the boundary of the fundamental domains of \mathbb{T}^2/D_{12} , however, only the orbit of $(1/4, 1/2)$ lies on the boundary of the fundamental domains of \mathbb{T}^2/S_3 , illustrated in Figure 1. Then $6K(0, 1/4)\Omega^W(0, 1/4) = 16 \int_{\mathbb{T}^2} \Omega(\theta_1, \theta_2) K(\theta_1, \theta_2) d_4 \times d_4$ where K is an S_3 -invariant function on \mathbb{T}^2 which is zero on the boundaries of the fundamental domains of \mathbb{T}^2/S_3 . Such an S_3 -invariant function is given by the square \tilde{J}^2 of the Jacobian which appeared for the A_2 spectral measures in [8, 9], which is given by $\tilde{J}(\theta_1, \theta_2) = 4\pi^2(\sin(2\pi(\theta_1 + \theta_2)) - \sin(2\pi(2\theta_1 - \theta_2)) - \sin(2\pi(2\theta_2 - \theta_1)))$. Now $|\tilde{J}(0, 1/4)| = 8\pi^2$, thus we take

$$K(\theta_1, \theta_2) = \frac{16\tilde{J}^2(\theta_1, \theta_2)}{64\pi^4} = \frac{\tilde{J}^2(\theta_1, \theta_2)}{4\pi^4}, \quad (18)$$

and we have $K(0, 1/4) = 16$, so that

$$6\Omega^W(0, 1/4) = \int_{\mathbb{T}^2} \Omega(\theta_1, \theta_2) K(\theta_1, \theta_2) d_4 \times d_4. \quad (19)$$

We can also obtain an expression for $\Omega^W(1/4, 1/2)$, where from Figure 6 we see that

$$6\Omega^W(1/4, 1/2) = \int_{\mathbb{T}^2} \Omega(\theta_1, \theta_2) ((16 - K(\theta_1, \theta_2)) d_4 \times d_4 - 4 d_2 \times d_2). \quad (20)$$

Now

$$12\Omega^W(1/6, 1/2) = \int_{\mathbb{T}^2} \Omega(\theta_1, \theta_2) \frac{J^2(\theta_1, \theta_2)}{16\pi^4} d_6 \times d_6, \quad (21)$$

as illustrated in Figure 7 since the Jacobian $J = 0$ along the boundaries of the orbit of the fundamental domain, whilst $J^2(g(1/6, 1/2)/64\pi^4) = 9$ for all $g \in D_{12}$. The points \circ in Figure 7 give the measure $3 d^{(1)}$ whilst the points \diamond in Figure 6 give the measure $4 d_2 \times d_2 - d_1 \times d_1$, thus we see that

$$6\Omega^W(1/6, 1/3) = \int_{\mathbb{T}^2} \Omega(\theta_1, \theta_2) (12 d^{(2)} - 3 d^{(1)} - 4 d_2 \times d_2 + d_1 \times d_1). \quad (22)$$

We have

$$12\Omega^W(1/7, 3/7) = \int_{\mathbb{T}^2} \Omega(\theta_1, \theta_2) \frac{J^2(\theta_1, \theta_2)}{16\pi^4} d_7 \times d_7, \quad (23)$$

as illustrated in Figure 8 since the Jacobian $J = 0$ along the boundaries of the orbit of the fundamental domain, whilst $J^2(g(1/7, 3/7)/64\pi^4) = 49/4$ for all $g \in D_{12}$.

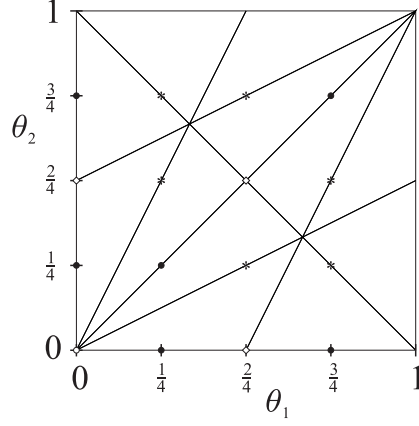


Figure 6: The orbits of $(0, 1/4) \bullet$ and $(1/4, 1/2) *$.

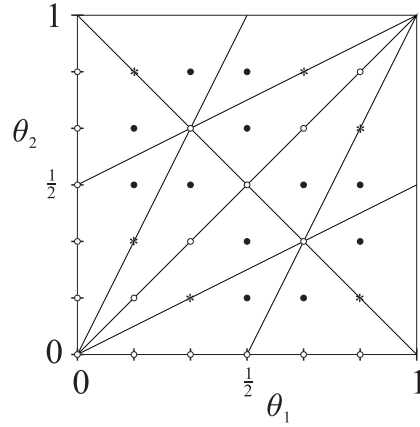


Figure 7: The orbits of $(1/6, 1/2) \bullet$ and $(1/6, 1/3) *$.

The orbits of $(1/8, 3/8)$, $(1/8, 1/2)$ are illustrated in Figure 9, represented by \bullet , $*$ respectively. We have

$$12\Omega^W(1/8, 3/8) = \int_{\mathbb{T}^2} \Omega(\theta_1, \theta_2) \left(\sum_{g \in D_{12}} \delta_{g(e^{\pi i/4}, e^{3\pi i/4})} \right), \quad (24)$$

$$12\Omega^W(1/8, 1/2) = \int_{\mathbb{T}^2} \Omega(\theta_1, \theta_2) \left(\sum_{g \in D_{12}} \delta_{g(e^{\pi i/4}, -1)} \right), \quad (25)$$

where neither orbit gives a linear combination of the measures in Definition 2.1.

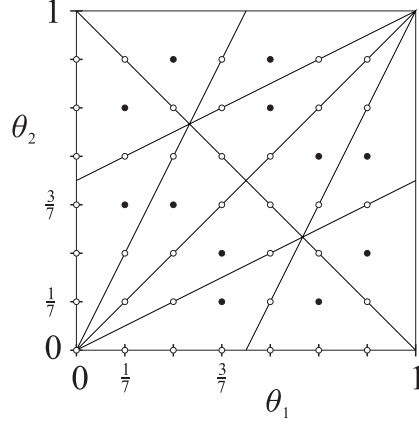


Figure 8: The orbit of $(1/7, 3/7)$.

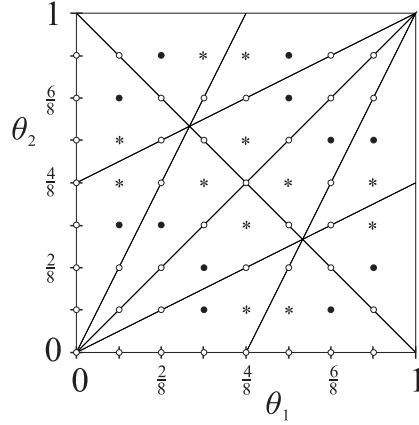


Figure 9: The orbits of $(1/8, 3/8)$ \bullet and $(1/8, 1/2)$ $*$.

Now

$$12a_1\Omega^W(1/9, 4/9) = 2 \int_{\mathbb{T}^2} \Omega(\theta_1, \theta_2) \frac{J^2(\theta_1, \theta_2)}{16\pi^4} d^{((9/2))} \quad (26)$$

$$12a_2\Omega^W(1/9, 1/3) = 2 \int_{\mathbb{T}^2} \Omega(\theta_1, \theta_2) \frac{J^2(\theta_1, \theta_2)}{16\pi^4} d^{((9/4))} \quad (27)$$

$$12a_3\Omega^W(2/9, 5/9) = 2 \int_{\mathbb{T}^2} \Omega(\theta_1, \theta_2) \frac{J^2(\theta_1, \theta_2)}{16\pi^4} d^{((9))}, \quad (28)$$

as illustrated in Figure 10 since the Jacobian $J = 0$ along the boundaries of the orbit of the fundamental domain.

We have

$$12\Omega^W(1/12, 5/12) = \int_{\mathbb{T}^2} \Omega(\theta_1, \theta_2) \frac{J^2(\theta_1, \theta_2)}{16\pi^4} d^{(4)}, \quad (29)$$

as illustrated in Figure 11 since the Jacobian $J = 0$ along the boundaries of the orbit of the fundamental domain, whilst $J^2(g(1/12, 5/12)/16\pi^4) = 48$ for all $g \in D_{12}$.

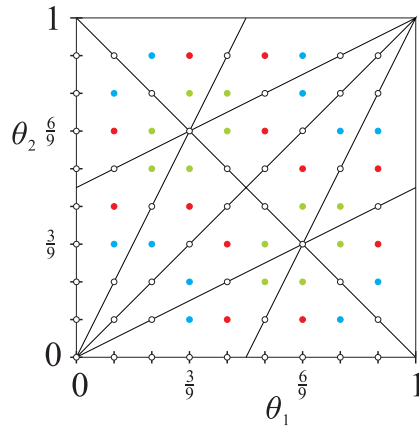


Figure 10: The orbits of $(1/9, 4/9)$, $(1/9, 1/3)$, $(2/9, 5/9)$.

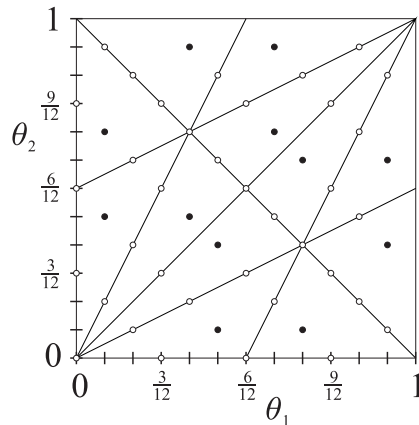


Figure 11: The orbit of $(1/12, 5/12)$.

Finally, the orbit of the points $(1/13, 4/13)$, $(2/13, 7/13)$, illustrated in Figure 12, do not give a linear combination of the measures in Definition 2.1, thus we have

$$12\Omega^W(1/13, 4/13) = \int_{\mathbb{T}^2} \Omega(\theta_1, \theta_2) \left(\sum_{g \in D_{12}} \delta_{g(\zeta, \zeta^4)} \right), \quad (30)$$

$$12\Omega^W(2/13, 7/13) = \int_{\mathbb{T}^2} \Omega(\theta_1, \theta_2) \left(\sum_{g \in D_{12}} \delta_{g(\zeta^2, \zeta^7)} \right). \quad (31)$$

4 Explicit results

4.1 Group $PSL(2; 7) \cong GL(3; 2) \cong \Sigma(168)$

The subgroup $PSL(2; 7)$ of G_2 is an irreducible imprimitive group of order 168 which is isomorphic to the group $GL(3; 2)$, and also to the subgroup $\Sigma(168)$ of $SU(3)$ which was considered in [9].

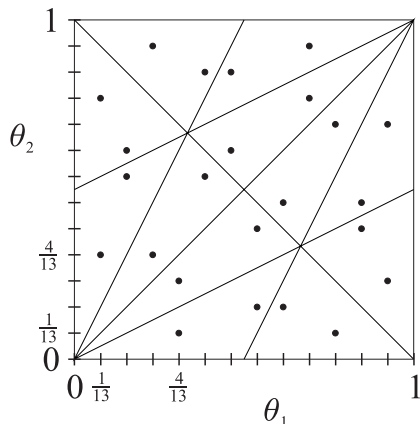


Figure 12: The orbits of $(1/13, 4/13)$ and $(2/13, 7/13)$.

C	C_1	C_2	(C_7, C_7^2, C_7^4)	(C_7^3, C_7^5, C_7^6)	(C_4, C_4^3)	(C_3, C_3^2)
$ C $	1	21	24	24	42	56
Σ_1	1	1	1	1	1	1
Σ_3	3	-1	w	\bar{w}	1	0
Σ_3^*	3	-1	\bar{w}	w	1	0
Σ_6	6	2	-1	-1	0	0
Σ_7	7	-1	0	0	-1	1
Σ_8	8	0	1	1	0	-1

Table 3: Character table for group $PSL(2; 7)$, where $w = \eta + \eta^2 + \eta^4 = (-1 + i\sqrt{7})/2$, $\eta = e^{2\pi i/7}$.

The group $PSL(2; 7)$ has irreducible real representations Σ_d of dimensions $d = 1, 6, 7, 8$, and two complex conjugate irreducible representations Σ_3, Σ_3^* of dimension 3. Its character table is given in Table 3 [26].

There are two non-conjugate embeddings of $PSL(2; 7)$ in G_2 [22], given by $\varrho_1^{(1)} = \Sigma_7$ and $\varrho_1^{(2)} = \Sigma_1 + \Sigma_3 + \Sigma_3^*$. The McKay graph $\mathcal{G}_{PSL(2;7)}^{\varrho_1^{(1)}}$ for $\varrho_1^{(1)}$ is given in [17, Figure 1]. We reproduce it in Figure 13 for completeness, along with the McKay graph $\mathcal{G}_{PSL(2;7)}^{\varrho_1^{(2)}}$ for $\varrho_1^{(2)}$. We use the notation n, n^* to label the vertices corresponding to the irreducible representations Σ_n, Σ_n^* respectively.

The eigenvalues of the representation matrices are given in [22, Tables 4a,b]. The decomposition of the Kronecker square of $\varrho_1^{(i)}$ into irreducibles is given by

$$(\varrho_1^{(1)})^2 = \text{id} + \varrho_1^{(1)} + \Sigma_3 + \Sigma_3^* + 2\Sigma_6 + \Sigma_7 + 2\Sigma_8, \quad (\varrho_1^{(2)})^2 = \text{id} + \varrho_1^{(2)} + \Sigma_1 + 2\Sigma_3 + 2\Sigma_3^* + 2\Sigma_6 + 2\Sigma_8,$$

where $\text{id} = \Sigma_1$. From dimension considerations there are thus two candidates for the fourteen-dimensional representation $\varrho_2^{(i)}$, which are given by $\Sigma_3 + \Sigma_3^* + \Sigma_8$ and $\Sigma_6 + \Sigma_8$ for both $i = 1, 2$. However, as discussed in Section 3, since $\chi_{\varrho_2}(C) = \Phi_2(t_1^C, t_2^C)$, where (t_1^C, t_2^C) is a pair from the set of eigenvalues of group elements from the conjugacy class C , from

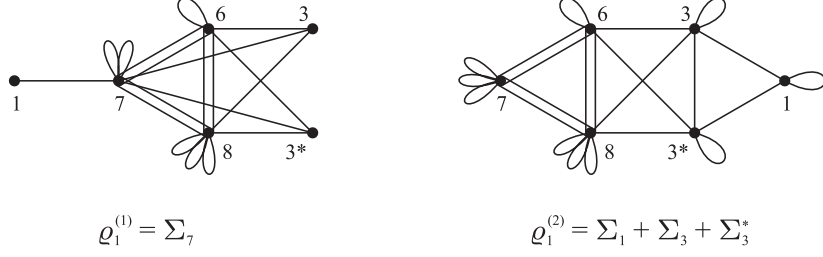


Figure 13: The McKay graphs $\mathcal{G}_{PSL(2;7)}^{\varrho_1^{(i)}}$, $i = 1, 2$.

C	C_1	C_2	(C_7, C_7^2, C_7^4)	(C_7^3, C_7^5, C_7^6)	(C_4, C_4^3)	(C_3, C_3^2)
$\chi_{\varrho_1^{(1)}}(C) \in [-2, 7]$	7	-1	0	0	-1	1
$\chi_{\varrho_1^{(2)}}(C) \in [-2, 7]$	7	-1	0	0	3	1
$\chi_{\varrho_2}(C) \in [-2, 14]$	14	-2	0	0	2	-1
$J^2(\theta_1, \theta_2)/64\pi^4$	0	0	49/4	49/4	0	0

Table 4: $\chi_{\varrho_j}(C)$ for group $PSL(2; 7)$, $j = 1, 2$.

knowledge of the eigenvalues from [22] we see that the decomposition of the fundamental fourteen-dimensional representation into irreducible representations of $PSL(2; 7)$ is given by $\varrho_2 := \varrho_2^{(i)} = \Sigma_3 + \Sigma_3^* + \Sigma_8$ for both $i = 1, 2$. The values of $x^{(i)} = \chi_{\varrho_1^{(i)}}(C) \in [-2, 7]$, $y = \chi_{\varrho_2}(C) \in [-2, 14]$ for $PSL(2; 7)$ are given in Table 4, along with the values of $J^2/64\pi^4$ for the corresponding pairs $(\theta_1, \theta_2) \in [0, 1]^2$ obtained from Table 2.

From (13) and Tables 3, 4 and 2, we see that

$$\varsigma_{m,n} = \frac{1}{168}\Omega^W(0, 0) + \frac{21}{168}\Omega^W(0, 1/2) + \frac{56}{168}\Omega^W(0, 1/3) + \frac{42}{168}\Omega' + \frac{24+24}{168}\Omega^W(1/7, 3/7),$$

where Ω' is $\Omega^W(1/4, 1/2)$ for $\varrho_1^{(1)}$, and $\Omega^W(0, 1/4)$ for $\varrho_1^{(2)}$. Then we obtain

Theorem 4.1. The joint spectral measure (over \mathbb{T}^2) for the non-conjugate embeddings of the projective special linear group $PSL(2; 7)$ over the finite field \mathbb{F}_7 into the fundamental representations of G_2 is

$$d\varepsilon = \frac{1}{672\pi^4}J^2 d_7 \times d_7 + \frac{1}{24}K' d_4 \times d_4 + \frac{1}{2}d_3 \times d_3 + \frac{1}{6}\beta d_2 \times d_2 - \frac{1}{28}d_1 \times d_1 - \frac{1}{6}d^{(1)} \quad (32)$$

where for the embedding of $PSL(2; 7)$ in G_2 given by $\varrho_1^{(1)} = \Sigma_7$ we have $K' = 16 - K$ and $\beta = 1$, whilst for the embedding given by $\varrho_1^{(2)} = \Sigma_1 + \Sigma_3 + \Sigma_3^*$ we have $K' = K$ and $\beta = 0$, where $K(\theta_1, \theta_2) = (\sin(2\pi(\theta_1 + \theta_2)) - \sin(2\pi(2\theta_1 - \theta_2)) - \sin(2\pi(2\theta_2 - \theta_1)))^2$, d_m is the uniform measure over m^{th} roots of unity and $d^{(k+4)}$ is the uniform measure on the points in C_k^W .

Remark 4.2. Note that measure in Theorem 4.1 for the second embedding $\varrho_1^{(2)}$ of $PSL(2; 7)$ in G_2 is precisely that for $\Sigma(168) \subset SU(3)$ given in [9, Theorem 16]. However,

C	C_1	C_2	C_4	C'_4	C_8	C'_8	C'_2	C_7	C'_7	C_6	C_3
$ C $	1	7	42	42	168	168	84	192	192	224	224
Σ_1	1	1	1	1	1	1	1	1	1	1	1
Σ_3	3	3	-1	-1	1	1	-1	w	\bar{w}	0	0
Σ_3^*	3	3	-1	-1	1	1	-1	\bar{w}	w	0	0
Σ_6	6	6	2	2	0	0	2	-1	-1	0	0
$\Sigma_7^{(1)}$	7	-1	-1	3	1	-1	-1	0	0	-1	1
$\Sigma_7^{(1)'}$	7	-1	3	-1	-1	1	-1	0	0	-1	1
$\Sigma_7^{(2)}$	7	7	-1	-1	-1	-1	-1	0	0	1	1
Σ_8	8	8	0	0	0	0	0	1	1	-1	-1
Σ_{14}	14	-2	2	2	0	0	-2	0	0	1	-1
Σ_{21}	21	-3	-3	1	-1	1	1	0	0	0	0
Σ'_{21}	21	-3	1	-3	1	-1	1	0	0	0	0

Table 5: Character table for $\mathbb{Z}_2^3 \cdot PSL(2; 7)$, where $w = \eta + \eta^2 + \eta^4 = (-1 + i\sqrt{7})/2$, $\eta = e^{2\pi i/7}$.

(32) has a cleaner expression than that given in [9], because here we were able to use the Jacobian J for G_2 which is also 0 along the diagonal, whereas the Jacobian for $SU(3)$ (essentially K in Theorem 4.1) is non-zero along the diagonal.

4.2 Group $\mathbb{Z}_2^3 \cdot PSL(2; 7)$

The subgroup $\mathbb{Z}_2^3 \cdot PSL(2; 7)$ of G_2 is a nonsplit extension of \mathbb{Z}_2^3 by $PSL(2; 7)$. It is an irreducible imprimitive group of order 1344. It has eleven irreducible representations (nine real and two complex conjugate representations). As a subgroup of $GL(7)$ it is generated by the permutations (1234567) and (124)(365), and the signed permutation $\delta_{\{1,2,4,7\}}^-(12)(36)$, where for a set S the notation δ_S^- means that the entries in the rows indexed by elements of S are multiplied by -1 [6, 3]. Its character table, generated by the GAP computational discrete algebra system, is given in Table 5 (see also [17]) where elements in $C_4, C'_4, C_8, C'_8, C_7, C'_7, C_6, C_3$ are of cycle type $(C_4, C_4^3), (C'_4, C_4^3), (C_8, C_8^3, C_8^5, C_8^7), (C'_8, C_8^3, C_8^5, C_8^7), (C_7, C_7^2, C_7^4), (C_7^3, C_7^5, C_7^6), (C_6, C_6^5), (C_3, C_3^2)$ respectively.

There are five non-conjugate seven-dimensional representations, $\Sigma_1 + \Sigma_3 + \Sigma_3^*, \Sigma_1 + \Sigma_6, \Sigma_7^{(1)}, \Sigma_7^{(1)'}$ and $\Sigma_7^{(2)}$, however of these, only $\rho_1^{(1)} := \Sigma_7^{(1)}$ and $\rho_1^{(2)} := \Sigma_7^{(1)'}$ are faithful. Both these satisfy the condition that $\rho_1^{(i)}$ appears in the decomposition of $(\rho_1^{(i)})^2$, and their eigenvalues are given in Table 6. As described in Section 3, these eigenvalues can be determined from the character table of $\mathbb{Z}_2^3 \cdot PSL(2; 7)$. The additional information that is needed is to note that the eigenvalues for group elements in C_4 and C'_4 square to those for elements in C_2 , those for C_8 square to those for C_4 , those for C'_8 square to those for C'_4 , whilst those for C_6 square to those for C_3 and also cube to those for C_2 .

We present one choice of eigenvalues (t_1^C, t_2^C) in Table 7, along with the characters of the corresponding fourteen-dimensional representation for these eigenvalues. In both cases we see that the character agrees with the irreducible representation Σ_{14} .

	$\varrho_1^{(1)} = \Sigma_7^{(1)}$	$\varrho_1^{(2)} = \Sigma_7^{(1)'}$
C_1	(1, 1, 1, 1, 1, 1, 1)	(1, 1, 1, 1, 1, 1, 1)
C_2	(1, 1, 1, -1, -1, -1, -1)	(1, 1, 1, -1, -1, -1, -1)
C_4	(1, -1, -1, $i, i, -i, -i$)	(1, 1, 1, $i, i, -i, -i$)
C'_4	(1, 1, 1, $i, i, -i, -i$)	(1, -1, -1, $i, i, -i, -i$)
C_8	(1, -1, -1, ν, ν^3, ν^5, ν^7)	(1, $i, -i, \nu, \nu^3, \nu^5, \nu^7$)
C'_8	(1, $i, -i, \nu, \nu^3, \nu^5, \nu^7$)	(1, -1, -1, ν, ν^3, ν^5, ν^7)
C'_2	(1, 1, 1, -1, -1, -1, -1)	(1, 1, 1, -1, -1, -1, -1)
C_7	(1, $\eta, \eta^2, \eta^3, \eta^4, \eta^5, \eta^6$)	(1, $\eta, \eta^2, \eta^3, \eta^4, \eta^5, \eta^6$)
C'_7	(1, $\eta, \eta^2, \eta^3, \eta^4, \eta^5, \eta^6$)	(1, $\eta, \eta^2, \eta^3, \eta^4, \eta^5, \eta^6$)
C_6	(1, -1, -1, $\omega, -\omega, \omega^2, -\omega^2$)	(1, -1, -1, $\omega, -\omega, \omega^2, -\omega^2$)
C_3	(1, 1, 1, $\omega, \omega, \omega^2, \omega^2$)	(1, 1, 1, $\omega, \omega, \omega^2, \omega^2$)

Table 6: Eigenvalues of group elements in each conjugacy class of $\mathbb{Z}_2^3 \cdot PSL(2; 7)$ for the irreducible representations $\varrho_1^{(i)}$, $i = 1, 2$, where $\omega = e^{2\pi i/3}$, $\nu = e^{2\pi i/8}$ and $\eta = e^{2\pi i/7}$.

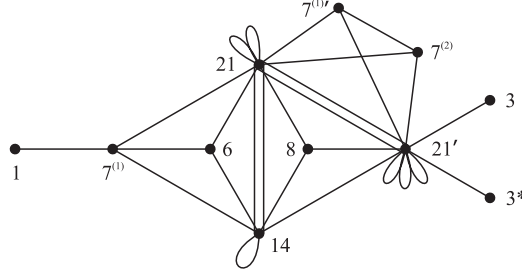


Figure 14: The McKay graph $\mathcal{G}_{\mathbb{Z}_2^3 \cdot PSL(2; 7)}^{\varrho_1^{(1)}}$.

Thus there are precisely two non-conjugate faithful irreducible representations $\varrho_1^{(i)}$, $i = 1, 2$, giving embeddings of $\mathbb{Z}_2^3 \cdot PSL(2; 7)$ in G_2 . The McKay graph $\mathcal{G}_{PSL(2; 7)}^{\varrho_1^{(1)}}$ for $\varrho_1^{(1)}$ is given in [17, Figure 1]. We reproduce it in Figure 14 for completeness. We use the notation $n, n^*, n^{(i)'}$ to label the vertices corresponding to the irreducible representations $\Sigma_n, \Sigma_n^*, \Sigma_n^{(i)'}$ respectively. The McKay graph for $\varrho_1^{(2)}$ is similar.

The values of $x^{(i)} = \chi_{\varrho_1^{(i)}}(C) \in [-2, 7]$, $y = \chi_{\rho_2}(C) \in [-2, 14]$ for $\mathbb{Z}_2^3 \cdot PSL(2; 7)$ are given in Table 4, along with the values of $J^2/64\pi^4$ for the corresponding pairs $(\theta_1, \theta_2) \in [0, 1]^2$ obtained from Table 2. Then from (13) and Tables 5, 8 and 2, we see that

$$\begin{aligned} \varsigma_{m,n} &= \frac{1}{1344} \Omega^W(0, 0) + \frac{7 + 84}{1344} \Omega^W(0, 1/2) + \frac{42}{1344} \Omega^W(0, 1/4) + \frac{42}{1344} \Omega^W(1/4, 1/2) \\ &+ \frac{168}{1344} \Omega^W(1/8, 1/2) + \frac{168}{1344} \Omega^W(1/8, 3/8) + \frac{192 + 192}{1344} \Omega^W(1/7, 3/7) \\ &+ \frac{224}{1344} \Omega^W(1/6, 1/2) + \frac{224}{1344} \Omega^W(0, 1/3), \end{aligned}$$

where $\Omega^W(\theta_1, \theta_2)$ is as in Section 4.1. Note that $12(\Omega^W(1/8, 1/2) + \Omega^W(1/8, 3/8)) = 8 \int_{\mathbb{T}^2} \Omega(\theta_1, \theta_2)(J^2/64\pi^4) d_8 \times d_8$, since the Jacobian $J = 0$ along the boundaries of the orbit

C	(t_1^C, t_2^C)	$\varrho_1^{(1)} \chi_{\varrho_2^{(1)}}(C)$	(t_1^C, t_2^C)	$\varrho_1^{(2)} \chi_{\varrho_2^{(2)}}(C)$	$\chi_{\Sigma_{14}}(C)$
C_1	(1, 1)	14	(1, 1)	14	14
C_2	(1, -1)	-2	(1, -1)	-2	-2
C_4	(-1, i)	2	(1, i)	2	2
C'_4	(1, i)	2	(-1, i)	2	2
C_8	(-1, ν)	0	(i , ν)	0	0
C'_8	(i , ν)	0	(-1, ν)	0	0
C'_2	(1, -1)	-2	(1, -1)	-2	-2
C_7	(η , η^5)	0	(η , η^5)	0	0
C'_7	(η^2 , η^3)	0	(η^2 , η^3)	0	0
C_6	(-1, ω)	1	(-1, ω)	1	1
C_3	(1, ω)	-1	(1, ω)	-1	-1

Table 7: Choice of eigenvalues (t_1^C, t_2^C) for $\varrho_1^{(i)}$, $i = 1, 2$, and corresponding values of $\chi_{\varrho_2^{(i)}}(C)$.

C	C_1	C_2	C_4	C'_4	C_8	C'_8	C'_2	C_7	C'_7	C_6	C_3
$\chi_{\varrho_1^{(1)}}(\Gamma_j) \in [-2, 7]$	7	-1	-1	3	1	-1	-1	0	0	-1	1
$\chi_{\varrho_1^{(2)}}(\Gamma_j) \in [-2, 7]$	7	-1	3	-1	-1	1	-1	0	0	-1	1
$\chi_{\varrho_2^{(2)}}(\Gamma_j) \in [-2, 14]$	14	-2	2	2	0	0	-2	0	0	1	-1
$J^2(\theta_1, \theta_2)/64\pi^4$	0	0	0	0	8	8	0	49/4	49/4	9	0

Table 8: $\chi_{\varrho_j}(C)$ for group $\mathbb{Z}_2^3 \cdot PSL(2; 7)$, $j = 1, 2$.

of the fundamental domain, whilst $J^2(g(1/8, 1/2))/64\pi^4 = J^2(g(1/8, 3/8))/64\pi^4 = 8$ for all $g \in D_{12}$. Thus we obtain:

Theorem 4.3. The joint spectral measure (over \mathbb{T}^2) for both non-conjugate embeddings of $\mathbb{Z}_2^3 \cdot PSL(2; 7)$ into the fundamental representations of G_2 are given by

$$\begin{aligned}
d\varepsilon = & \frac{1}{768\pi^4} J^2 d_8 \times d_8 + \frac{1}{672\pi^4} J^2 d_7 \times d_7 + \frac{1}{1152\pi^4} J^2 d_6 \times d_6 + \frac{1}{12} d_4 \times d_4 \\
& + \frac{1}{4} \beta d_3 \times d_3 + \frac{5}{72} d_2 \times d_2 - \frac{11}{504} d_1 \times d_1 - \frac{1}{12} d^{(1)}
\end{aligned} \tag{33}$$

where d_m is the uniform measure over m^{th} roots of unity and $d^{(k+4)}$ is the uniform measure on the points in C_k^W .

4.3 Group $PGL(2; 7)$

The subgroup $PGL(2; 7)$ of G_2 is an irreducible primitive group of order 336. It has nine irreducible representations, all real, and its character table is given in Table 9 [4, 17].

There are eight non-conjugate seven-dimensional representations, $\gamma_1^{(1)} = \Sigma_1 + \Sigma_6^{(1)}$, $\gamma_1^{(2)} = \Sigma_1 + \Sigma_6^{(2)}$, $\gamma_1^{(3)} = \Sigma_1 + \Sigma_6^{(2)'}$, $\gamma_1^{(4)} = \Sigma'_1 + \Sigma_6^{(1)}$, $\gamma_1^{(5)} = \Sigma'_1 + \Sigma_6^{(2)}$, $\gamma_1^{(6)} = \Sigma'_1 + \Sigma_6^{(2)'}$,

C	C_1	C_2	C'_2	C_8	C'_8	C_4	C_7	C_6	C_3
$ C $	1	21	28	42	42	42	48	56	56
Σ_1	1	1	1	1	1	1	1	1	1
Σ'_1	1	1	-1	-1	-1	1	1	-1	1
$\Sigma_6^{(1)}$	6	-2	0	0	0	2	-1	0	0
$\Sigma_6^{(2)}$	6	2	0	$\sqrt{2}$	$-\sqrt{2}$	0	-1	0	0
$\Sigma_6^{(2)'}$	6	2	0	$-\sqrt{2}$	$\sqrt{2}$	0	-1	0	0
Σ_7	7	-1	1	-1	-1	-1	0	1	1
Σ'_7	7	-1	-1	1	1	-1	0	-1	1
Σ_8	8	0	2	0	0	0	1	-1	-1
Σ'_8	8	0	-2	0	0	0	1	1	-1

Table 9: Character table for $PGL(2; 7)$.

$\gamma_1^{(7)} = \Sigma_7$ and $\gamma_1^{(8)} = \Sigma'_7$. The decomposition of the Kronecker squares of the $\gamma_1^{(i)}$ are given by:

$$\begin{aligned}
(\gamma_1^{(1)})^2 &= \text{id} + \gamma_1^{(1)} + \Sigma'_1 + 2\Sigma_6^{(1)} + \Sigma_6^{(2)} + \Sigma_6^{(2)'} + \Sigma_8 + \Sigma'_8, \\
(\gamma_1^{(2)})^2 &= \text{id} + \gamma_1^{(2)} + 2\Sigma_6^{(2)} + \Sigma_6^{(2)'} + \Sigma'_7 + \Sigma_8 + \Sigma'_8, \\
(\gamma_1^{(4)})^2 &= \text{id} + \gamma_1^{(4)} + \Sigma_1 + 2\Sigma_6^{(1)} + \Sigma_6^{(2)} + \Sigma_6^{(2)'} + \Sigma_8 + \Sigma'_8, \\
(\gamma_1^{(5)})^2 &= \text{id} + \Sigma_1 + 3\Sigma_6^{(2)} + \Sigma_6^{(2)'} + \Sigma'_7 + \Sigma_8 + \Sigma'_8, \\
(\gamma_1^{(7)})^2 &= \text{id} + \gamma_1^{(7)} + \Sigma_1 + \Sigma_6^{(1)} + \Sigma_6^{(2)} + \Sigma_6^{(2)'} + \Sigma'_7 + \Sigma_8 + \Sigma'_8, \\
(\gamma_1^{(8)})^2 &= \text{id} + \gamma_1^{(8)} + \Sigma_1 + \Sigma_6^{(1)} + \Sigma_6^{(2)} + \Sigma_6^{(2)'} + \Sigma_7 + \Sigma_8 + \Sigma'_8,
\end{aligned}$$

where $\text{id} = \Sigma_1$. Note, we have omitted the decompositions for $\gamma_1^{(3)}$, $\gamma_1^{(6)}$, since from the character table we see that these representations are essentially the same as $\gamma_1^{(2)}$, $\gamma_1^{(5)}$ respectively. Then we see that $\gamma_1^{(i)}$ does not appear in the decomposition of $(\gamma_1^{(i)})^2$ for $i = 5, 6$, therefore they are not embeddings of $PGL(2; 7)$ in G_2 . From dimension considerations, we see that candidates $\gamma_2^{(j)}$ for ϱ_2 are $\gamma_2^{(1)} = \Sigma_6^{(1)} + \Sigma_8$, $\gamma_2^{(2)} = \Sigma_6^{(1)} + \Sigma'_8$, $\gamma_2^{(3)} = \Sigma_6^{(2)} + \Sigma_8$, $\gamma_2^{(4)} = \Sigma_6^{(2)} + \Sigma'_8$, $\gamma_2^{(5)} = \Sigma_6^{(2)'} + \Sigma_8$ and $\gamma_2^{(6)} = \Sigma_6^{(2)'} + \Sigma'_8$, where $1 \leq j \leq 6$ for $\gamma_1^{(i)}$ when $i = 1, 4, 7, 8$, and $3 \leq j \leq 6$ when $i = 2$. Then with $(x_i^C, y_j^C) = (\chi_{\gamma_1^{(i)}}(C), \chi_{\gamma_2^{(j)}}(C))$, we see that $(x_i^C, y_j^C) \notin \mathfrak{D}$ for any candidate $\gamma_2^{(j)}$ in the cases $i = 1, 2, 3, 7$ when $C = C'_2$. Thus $\gamma_1^{(i)}$ cannot define an embedding of $PGL(2; 7)$ in G_2 for $i = 1, 2, 3, 7$. We also see that $(x_8^C, y_j^C) \notin \mathfrak{D}$ for $j = 3, 4, 5, 6$.

Thus we have candidates $(\gamma_1^{(i)}, \gamma_2^{(j)})$ for (ϱ_1, ϱ_2) when $i = 4, 8$, where $1 \leq j \leq 6$ for $i = 4$ and $j \in \{1, 2\}$ for $i = 8$.

We now consider the eigenvalues of the representation matrices to determine which of the remaining $\gamma_1^{(i)}$ are embeddings of $PSL(2; 7)$ in G_2 . These eigenvalues are given in Table 10. As described in Section 3, these eigenvalues can be determined from the character table of $PGL(2; 7)$. The additional information that is needed is to note that the eigenvalues for group elements in C_4 square to those for elements in C_2 , those for C_8 ,

	Σ_1	Σ'_1	$\Sigma_6^{(1)}$	$\Sigma_6^{(2)}$	Σ_7
C_1	1	1	(1, 1, 1, 1, 1, 1)	(1, 1, 1, 1, 1, 1)	(1, 1, 1, 1, 1, 1, 1)
C_2	1	1	(1, 1, -1, -1, -1, -1)	(1, 1, 1, 1, -1, -1)	(1, 1, 1, -1, -1, -1, -1)
C'_2	1	-1	(1, 1, 1, -1, -1, -1)	(1, 1, 1, -1, -1, -1)	(1, 1, 1, 1, -1, -1, -1)
C_8	1	-1	(1, -1, ν, ν^3, ν^5, ν^7)	(1, -1, $i, -i, \nu, \nu^7$)	(-1, $i, -i, \nu, \nu^3, \nu^5, \nu^7$)
C'_8	1	-1	(1, -1, ν, ν^3, ν^5, ν^7)	(1, -1, $i, -i, \nu^3, \nu^5$)	(-1, $i, -i, \nu, \nu^3, \nu^5, \nu^7$)
C_4	1	1	(1, 1, $i, i, -i, -i$)	(1, 1, -1, -1, $i, -i$)	(1, -1, -1, $i, i, -i, -i$)
C_7	1	1	($\eta, \eta^2, \eta^3, \eta^4, \eta^5, \eta^6$)	($\eta, \eta^2, \eta^3, \eta^4, \eta^5, \eta^6$)	(1, $\eta, \eta^2, \eta^3, \eta^4, \eta^5, \eta^6$)
C_6	1	-1	(1, -1, μ, μ^2, μ^4, μ^5)	(1, -1, μ, μ^2, μ^4, μ^5)	(1, 1, -1, μ, μ^2, μ^4, μ^5)
C_3	1	1	(1, 1, $\omega, \omega, \omega^2, \omega^2$)	(1, 1, $\omega, \omega, \omega^2, \omega^2$)	(1, 1, 1, $\omega, \omega, \omega^2, \omega^2$)

	Σ'_7	Σ_8	Σ'_8
C_1	(1, 1, 1, 1, 1, 1, 1)	(1, 1, 1, 1, 1, 1, 1, 1)	(1, 1, 1, 1, 1, 1, 1, 1)
C_2	(1, 1, 1, -1, -1, -1, -1)	(1, 1, 1, 1, -1, -1, -1, -1)	(1, 1, 1, 1, -1, -1, -1, -1)
C'_2	(1, 1, 1, -1, -1, -1, -1)	(1, 1, 1, 1, 1, -1, -1, -1)	(1, 1, 1, -1, -1, -1, -1, -1)
C_8	(1, $i, -i, \nu, \nu^3, \nu^5, \nu^7$)	(1, -1, $i, -i, \nu, \nu^3, \nu^5, \nu^7$)	(1, -1, $i, -i, \nu, \nu^3, \nu^5, \nu^7$)
C'_8	(1, $i, -i, \nu, \nu^3, \nu^5, \nu^7$)	(1, -1, $i, -i, \nu, \nu^3, \nu^5, \nu^7$)	(1, -1, $i, -i, \nu, \nu^3, \nu^5, \nu^7$)
C_4	(1, -1, -1, $i, i, -i, -i$)	(1, 1, -1, -1, $i, i, -i, -i$)	(1, 1, -1, -1, $i, i, -i, -i$)
C_7	(1, $\eta, \eta^2, \eta^3, \eta^4, \eta^5, \eta^6$)	(1, 1, $\eta, \eta^2, \eta^3, \eta^4, \eta^5, \eta^6$)	(1, 1, $\eta, \eta^2, \eta^3, \eta^4, \eta^5, \eta^6$)
C_6	(1, -1, -1, μ, μ^2, μ^4, μ^5)	(1, -1, $\mu, \mu^2, \mu^2, \mu^4, \mu^4, \mu^5$)	(1, -1, $\mu, \mu, \mu^2, \mu^4, \mu^5, \mu^5$)
C_3	(1, 1, 1, $\omega, \omega, \omega^2, \omega^2$)	(1, 1, $\omega, \omega, \omega, \omega^2, \omega^2, \omega^2$)	(1, 1, $\omega, \omega, \omega, \omega^2, \omega^2, \omega^2$)

Table 10: Eigenvalues of group elements in each conjugacy class of $PGL(2; 7)$, where $\omega = e^{2\pi i/3}$, $\mu = e^{2\pi i/6}$ and $\eta = e^{2\pi i/7}$.

C'_8 square to those for C_4 , those for C_6 square to those for C_3) and also cube to those for C'_2 . These observations follow from the fact that, for example, it is impossible to choose eigenvalues for group elements in C_4 which square to those for elements in C'_2 for all irreducible representations. Note also that we have omitted the eigenvalues for matrices in the irreducible representation $\Sigma_6^{(2)'$. These eigenvalues are identical to those for $\Sigma_6^{(2)}$, except for elements in the conjugacy classes C_8, C'_8 , where for $\Sigma_6^{(2)'$ the eigenvalues for elements in C_8, C'_8 respectively are given by those for elements in C'_8, C_8 respectively in the representation $\Sigma_6^{(2)}$.

From considering the set of eigenvalues X_C for group elements in C in the representation $\gamma_1^{(i)}$, $i = 4, 8$, we see that there is a choice of $(t_1^C, t_2^C) \in X_C$ such that $\mathcal{E}_{t_1^C, t_2^C} = X_C$, for all conjugacy classes C . We present one such choice of eigenvalues (t_1^C, t_2^C) in Table 11.

Thus we obtain two non-conjugate embeddings of $PGL(2; 7)$ in G_2 . We set $\varrho_1^{(1)} = \Sigma'_7$, $\varrho_1^{(2)} = \Sigma'_1 + \Sigma_6^{(1)}$. The McKay graphs $\mathcal{G}_{PGL(2;7)}^{\varrho_1^{(1)}}$, $\mathcal{G}_{PGL(2;7)}^{\varrho_1^{(2)}}$ for $\varrho_1^{(1)}, \varrho_1^{(2)}$ are given in Figure 15, where we use the same notation as previously. Note that the graph given in [17, Figure 2] is not the McKay graph for the restriction ϱ_1 of the fundamental seven-dimensional representation of G_2 as claimed in [17], but is rather the McKay graph for the seven-dimensional representation Σ_7 , which as shown above does not define an embedding of $PGL(2; 7)$ in G_2 .

Since $\chi_{\varrho_2^{(i)}}(C) = \Phi_2(t_1^C, t_2^C)$, we see from Table 7 that the decomposition of the funda-

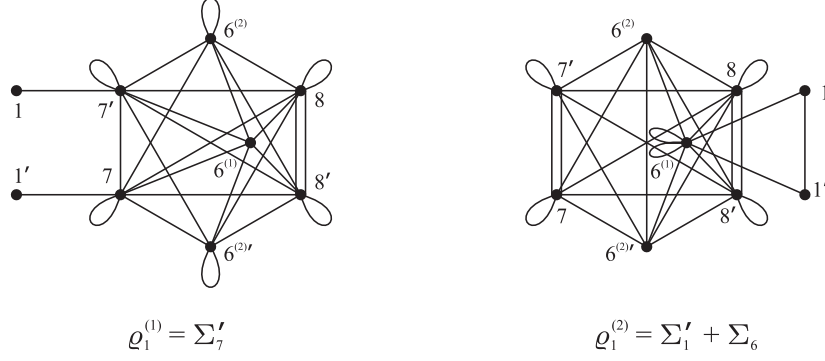


Figure 15: The McKay graphs $\mathcal{G}_{PGL(2;7)}^{\varrho_1^{(i)}}$, $i = 1, 2$.

C	$\varrho_1^{(1)}$		$\varrho_1^{(2)}$		$\chi_{\gamma_2^{(1)}}(C)$	$\chi_{\gamma_2^{(2)}}(C)$	$\chi_{\gamma_2^{(3)}}(C)$	$\chi_{\gamma_2^{(4)}}(C)$
	(t_1^C, t_2^C)	$\chi_{\varrho_2^{(1)}}(C)$	(t_1^C, t_2^C)	$\chi_{\varrho_2^{(2)}}(C)$				
C_1	(1, 1)	14	(1, 1)	14	14	14	14	14
C_2	(1, -1)	-2	(1, -1)	-2	-2	-2	2	2
C_2'	(1, -1)	-2	(1, -1)	-2	2	-2	2	-2
C_8	(i, ν^7)	0	$(-1, \nu)$	0	0	0	$\sqrt{2}$	$\sqrt{2}$
C_8'	(i, ν^5)	0	$(-1, \nu)$	0	0	0	$-\sqrt{2}$	$-\sqrt{2}$
C_4	$(-1, i)$	2	$(1, i)$	2	2	2	0	0
C_7	(η, η^5)	0	(η, η^5)	0	0	0	0	0
C_6	$(-1, \mu)$	1	$(-1, \mu)$	1	-1	1	-1	1
C_3	$(1, \omega)$	-1	$(1, \omega)$	-1	-1	-1	-1	-1

Table 11: Choice of eigenvalues (t_1^C, t_2^C) for $\varrho_1^{(i)}$, $i = 1, 2$, and corresponding values of $\chi_{\varrho_2^{(i)}}(C)$.

mental fourteen-dimensional representation into irreducible representations of $PGL(2; 7)$ is given by $\varrho_2 = \gamma_2^{(2)} = \Sigma_6^{(1)} + \Sigma_8'$ for both $i = 1, 2$. The values of $x^{(i)} = \chi_{\varrho_1^{(i)}}(C) \in [-2, 7]$, $y = \chi_{\varrho_2}(C) \in [-2, 14]$ for $PGL(2; 7)$ are given in Table 12.

Then from (13) and Tables 9, 12 and 2, we see that

$$\begin{aligned} \varsigma_{m,n} = & \frac{1}{336} \Omega^W(0, 0) + \frac{21 + 28}{336} \Omega^W(0, 1/2) + \frac{56}{336} \Omega^W(0, 1/3) \\ & + \frac{56}{336} \Omega^W(1/6, 1/2) + \frac{48}{336} \Omega^W(1/7, 3/7) + \frac{42}{336} \Omega' + \frac{42 + 42}{336} \Omega'', \end{aligned}$$

where $\Omega^W(\theta_1, \theta_2)$ is as in Section 4.1 and (Ω', Ω'') is $(\Omega^W(1/4, 1/2), \Omega^W(1/8, 3/8))$ for $\varrho_1^{(1)}$, and $(\Omega^W(0, 1/4), \Omega^W(1/8, 1/2))$ for $\varrho_1^{(2)}$. Then we obtain:

Theorem 4.4. The joint spectral measure (over \mathbb{T}^2) for the non-conjugate embeddings

C	C_1	C_2	C'_2	C_8	C'_8	C_4	C_7	C_6	C_3
$\chi_{\varrho_1^{(1)}}(C) \in [-2, 7]$	7	-1	-1	1	1	-1	0	-1	1
$\chi_{\varrho_1^{(2)}}(C) \in [-2, 7]$	7	-1	-1	-1	-1	3	0	-1	1
$\chi_{\varrho_2}(C) \in [-2, 14]$	14	-2	-2	0	0	2	0	1	-1

Table 12: $\chi_{\varrho_j}(C)$ for group $PGL(2; 7)$, $j = 1, 2$.

C	C_1	C_3	C_9	C'_9	C''_9	C_2	C_7	C'_7	C''_7
$ C $	1	56	56	56	56	63	72	72	72
Σ_1	1	1	1	1	1	1	1	1	1
$\Sigma_7^{(1)}$	7	1	$-p$	$-q$	$p+q$	-1	0	0	0
$\Sigma_7^{(1)'}$	7	1	$p+q$	$-p$	$-q$	-1	0	0	0
$\Sigma_7^{(1)''}$	7	1	$-q$	$p+q$	$-p$	-1	0	0	0
$\Sigma_7^{(2)}$	7	-2	1	1	1	-1	0	0	0
Σ_8	8	-1	-1	-1	-1	0	1	1	1
Σ_9	9	0	0	0	0	1	$2 \cos(2\pi/7)$	$2 \cos(4\pi/7)$	$2 \cos(6\pi/7)$
Σ'_9	9	0	0	0	0	1	$2 \cos(4\pi/7)$	$2 \cos(6\pi/7)$	$2 \cos(2\pi/7)$
Σ''_9	9	0	0	0	0	1	$2 \cos(6\pi/7)$	$2 \cos(2\pi/7)$	$2 \cos(4\pi/7)$

Table 13: Character table for $PSL(2; 8)$, where $p = 2 \cos(4\pi/9)$, $q = 2 \cos(8\pi/9)$.

of $PGL(2; 7)$ into the fundamental representations of G_2 is

$$\begin{aligned}
d\varepsilon = & \frac{1}{1344\pi^4} J^2 d_7 \times d_7 + \frac{1}{1152\pi^4} J^2 d_6 \times d_6 + \frac{1}{48} K' d_4 \times d_4 + \frac{1}{4} d_3 \times d_3 \\
& + \frac{1}{36} \beta d_2 \times d_2 - \frac{23}{504} d_1 \times d_1 + \frac{1}{48} \sum_{g \in D_{12}} \delta_{g(\mu, t)},
\end{aligned} \tag{34}$$

for μ a primitive 8-th root of unity, where for the embedding of $PGL(2; 7)$ in G_2 given by $\varrho_1^{(1)} = \Sigma'_7$ we have $K' = 16 - K$, $\beta = 4$ and $t = \mu^3$, whilst for the embedding given by $\varrho_1^{(2)} = \Sigma'_1 + \Sigma_6$ we have $K' = K$, $\beta = 7$ and $t = -1$, with K as in Theorem 4.1, and where d_m is the uniform measure over m^{th} roots of unity and δ_x is the Dirac measure at the point x .

4.4 Group $PSL(2; 8)$

The subgroup $PSL(2; 8)$ of G_2 is an irreducible primitive group of order 504. It has nine irreducible representations, all real, five of which have dimension less than or equal to 7. The character table for $PSL(2; 8)$ is given in Table 13 [18] (the orders of the elements in each conjugacy class can be obtained from [27]). The spectral measure for the first three seven-dimensional representations are equal, since the conjugacy classes C_9 , C'_9 , C''_9 each have the same order.

We thus have four candidates for the restriction ϱ_1 of the fundamental representation ρ_1 of G_2 to $PSL(2; 8)$. The Kronecker squares of $\Sigma_7^{(1)}$, $\Sigma_7^{(2)}$ decompose into irreducibles

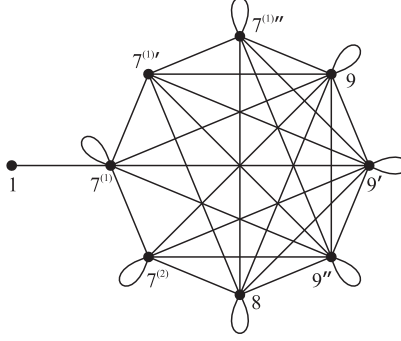


Figure 16: The McKay graph $\mathcal{G}_{PSL(2;8)}^{\varrho_1^{(1)}}$.

C	C_1	C_3	C_9	C_9'	C_9''	C_2	C_7	C_7'	C_7''
$\chi_{\varrho_1^{(j)}}(C) \in [-2, 7]$	7	1	$-p$	$-q$	$p+q$	-1	0	0	0
$\chi_{\varrho_2}(C) \in [-2, 14]$	14	-1	$p+q+1$	$1-p$	$1-q$	-2	0	0	0

Table 14: $\chi_{\varrho_1^{(j)}}(C)$ and $\chi_{\varrho_2}(C)$ for group $PSL(2;8)$, $j \in \{1, 2, 3\}$.

as

$$(\Sigma_7^{(1)})^2 = \text{id} + \Sigma_7^{(1)} + \Sigma_7^{(1)'} + \Sigma_7^{(2)} + \Sigma_9 + \Sigma_9' + \Sigma_9'', \quad (\Sigma_7^{(2)})^2 = \text{id} + \Sigma_7^{(1)} + \Sigma_7^{(1)'} + \Sigma_7^{(1)''} + \Sigma_9 + \Sigma_9' + \Sigma_9'',$$

where $\text{id} = \Sigma_1$. The irreducible representation $\Sigma_7^{(2)}$ does not appear in the decomposition of its Kronecker square into irreducibles, and therefore does not give an embedding of $PSL(2;8)$ into G_2 .

Thus we obtain that there are precisely three seven-dimensional representations give non-conjugate embeddings of $PSL(2;8)$ into G_2 . Let $\varrho_1^{(1)} = \Sigma_7^{(1)}$, $\varrho_1^{(2)} = \Sigma_7^{(1)'}$ and $\varrho_1^{(3)} = \Sigma_7^{(1)''}$. The McKay graph for $\varrho_1^{(1)}$ is given in Figure 16, where we use the same notation as previously. The McKay graphs for $\varrho_1^{(j)}$, $j = 2, 3$, are similar. Note that the graph given in [17, Figure 2] is not the McKay graph for a restriction $\varrho_1^{(j)}$, $j \in \{1, 2, 3\}$, of the fundamental seven-dimensional representation of G_2 as claimed in [17], but is rather the McKay graph for the seven-dimensional representation $\Sigma_7^{(2)}$.

From the decomposition of the Kronecker square of $\varrho_1^{(j)}$ and dimension considerations there is only one possibility for the fourteen-dimensional representation ϱ_2 , that is, $\varrho_2 = \Sigma_7^{(1)'} + \Sigma_7^{(2)}$. The values of $x = \chi_{\varrho_1^{(j)}}(C) \in [-2, 7]$, $j \in \{1, 2, 3\}$, and $y = \chi_{\varrho_2}(C) \in [-2, 14]$ for $PSL(2;8)$ are given in Table 14.

Then from (13) and Tables 13 and 2, we see that

$$\begin{aligned} \varsigma_{m,n} = & \frac{1}{504} \Omega^W(0,0) + \frac{56}{504} \Omega^W(0,1/3) + \frac{63}{504} \Omega^W(0,1/2) + \frac{72+72+72}{504} \Omega^W(1/7,3/7) \\ & + \frac{56}{504} \Omega^W(1/9,4/9) + \frac{56}{504} \Omega^W(1/9,1/3) + \frac{56}{504} \Omega^W(2/9,5/9), \end{aligned}$$

where $\Omega^W(\theta_1, \theta_2)$ is as in Section 4.1. Then we obtain:

C	C_1	C_{13}	C'_{13}	C_2	C_7	C'_7	C''_7	C_6	C_3
$ C $	1	84	84	91	156	156	156	182	182
Σ_1	1	1	1	1	1	1	1	1	1
Σ_7	7	p_-	p_+	-1	0	0	0	-1	1
Σ'_7	7	p_+	p_-	-1	0	0	0	-1	1
Σ_{12}	12	-1	-1	0	r_1	r_2	r_3	0	0
Σ'_{12}	12	-1	-1	0	r_2	r_3	r_1	0	0
Σ''_{12}	12	-1	-1	0	r_3	r_1	r_2	0	0
Σ_{13}	13	0	0	1	-1	-1	-1	1	1
Σ_{14}	14	1	1	-2	0	0	0	1	-1
Σ'_{14}	14	1	1	2	0	0	0	-1	-1

Table 15: Character table for $PSL(2; 13)$, where $p_{\pm} = (1 \pm \sqrt{13})/2$, $r_j = -2 \cos(2j\pi/7)$.

Theorem 4.5. The joint spectral measure (over \mathbb{T}^2) for all embeddings of $PSL(2; 8)$ into the fundamental representations of G_2 is

$$\begin{aligned}
d\varepsilon = & \frac{1}{448\pi^4} J^2 d_7 \times d_7 + \frac{1}{6} d_3 \times d_3 + \frac{1}{6} d_2 \times d_2 - \frac{5}{126} d_1 \times d_1 - \frac{1}{18} d^{(1)} \\
& + \frac{1}{864\pi^4} a_3^{-1} J^2 d^{((9))} + \frac{1}{864\pi^4} a_1^{-1} J^2 d^{((9/2))} + \frac{1}{864\pi^4} a_2^{-1} J^2 d^{((9/4))},
\end{aligned} \tag{35}$$

where $d^{((n))}$ is as in Definition 2.1, d_m is the uniform measure over m^{th} roots of unity and $d^{(k+4)}$ is the uniform measure on the points in C_k^W .

4.5 Group $PSL(2; 13)$

The subgroup $PSL(2; 13)$ of G_2 is an irreducible primitive group of order 1092. It has nine irreducible representations, all real, and its character table is given in Table 15 [2].

Only three of the irreducible representation have dimension less than or equal to 7. These are the identity representation, and two seven-dimensional representations Σ_7 , Σ'_7 whose character values only differ on the two conjugacy classes C_{13} , C'_{13} whose elements have order 13. Here $\chi_{\Sigma_7}(g) = p, q$ for $g \in C_{13}, C'_{13}$ respectively, whilst $\chi_{\Sigma'_7}(g) = q, p$ respectively, where $p = (1 + \sqrt{13})/2 = 1 + \zeta + \zeta^3 + \zeta^4 + \zeta^9 + \zeta^{10} + \zeta^{12}$, $q = (1 - \sqrt{13})/2 = 1 + \zeta^2 + \zeta^5 + \zeta^6 + \zeta^7 + \zeta^8 + \zeta^{11}$, for $\zeta = e^{2\pi i/13}$. Thus the spectral measure for both seven-dimensional representations are equal, since C_{13} , C'_{13} both have the same order.

There are thus two non-conjugate embeddings $\varrho_1^{(1)} := \Sigma_7$ and $\varrho_1^{(2)} := \Sigma'_7$ of $PSL(2; 13)$ in G_2 . The McKay graph for the fundamental seven-dimensional representation $\varrho_1^{(j)}$, $j \in \{1, 2\}$, is given in [17, Figure 2], and is reproduced (twice) in Figure 17 for completeness, where we use the same notation as previously. The figure on the right hand side illustrates the resemblance of $\mathcal{G}_{PSL(2;13)}^{\varrho_1^{(j)}}$ with the McKay graph of G_2 itself.

The set X_C of eigenvalues of group elements in each conjugacy class for $\varrho_1^{(1)} = \Sigma_7$ are given in Table 16, along with a choice of $(t_1^C, t_2^C) \in X_C$ such that $\mathcal{E}_{t_1^C, t_2^C} = X_C$. The eigenvalues for $\varrho_1^{(2)}$ are similar. The decomposition of the Kronecker square of $\varrho_1^{(j)}$,

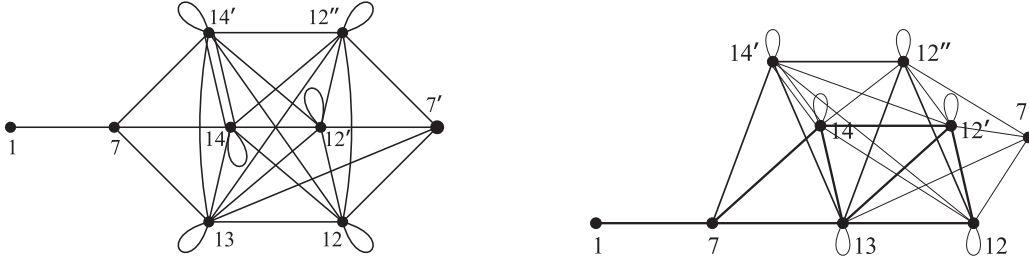


Figure 17: Two presentations of the McKay graph $\mathcal{G}_{PSL(2;13)}^{\varrho_1^{(j)}}$, $j \in \{1, 2\}$.

C	X_C	(t_1^C, t_2^C)	$\chi_{\varrho_1^{(1)}}(C)$	$\chi_{\varrho_2}(C)$	$\chi_{\Sigma_{14}}(C)$	$\chi_{\Sigma'_{14}}(C)$
C_1	$(1, 1, 1, 1, 1, 1, 1)$	$(1, 1)$	7	14	14	14
C_{13}	$(1, \zeta, \zeta^3, \zeta^4, \zeta^9, \zeta^{10}, \zeta^{12})$	(ζ, ζ^4)	$(1 + \sqrt{13})/2$	1	1	1
C'_{13}	$(1, \zeta^2, \zeta^5, \zeta^6, \zeta^7, \zeta^8, \zeta^{11})$	(ζ^2, ζ^7)	$(1 - \sqrt{13})/2$	1	1	1
C_2	$(1, 1, 1, -1, -1, -1, -1)$	$(1, -1)$	-1	-2	-2	2
C_7	$(1, \eta, \eta^2, \eta^3, \eta^4, \eta^5, \eta^6)$	(η, η^5)	0	0	0	0
C'_7	$(1, \eta, \eta^2, \eta^3, \eta^4, \eta^5, \eta^6)$	(η, η^5)	0	0	0	0
C''_7	$(1, \eta, \eta^2, \eta^3, \eta^4, \eta^5, \eta^6)$	(η, η^5)	0	0	0	0
C_6	$(1, -1, -1, \mu, \mu^2, \mu^4, \mu^5)$	$(-1, \mu)$	-1	1	1	-1
C_3	$(1, 1, 1, \omega, \omega, \omega^2, \omega^2)$	$(1, \omega)$	1	-1	-1	-1

Table 16: Choice of eigenvalues (t_1^C, t_2^C) for $\varrho_1^{(1)}$ and corresponding values of $\chi_{\varrho_2}(C)$, where $\omega = e^{2\pi i/3}$, $\mu = e^{2\pi i/6}$, $\eta = e^{2\pi i/7}$ and $\zeta = e^{2\pi i/13}$.

$j \in \{1, 2\}$, into irreducibles is given by

$$\left(\varrho_1^{(j)}\right)^2 = \text{id} + \varrho_1^{(j)} + \Sigma_{13} + \Sigma_{14} + \Sigma'_{14},$$

where as before the notation Σ_n denotes an irreducible representation of $PSL(2; 13)$ of dimension n . From dimension considerations there are thus two candidates, Σ_{14} and Σ'_{14} , for the fourteen-dimensional representation ϱ_2 , which are given by the two fourteen-dimensional irreducible representations. However, since $\chi_{\varrho_2}(C) = \Phi_2(t_1^C, t_2^C)$, we see from Table 16 that the decomposition of the fundamental fourteen-dimensional representation into irreducible representations of $PSL(2; 13)$ is given by $\varrho_2 = \Sigma_{14}$, not Σ'_{14} .

Then from (13) and Tables 16 and 2, we see that

$$\begin{aligned} \varsigma_{m,n} = & \frac{1}{1092} \Omega^W(0, 0) + \frac{91}{1092} \Omega^W(0, 1/2) + \frac{182}{1092} \Omega^W(0, 1/3) + \frac{182}{1092} \Omega^W(1/6, 1/2) \\ & + \frac{156 + 156 + 156}{1092} \Omega^W(1/7, 3/7) + \frac{84}{1092} \Omega^W(1/13, 4/13) + \frac{84}{1092} \Omega^W(2/13, 7/13), \end{aligned}$$

where $\Omega^W(\theta_1, \theta_2)$ is as in Section 4.1. Then we obtain:

Theorem 4.6. The joint spectral measure (over \mathbb{T}^2) for all embeddings of $PSL(2; 13)$

C	C_1	C_3	C_2	C_4	C'_4	C''_4	C_{12}	C'_{12}	C_6	C'_3	C_8	C'_8	C_7	C'_7
$ C $	1	56	63	63	63	378	504	504	504	672	756	756	864	864
Σ_1	1	1	1	1	1	1	1	1	1	1	1	1	1	1
Σ_6	6	-3	-2	-2	-2	2	1	1	1	0	0	0	-1	-1
Σ_7	7	-2	3	$-1 - 2i$	$-1 + 2i$	1	$-1 - i$	$-1 + i$	0	1	$-i$	i	0	0
Σ_7^*	7	-2	3	$-1 + 2i$	$-1 - 2i$	1	$-1 + i$	$-1 - i$	0	1	i	$-i$	0	0
Σ'_7	7	-2	-1	3	3	-1	0	0	2	1	-1	-1	0	0
Σ_{14}	14	5	-2	2	2	2	-1	-1	1	-1	0	0	0	0
Σ_{21}	21	3	1	$-3 - 2i$	$-3 + 2i$	-1	$-i$	i	1	0	i	$-i$	0	0
Σ_{21}^*	21	3	1	$-3 + 2i$	$-3 - 2i$	-1	i	$-i$	1	0	$-i$	i	0	0
Σ'_{21}	21	3	5	1	1	1	1	1	-1	0	-1	-1	0	0
Σ_{27}	27	0	3	3	3	-1	0	0	0	0	1	1	-1	-1
Σ_{28}	28	1	-4	$-4i$	$4i$	0	i	$-i$	-1	1	0	0	0	0
Σ_{28}^*	28	1	-4	$4i$	$-4i$	0	$-i$	i	-1	1	0	0	0	0
Σ_{32}	32	-4	0	0	0	0	0	0	0	-1	0	0	$\frac{1+i\sqrt{7}}{2}$	$\frac{1-i\sqrt{7}}{2}$
Σ_{32}^*	32	-4	0	0	0	0	0	0	0	-1	0	0	$\frac{1-i\sqrt{7}}{2}$	$\frac{1+i\sqrt{7}}{2}$

Table 17: Character table for $PU(3; 3)$.

into the fundamental representations of G_2 is

$$\begin{aligned}
d\varepsilon = & \frac{1}{448\pi^4} J^2 d_7 \times d_7 + \frac{1}{1152\pi^4} J^2 d_6 \times d_6 + \frac{1}{4} d_3 \times d_3 + \frac{1}{9} d_2 \times d_2 \\
& - \frac{22}{819} d_1 \times d_1 - \frac{1}{12} d^{(1)} + \frac{1}{156} \sum_{g \in D_{12}} (\delta_{g(\zeta, \zeta^4)} + \delta_{g(\zeta^2, \zeta^7)})
\end{aligned} \tag{36}$$

for ζ a primitive 13-th root of unity, where d_m is the uniform measure over m^{th} roots of unity, $d^{(k+4)}$ is the uniform measure on the points in C_k^W , and δ_x is the Dirac measure at the point x .

4.6 Group $PU(3; 3) \cong G_2(2)'$

The subgroup $PU(3; 3)$ of G_2 is an irreducible primitive group of order 6048. It has fourteen irreducible representations (seven real representations, one quaternionic representation Σ_6 , and three pairs of complex conjugate representations) and its character table is given in Table 17 [5].

There are two non-conjugate real seven-dimensional representations, $\gamma_1^{(1)} = \Sigma_1 + \Sigma_6$ and $\gamma_1^{(2)} = \Sigma'_7$. These both satisfy the condition that $\gamma_1^{(i)}$ appears in the decomposition of $(\gamma_1^{(i)})^2$. We thus consider the eigenvalues of the representation matrices to determine which of the $\gamma_1^{(i)}$ are embeddings of $PU(3; 3)$ in G_2 . These eigenvalues are given in Table 18 for the representations Σ_1 , Σ_6 and Σ'_7 . As described in Section 3, these eigenvalues can be determined from the character table of $PU(3; 3)$. The additional information that is needed is to note that the eigenvalues for group elements in C_4 , C'_4 and C''_4 all square to those for elements in C_2 , those for C_8 , C'_8 square to those for C_4 , C'_4 respectively, those

	Σ_1	Σ_6	Σ'_7	(t_1^C, t_2^C)
C_1	1	(1, 1, 1, 1, 1, 1)	(1, 1, 1, 1, 1, 1, 1)	(1, 1)
C_3	1	$(\omega, \omega, \omega, \omega^2, \omega^2, \omega^2)$	$(1, \omega, \omega, \omega, \omega^2, \omega^2, \omega^2)$	(ω, ω^2)
C_2	1	(1, 1, -1, -1, -1, -1)	(1, 1, 1, -1, -1, -1, -1)	(1, -1)
C_4	1	$(-1, -1, i, i, -i, -i)$	$(1, 1, 1, i, i, -i, -i)$	(1, i)
C'_4	1	$(-1, -1, i, i, -i, -i)$	$(1, 1, 1, i, i, -i, -i)$	(1, i)
C''_4	1	$(1, 1, i, i, -i, -i)$	$(1, -1, -1, i, i, -i, -i)$	$(-1, i)$
C_{12}	1	$(\xi, \xi^2, \xi^5, \xi^7, \xi^{10}, \xi^{11})$	$(1, \xi, \xi^4, \xi^5, \xi^7, \xi^8, \xi^{11})$	(ξ, ξ^5)
C'_{12}	1	$(\xi, \xi^2, \xi^5, \xi^7, \xi^{10}, \xi^{11})$	$(1, \xi, \xi^4, \xi^5, \xi^7, \xi^8, \xi^{11})$	(ξ, ξ^5)
C_6	1	$(\mu, \mu, \mu^2, \mu^4, \mu^5, \mu^5)$	$(1, \mu, \mu, \mu^2, \mu^4, \mu^5, \mu^5)$	(μ, μ^2)
C'_3	1	$(1, 1, \omega, \omega, \omega^2, \omega^2)$	$(1, 1, 1, \omega, \omega, \omega^2, \omega^2)$	(1, ω)
C_8	1	$(i, -i, \nu, \nu^3, \nu^5, \nu^7)$	$(1, -1, -1, \nu, \nu^3, \nu^5, \nu^7)$	$(-1, \nu)$
C'_8	1	$(i, -i, \nu, \nu^3, \nu^5, \nu^7)$	$(1, -1, -1, \nu, \nu^3, \nu^5, \nu^7)$	$(-1, \nu)$
C_7	1	$(\eta, \eta^2, \eta^3, \eta^4, \eta^5, \eta^6)$	$(1, \eta, \eta^2, \eta^3, \eta^4, \eta^5, \eta^6)$	(η, η^5)
C'_7	1	$(\eta, \eta^2, \eta^3, \eta^4, \eta^5, \eta^6)$	$(1, \eta, \eta^2, \eta^3, \eta^4, \eta^5, \eta^6)$	(η, η^5)

Table 18: Eigenvalues of group elements in each conjugacy class of $PU(3;3)$ for the irreducible representations Σ_1 , Σ_6 and Σ'_7 , where $\omega = e^{2\pi i/3}$, $\mu = e^{2\pi i/6}$, $\eta = e^{2\pi i/7}$, $\nu = e^{2\pi i/8}$ and $\xi = e^{2\pi i/12}$

for C_6 square to those for C_3 and also cube to those for C_2 , those for C_{12} square to those for C_6 and cube to those for C_4 whilst those for C'_{12} also square to those for C_6 but cube to those for C'_4 (see [5]).

From considering the set of eigenvalues X_C for group elements in C in the representation $\gamma_1^{(i)}$, we see that there is no choice of $(t_1^C, t_2^C) \in X_C$ such that $\mathcal{E}_{t_1^C, t_2^C} = X_C$ for $i = 1$ when $C = C_{12}, C'_{12}$. However, such a choice does exist for all conjugacy classes C for $i = 2$, thus we have $\varrho_1 = \gamma_1^{(2)}$. We present one such choice of eigenvalues (t_1^C, t_2^C) in the final column of Table 18.

Thus, up to conjugacy, there is only one embedding of $PU(3;3)$ in G_2 . The McKay graph $\mathcal{G}_{PU(3;3)}^{\varrho_1}$ is given in [17, Figure 2], and we reproduce it here in Figure 18 for completeness.

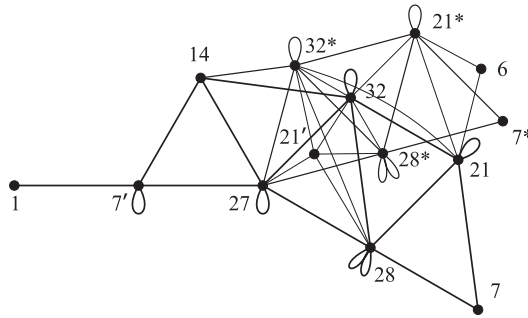


Figure 18: The McKay graphs $\mathcal{G}_{PU(3;3)}^{\varrho_1}$.

The decomposition of the Kronecker square of ϱ_1 into irreducibles is given by

$$\varrho_1^2 = \text{id} + \varrho_1 + \Sigma_{14} + \Sigma_{27},$$

where $\text{id} = \Sigma_1$. Thus the fourteen-dimensional representation ϱ_2 must be Σ_{14} , and we note that $\chi_{\varrho_2}(C) = \Phi_2(t_1^C, t_2^C)$ as required.

Then from (13) and Tables 17 and 2, we see that

$$\begin{aligned} \varsigma_{m,n} = & \frac{1}{6048} \Omega^W(0, 0) + \frac{56}{6048} \Omega^W(1/3, 2/3) + \frac{63}{6048} \Omega^W(0, 1/2) + \frac{672}{6048} \Omega^W(0, 1/3) \\ & + \frac{63+63}{6048} \Omega^W(0, 1/4) + \frac{378}{6048} \Omega^W(1/4, 1/2) + \frac{504}{6048} \Omega^W(1/6, 1/3) \\ & + \frac{864+864}{6048} \Omega^W(1/7, 3/7) + \frac{756+756}{6048} \Omega^W(1/8, 1/2) + \frac{504+504}{6048} \Omega^W(1/12, 5/12), \end{aligned}$$

where $\Omega^W(\theta_1, \theta_2)$ is as in Section 4.1. Then we obtain:

Theorem 4.7. The joint spectral measure (over \mathbb{T}^2) for all embeddings of $PU(3; 3)$ into the fundamental representations of G_2 is

$$\begin{aligned} d\varepsilon = & \frac{1}{672\pi^4} J^2 d_7 \times d_7 + \frac{1}{144} (24 - K) d_4 \times d_4 + \frac{1}{6} d_3 \times d_3 - \frac{1}{12} d_2 \times d_2 \\ & + \frac{1}{168} d_1 \times d_1 + \frac{1}{1152\pi^4} J^2 d^{(4)} + \frac{1}{6} d^{(2)} - \frac{1}{12} d^{(1)} + \frac{1}{48} \sum_{g \in D_{12}} \delta_{g(\mu, -1)}, \end{aligned} \quad (37)$$

for μ primitive 8-th root of unity, where $K(\theta_1, \theta_2) = (\sin(2\pi(\theta_1 + \theta_2)) - \sin(2\pi(2\theta_1 - \theta_2)) - \sin(2\pi(2\theta_2 - \theta_1)))^2$, d_m is the uniform measure over m^{th} roots of unity, $d^{(k+4)}$ is the uniform measure on the points in C_k^W , and δ_x is the Dirac measure at the point x .

4.7 Group $G_2(2)$

The subgroup $G_2(2) = G_2(\mathbb{F}_2)$ of $G_2 = G_2(\mathbb{C})$ is an irreducible primitive group of order 12096. It is the group G_2 defined over the Galois field \mathbb{F}_2 . It has sixteen irreducible representations (fourteen real and two complex conjugate representations), and its character table is given in [17] (the orders of the elements in each conjugacy class can be obtained from [24]).

The group $G_2(2)$ has only two seven-dimensional real representations, both of which are irreducible. Of these two, only Σ_7 has character values in $[-2, 7]$ for all $g \in G_2(2)$, making it the restriction ϱ_1 of the seven-dimensional fundamental representation ρ_1 of G_2 to $G_2(2)$.

Thus, up to conjugacy, there is only one embedding of $G_2(2)$ in G_2 . The McKay graph for ϱ_1 is given in Figure 19. This graph is a \mathbb{Z}_2 -orbifold of the McKay graph $\mathcal{G}_{PU(3;3)}^{\varrho_1}$ for $PU(3; 3)$. Note that the McKay graph given in [17, Figure 2] is not that for ϱ_1 , but rather for the other irreducible seven-dimensional representation Σ_7' of $G_2(2)$.

The decomposition of the Kronecker square of ϱ_1 into irreducibles is given by

$$\varrho_1^2 = \text{id} + \varrho_1 + \Sigma_{14} + \Sigma_{27},$$

thus the fourteen-dimensional representation ϱ_2 is given by the irreducible representation Σ_{14} . We note that the eigenvalues of the representation matrices for C_4, C_4', C_4'' all square

C	C_1	C_3	C_2	C_4	C'_2	C'_4	C''_4	C_6	C'_3	C_{12}	C'_{12}	C''_{12}	C_8	C'_8	C_7	C'_6
$ C $	1	56	63	126	252	252	378	504	672	1008	1008	1008	1512	1512	1728	2016
Σ_1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1
Σ'_1	1	1	1	1	-1	-1	1	1	1	-1	-1	1	-1	1	1	-1
Σ_6	6	-3	-2	-2	0	0	2	1	0	$i\sqrt{3}$	$-i\sqrt{3}$	1	0	0	-1	0
Σ^*_6	6	-3	-2	-2	0	0	2	1	0	$-i\sqrt{3}$	$i\sqrt{3}$	1	0	0	-1	0
Σ_7	1	-2	-1	3	-1	3	-1	2	1	0	0	0	1	-1	0	-1
Σ'_7	1	-2	-1	3	1	-3	-1	2	1	0	0	0	-1	-1	0	1
Σ_{14}	14	5	-2	2	-2	2	2	1	-1	-1	-1	-1	0	0	0	1
Σ'_{14}	14	-4	6	-2	0	0	2	0	2	0	0	-2	0	0	0	0
Σ''_{14}	14	5	-2	2	2	-2	2	1	-1	1	1	-1	0	0	0	-1
Σ_{21}	21	3	5	1	3	-1	1	-1	0	-1	-1	1	1	-1	0	0
Σ'_{21}	21	3	5	1	-3	1	1	-1	0	1	1	1	-1	-1	0	0
Σ_{27}	27	0	3	3	3	3	-1	0	0	0	0	0	-1	1	-1	0
Σ'_{27}	27	0	3	3	-3	-3	-1	0	0	0	0	0	1	1	-1	0
Σ_{42}	42	6	2	-6	0	0	-2	2	0	0	0	0	0	0	0	0
Σ_{56}	56	2	-8	0	0	0	0	-2	2	0	0	0	0	0	0	0
Σ_{64}	64	-8	0	0	0	0	0	0	-2	0	0	0	0	0	1	0

Table 19: Character table for $G_2(2)$.

to those for C_2 , those for C_6 square to those for C_3 , those for C'_6 square to those for C'_3 , those for C_8 square to those for C_4 , those for C'_8 square to those for C'_4 , and those for $C_{12}, C'_{12}, C''_{12}$ all square to those for C_6 .

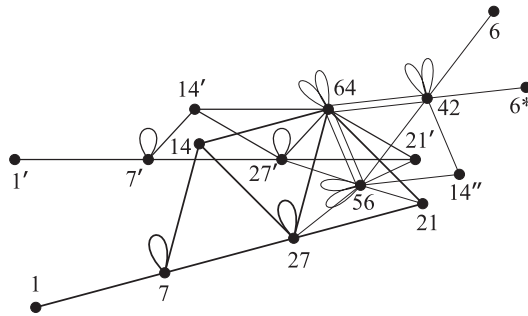


Figure 19: The McKay graphs $\mathcal{G}_{G_2(2)}^{e1}$.

Then from (13) and Tables 19 and 2, we see that

$$\begin{aligned}
\varsigma_{m,n} = & \frac{1}{12096} \Omega^W(0,0) + \frac{56}{12096} \Omega^W(1/3,2/3) + \frac{63+252}{12096} \Omega^W(0,1/2) + \frac{672}{12096} \Omega^W(0,1/3) \\
& + \frac{126+252}{12096} \Omega^W(0,1/4) + \frac{378}{12096} \Omega^W(1/4,1/2) + \frac{504}{12096} \Omega^W(1/6,1/3) \\
& + \frac{1728}{12096} \Omega^W(1/7,3/7) + \frac{1512}{12096} \Omega^W(1/8,1/2) + \frac{1512}{12096} \Omega^W(1/8,3/8) \\
& + \frac{1008+1008+1008}{12096} \Omega^W(1/12,5/12) + \frac{2016}{12096} \Omega^W(1/6,1/2),
\end{aligned}$$

where $\Omega^W(\theta_1, \theta_2)$ is as in Section 4.1. Thus we obtain:

Theorem 4.8. The joint spectral measure (over \mathbb{T}^2) for all embeddings of $G_2(2)$ into the fundamental representations of G_2 is

$$\begin{aligned}
d\varepsilon = & \frac{1}{768\pi^4} J^2 d_8 \times d_8 + \frac{1}{1344\pi^4} J^2 d_7 \times d_7 + \frac{1}{1152\pi^4} J^2 d_6 \times d_6 + \frac{1}{12} d_4 \times d_4 \\
& + \frac{1}{12} d_3 \times d_3 - \frac{1}{72} d_2 \times d_2 - \frac{1}{252} d_1 \times d_1 + \frac{1}{768\pi^4} J^2 d^{(4)} + \frac{1}{12} d^{(2)} - \frac{1}{24} d^{(1)},
\end{aligned} \tag{38}$$

where d_m is the uniform measure over m^{th} roots of unity and $d^{(k+4)}$ is the uniform measure on the points in C_k^W .

4.8 Conclusion and Outlook

The spectral measures studied here encode the character table of the groups. In the context of quantum subgroups, as discussed in the Introduction and Section 2.3, the spectral measure provides the appropriate generalisation of the character table. Whilst the spectral measures for known nimrep graphs associated to the G_2 modular invariant partition functions were studied in [10], the classification of quantum subgroups of G_2 is still open, including the classification of all nimrep graphs for G_2 . And beyond G_2 , there are many other rank two Lie groups where little is known about their quantum subgroups, including the semisimple, connected, simply-connected, compact group $SU(2) \times SU(2)$, as well as other connected (but not necessarily simply-connected) compact groups, e.g. $U(2)$, $PSU(3)$ or $\mathbb{T} \times SO(3)$ [12].

The spectral measures for rank two Lie groups and their quantum subgroups are naturally given as measures over \mathbb{T}^2 . The spectral measures presented in this paper for finite subgroups of G_2 have been given over \mathbb{T}^2 , however by transporting the measure on a suitable (possibly noncommutative) manifold the picture might simplify. This would be a valuable direction for future research in this area.

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References

- [1] T. Banica and D. Bisch, Spectral measures of small index principal graphs, *Comm. Math. Phys.* **269** (2007), 259–281.
- [2] A.M. Cohen, Finite subgroups of Lie groups, Lecture notes from Leiden Summer Course on Lie groups (June 30, 1998).
- [3] A.M. Cohen and D.B. Wales, Finite subgroups of $G_2(\mathbb{C})$, *Comm. Algebra* **11** (1983) 441–459.
- [4] M.J. Collins, *Representations and characters of finite groups*, Cambridge Studies in Advanced Mathematics **22**. Cambridge University Press, Cambridge, 1990.
- [5] J.H. Conway, R.T. Curtis, S.P. Norton, R.A. Parker and R.A. Wilson, *Atlas of finite groups*. Oxford University Press, Eynsham, 1985.
- [6] H.S.M. Coxeter, Integral Cayley numbers, *Duke Math. J.* **13** (1946), 561–578.
- [7] D. E. Evans and Y. Kawahigashi, *Quantum symmetries on operator algebras*, Oxford Mathematical Monographs. The Clarendon Press Oxford University Press, New York, 1998. Oxford Science Publications.
- [8] D.E. Evans and M. Pugh, Spectral measures and generating series for nimrep graphs in subfactor theory, *Comm. Math. Phys.* **295** (2010), 363–413.
- [9] D. E. Evans and M. Pugh, Spectral measures and generating series for nimrep graphs in subfactor theory II: $SU(3)$, *Comm. Math. Phys.* **301** (2011), 771–809.
- [10] D.E. Evans and M. Pugh, Spectral measures for G_2 , *Comm. Math. Phys.* **337** (2015), 1161–1197.
- [11] D.E. Evans and M. Pugh, Spectral measures for $Sp(2)$. Preprint, arXiv:1404.1912 [math.OA].
- [12] D.E. Evans and M. Pugh, Spectral measures associated to rank two Lie groups and finite subgroups of $GL(2, \mathbb{Z})$, *Comm. Math. Phys.* **343** (2016), 811–850.
- [13] F. M. Goodman, P. de la Harpe and V. F. R. Jones, *Coxeter graphs and towers of algebras* MSRI Publications, 14, Springer-Verlag, New York, 1989.
- [14] R.L. Griess, Basic conjugacy theorems for G_2 , *Invent. Math.* **121** (1995), 257–277.
- [15] A. Guionnet, Uses of free probability in random matrix theory, XVIth International Congress on Mathematical Physics (2010) pp. 106–122.
- [16] R. Haag and D. Kastler, An algebraic approach to quantum field theory, *Journ. Math. Phys.* **5** (1964), 848–861.
- [17] Y.-H. He, G_2 quivers, *J. High Energy Phys.* **2003**, 023, 14 pp.
- [18] G. James and M. Liebeck, *Representations and characters of groups*, Second Edition, Cambridge University Press, New York, 2001.
- [19] C. Jones, Quantum G_2 categories have property (T), *Int. J. Math.* **27** (2016), 1650015, 23 pp.
- [20] V. F. R. Jones, Planar algebras. I, *New Zealand J. Math.* (to appear).
- [21] T. Kawai, On the structure of fusion algebras, *Phys. Lett. B* **217** (1989), 47–251.
- [22] R.C. King, F. Toumazet and B.G. Wybourne, A finite subgroup of the exceptional Lie group G_2 , *J. Phys. A* **32** (1999), 8527–8537.
- [23] H. Kluberg-Stern, A. Morel and B. Petersson, Spectrum of lattice gauge theories with fermions from a $1/d$ expansion at strong coupling, *Nucl. Phys. B* **215** (1983), 527–554.
- [24] M. Koca and R. Koc, Automorphism groups of pure integral octonions, *J. Phys. A* **27** (1994), 2429–2442.
- [25] G. Kuperberg, Spiders for rank 2 Lie algebras, *Comm. Math. Phys.* **180** (1996), 109–151.
- [26] D.E. Littlewood, Group characters and the structure of groups, *Proc. London Math. Soc.* S2-39 (1934), 150.
- [27] J. Lopez Pena, S. Majid and K. Rietsch, Lie theory of finite simple groups and generalised Roth conjectures. Preprint, arXiv:1003.5611 [math.QA].
- [28] S. Morrison, E. Peters and N. Snyder, Categories generated by a trivalent vertex, *Selecta Mathematica* **23** (2017), 817–868.
- [29] S. Popa, An axiomatization of the lattice of higher relative commutants of a subfactor, *Invent. Math.* **120** (1995), 427–445.
- [30] N. Snyder and V. Ostrik, Diagrams and quantum G_2 at roots of unity, Special Session on Advances in Operator Algebras, AMS Spring Southeastern Sectional Meeting, Vanderbilt University, 14–15 April 2018.
- [31] M. Takesaki, *Theory of operator algebras. I*, Encyclopaedia of Mathematical Sciences **124**. Springer-Verlag, Berlin, 2002.
- [32] S. Uhlmann, R. Meinel and A. Wipf, Ward identities for invariant group integrals, *J. Phys. A* **40** (2007), 4367–4389.
- [33] D.B. Wales, Finite linear groups of degree seven II, *Pac. J. Math.* **34** (1970), 207–235.
- [34] A. Wassermann, Coactions and Yang-Baxter equations for ergodic actions and subfactors, in *Operator algebras and applications, Vol. 2, 203–236*, London Math. Soc. Lecture Note Ser. **136**, Cambridge Univ. Press, Cambridge, 1988.
- [35] B.H. Wellegehausen, A. Wipf and C. Wozar, Effective Polyakov loop dynamics for finite temperature G_2 gluodynamics, *Phys. Rev. D* **80** (2009), 065028.